



Comparative confinement studies in large stellarators

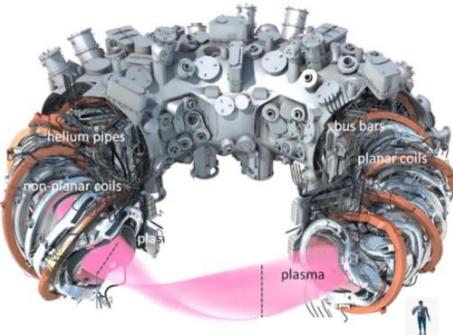
Hjördis Bouvain

Defense of the Master thesis
University of Greifswald

02.06.2023

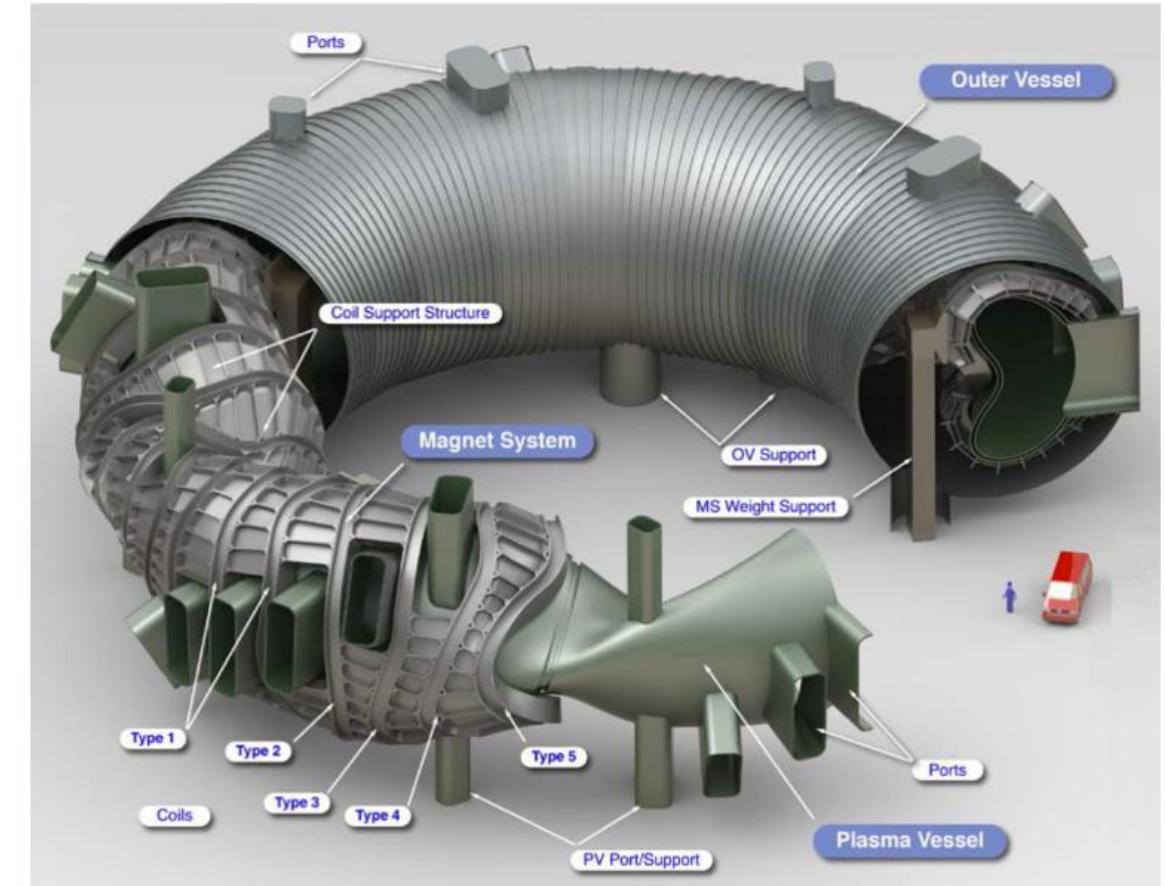


What can be learned from existing devices for future devices?



T. Klinger et al 2017 *Plasma Phys. Control. Fusion* **59** 014018

Extrapolating from today's fusion device to reactor size:



"HELIAS 5-B magnet system structure and maintenance concept", F. Schauer et al, *Fus. Eng. Des.* 88 (2013)

How to extrapolate from smaller to bigger scales



Wind Tunnel Test on Model Cessna - Laboratory Report, T.J. Sheng, 2018

Scaling from
model to full-
sized plane



https://commons.wikimedia.org/wiki/File:Cessna_172_%28D-EGUP%29_03.jpg

$$\text{Reynolds number: } R_e = \frac{\rho v d}{\mu}$$

medium density:	ρ
velocity:	v
characteristic length:	d
dynamic viscosity:	μ

Principle of similarity:
Two geometric similar bodies with
equal R_e possess equal flow physics.

How to extrapolate from smaller to bigger scales



Model size ten
times smaller

$$R_e = \frac{\rho v d}{\mu}$$



	full-sized plane	velocity scaling	medium: water
airfoil length	d (m)	1.6	0.16
velocity	v ($\frac{m}{s}$)	60	600
medium density	ρ ($\frac{kg}{m^3}$)	0.82	997.8
dynamic viscosity	μ ($10^{-5} Pa s$)	1.66	0.95
<hr/>			
Reynolds number R_e	4 700 000	4 700 000	4 700 000

How to extrapolate from smaller to bigger scales



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	full-sized plane	velocity scaling	medium: water
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Reynolds number R_e	4 700 000	4 700 000 316 144	4 700 000

How to extrapolate from smaller to bigger scales



Model size ten
times smaller

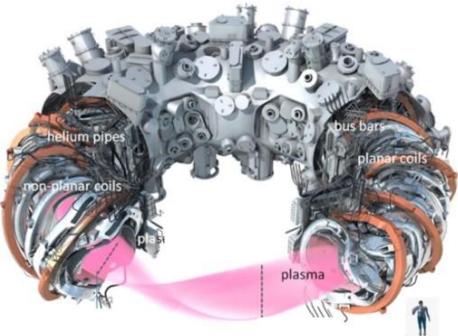
$$R_e = \frac{\rho v d}{\mu}$$



	full-sized plane	velocity scaling	medium: water
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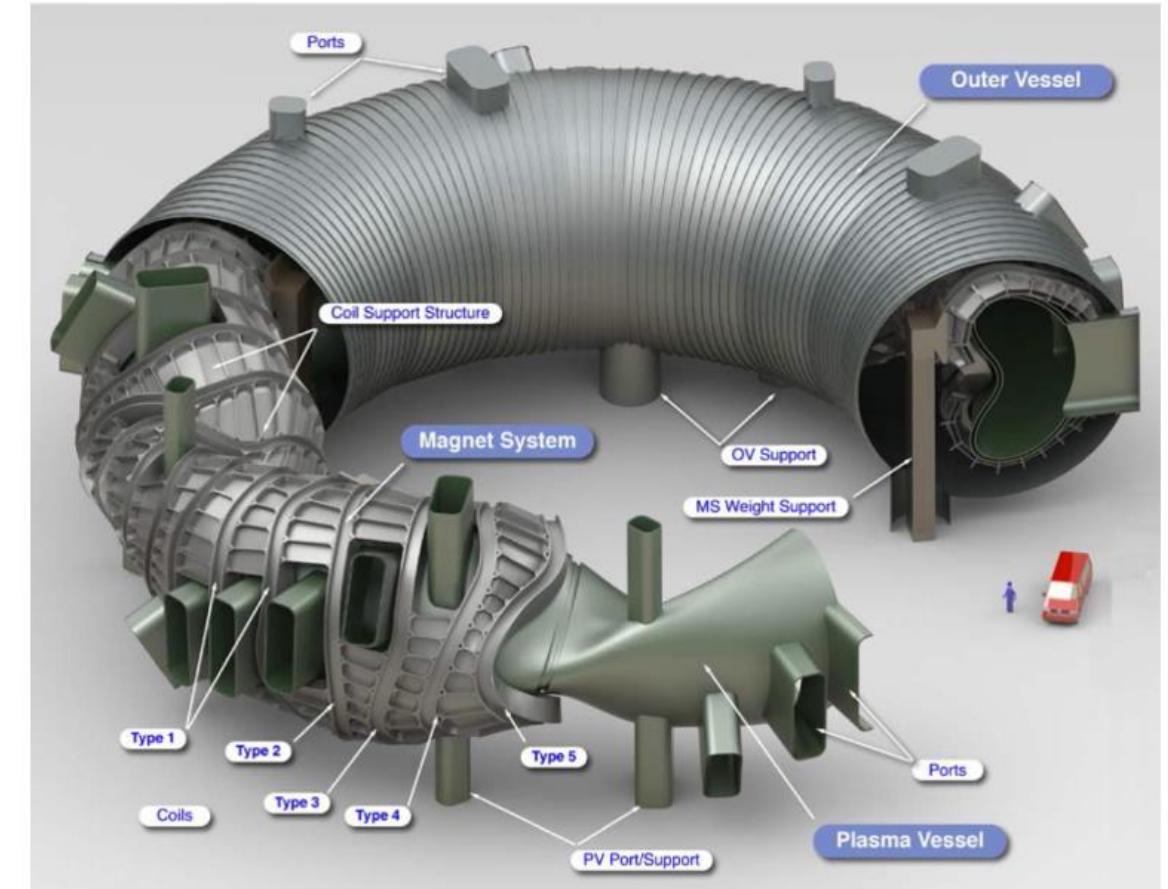
What can be learned from existing devices for future devices?



Extrapolating from today's fusion device to reactor size:



Scaling of the confinement with dimensionless parameters to extrapolate to larger machines





Extrapolating stellarators to ignition

Dimensionless parameters in fusion:

plasma beta $\beta = \frac{2\mu_0 n T}{B^2}$  Set by economical targets and physical limits

normalised
collisionality $\nu^* = \frac{R}{\ell} \frac{\nu_{ei}}{v_{th}} \propto \frac{n}{T^2}$

normalised
gyroradius $\rho^* = \frac{\sqrt{m_i T}}{a B Z e}$

plasma density:	n
plasma temperature:	T
magnetic field:	B



Extrapolating stellarators to ignition

Dimensionless parameters in fusion:

plasma beta

$$\beta = \frac{2\mu_0 n T}{B^2}$$



Set by economical targets and physical limits

normalised
collisionality

$$\nu^* = \frac{R}{l} \frac{\nu_{ei}}{v_{th}} \propto \frac{n}{T^2}$$



Set by limits in density and temperature

normalised
gyroradius

$$\rho^* = \frac{\sqrt{m_i T}}{a B Z e}$$

plasma density:	n
plasma temperature:	T
magnetic field:	B
major plasma radius:	R
electron-ion collision frequency:	ν_{ei}
thermal velocity:	v_{th}



Extrapolating stellarators to ignition

Dimensionless parameters in fusion:

plasma beta

$$\beta = \frac{2\mu_0 n T}{B^2}$$



Set by economical targets and physical limits

normalised
collisionality

$$\nu^* = \frac{R}{l} \frac{\nu_{ei}}{v_{th}} \propto \frac{n}{T^2}$$



Set by limits in density and temperature

normalised
gyroradius

$$\rho^* = \frac{\sqrt{m_i T}}{a B Z e}$$



Dependency on ρ^* gives insight on turbulence

$$\Omega_i \tau_E \propto \rho^{*\alpha_\rho} F(\beta, \nu^*, \dots)$$

plasma temperature:	T
magnetic field:	B
ion mass:	m_i
ion charge:	Z
minor plasma radius:	a
electron charge:	e

$\alpha_\rho = 1 \rightarrow$ gyro-Bohm scaling

$\alpha_\rho = 0 \rightarrow$ Bohm scaling

$\alpha_\rho = -1 \rightarrow \lambda \gg a$



Confinement

Global energy confinement time:

$$\tau_E = \frac{W_{dia}}{P_{heat} - \frac{dW_{dia}}{dt}}$$

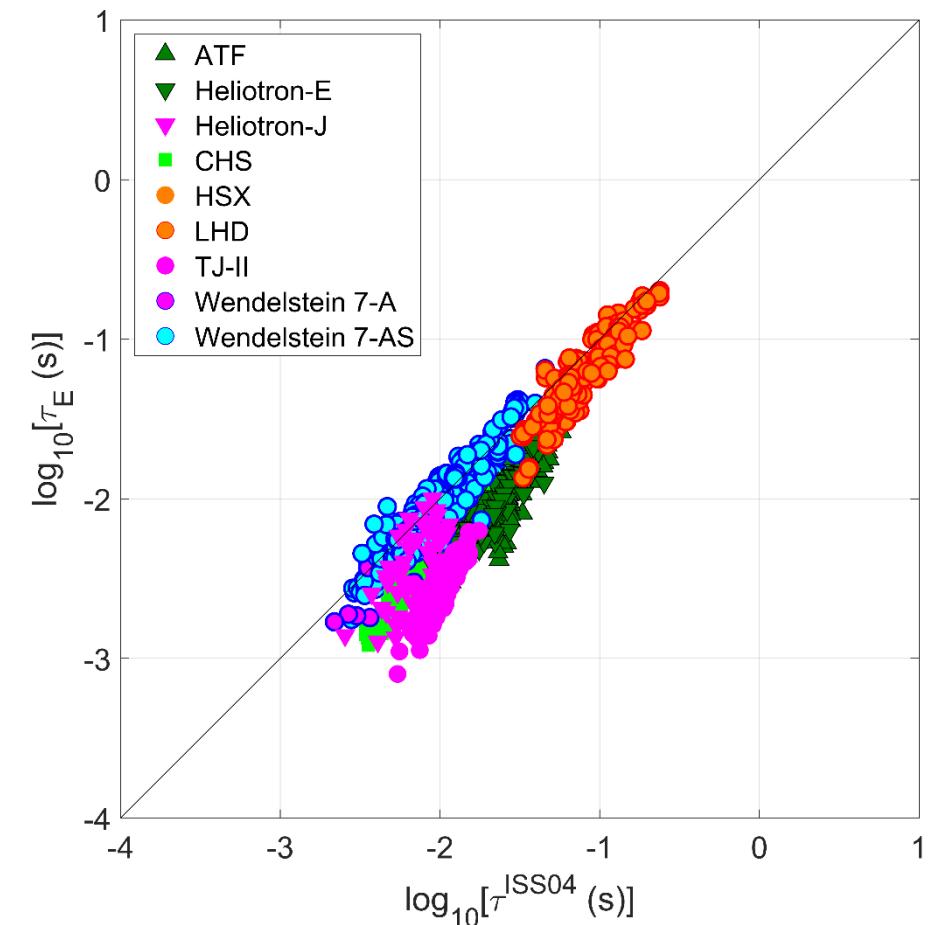
Confinement

Global energy confinement time:

$$\tau_E = \frac{W_{dia}}{P_{heat} - \frac{dW_{dia}}{dt}}$$

Empirical scaling law ISS04: [1]

$$\tau^{\text{ISS04}} = 0.134 a^{2.28} R^{0.64} B^{0.84} t_2^{\frac{0.41}{3}} P_{heat}^{-0.61} n_e^{0.54}$$



Recreated from A. Dinklage et al 2007 Nucl. Fusion **47** 1265

[1] H. Yamada et al 2005 Nucl. Fusion **45** 1684

Confinement

Global energy confinement time:

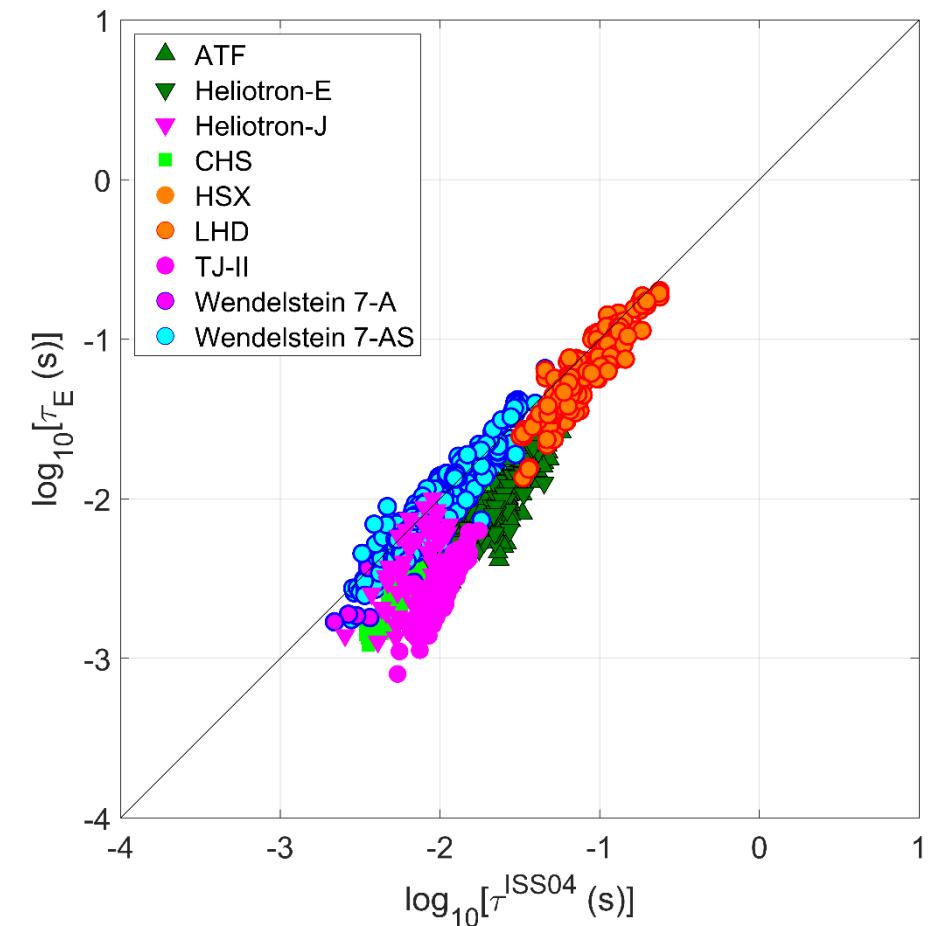
$$\tau_E = \frac{W_{dia}}{P_{heat} - \frac{dW_{dia}}{dt}}$$

Empirical scaling law ISS04: [1]

$$\tau^{\text{ISS04}} = 0.134 a^{2.28} R^{0.64} B^{0.84} t_2^{\frac{0.41}{3}} P_{heat}^{-0.61} n_e^{0.54}$$

Renormalisation factor:

$$f_{ren} = \frac{\tau_E}{\tau_{scaling}}$$

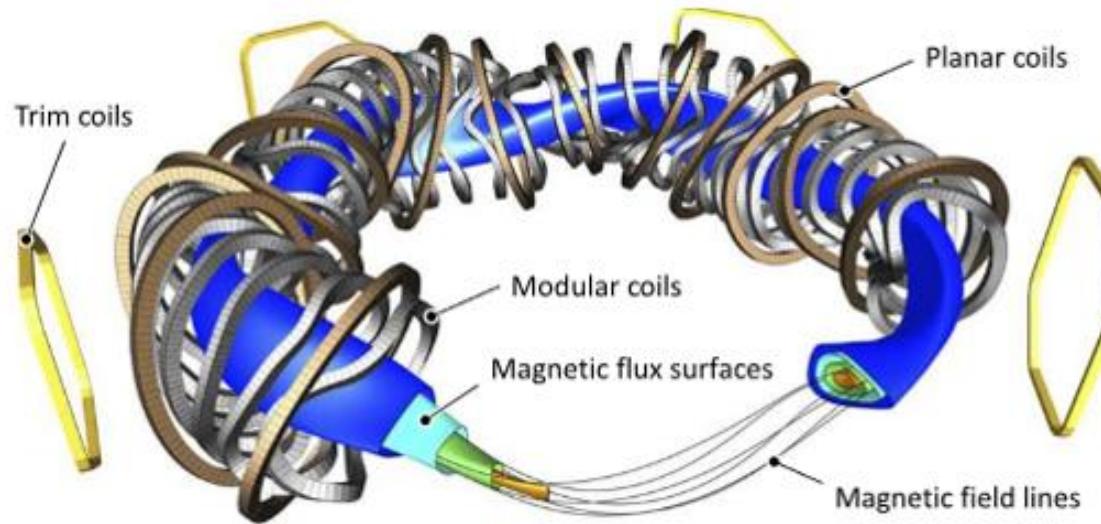


Recreated from A. Dinklage et al 2007 Nucl. Fusion **47** 1265

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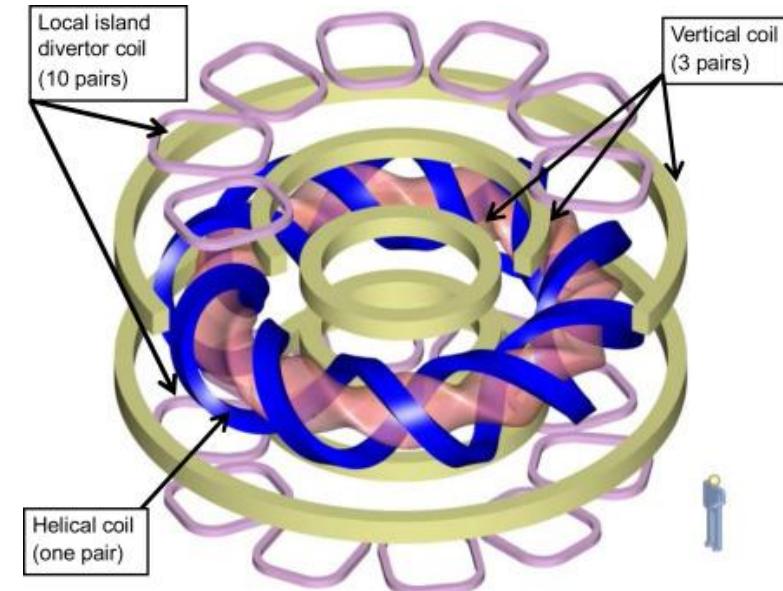
Machines

W7-X



R. C. Wolf, A. Alonso, S. Äkäslompolo, et al., "Performance of Wendelstein 7-X stellarator plasmas during the first divertor operation phase"

LHD



"Magnetic Fusion Energy From Experiments to Power Plants", George H. Neilson

W7-X

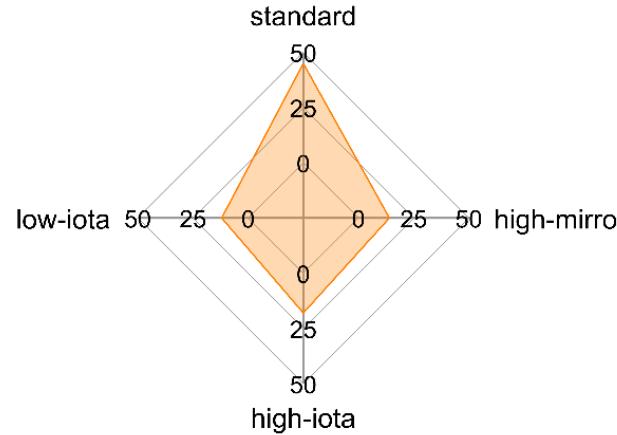
major plasma radius R (m)	5.5
minor plasma radius a (m)	0.55
plasma volume V (m^3)	30
magnetic field B (T)	2.5

LHD

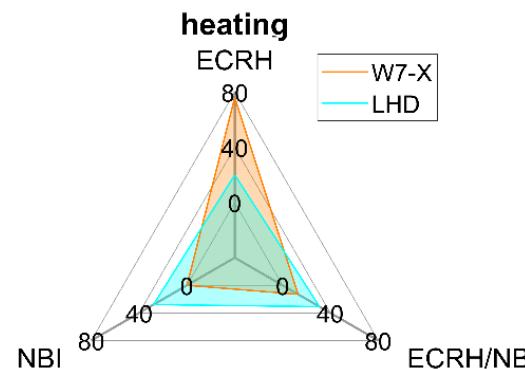
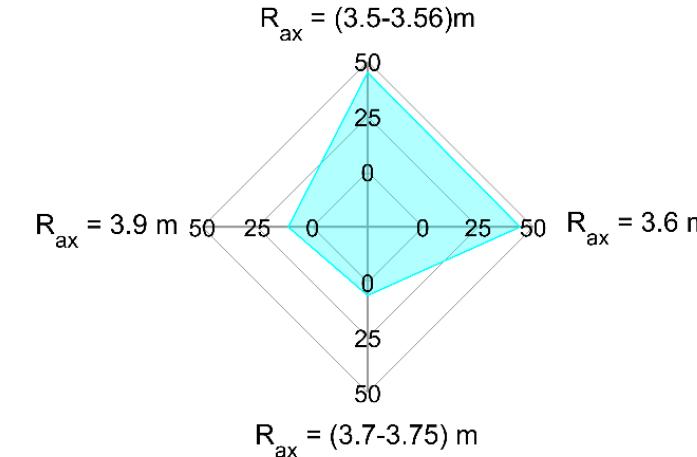
5.5	3.9
0.55	0.65
30	30
2.5	< 3

Dataset used

W7-X: configuration

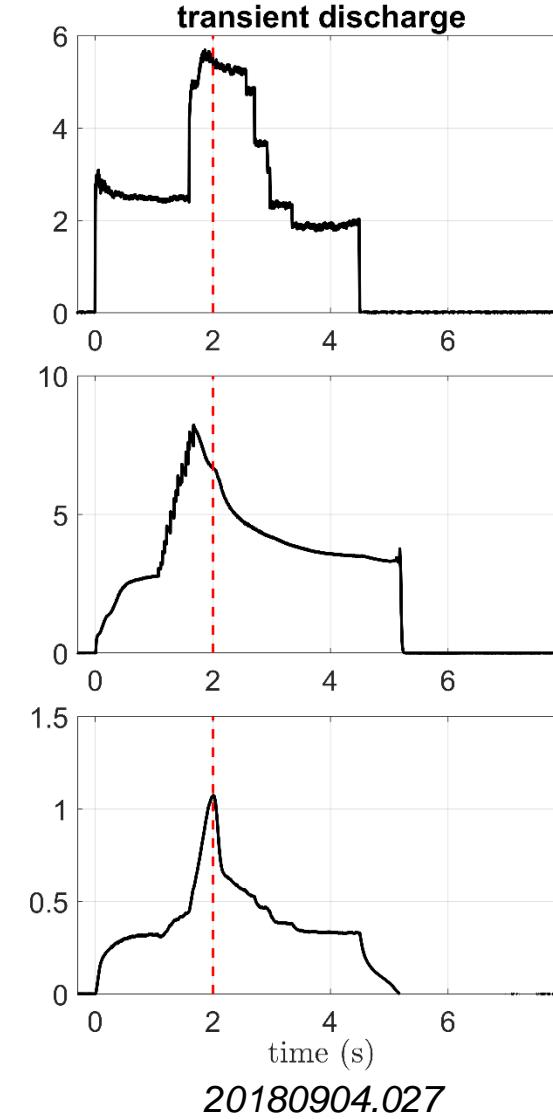
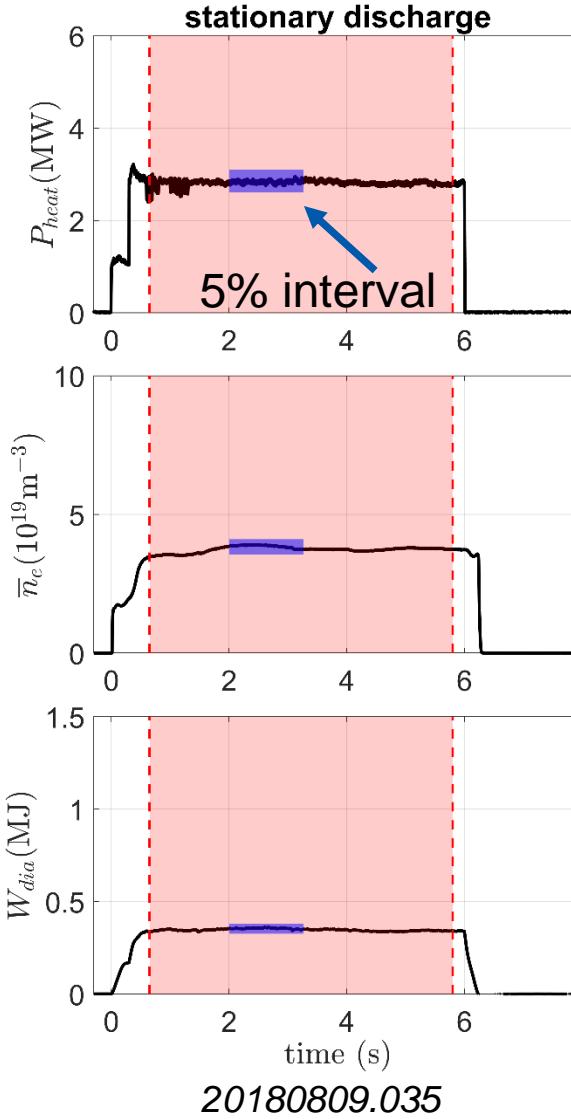


LHD: configuration



- **Across both machines:**
168 discharges
- **W7-X dataset:**
89 discharges from OP1.2b (2018)
- **LHD dataset:**
79 discharges from three different OP (2017-2022)

Calculation of τ_E and τ^{ISS04}



$$\tau_E = \frac{W_{dia}(t)}{P_{heat}(t) - \frac{dW_{dia}}{dt}}$$

$$\tau^{ISS04} = 0.134 a^{2.28} R^{0.64} B^{0.84} t_2^{0.41} P_{heat}(t)^{-0.61} n_e(t)^{0.54}$$

“Stationary“ discharges:

Employing average values of P_{heat} , n_e and W_{dia} during the stationary phase (red shaded region)

“Transient“ discharges:

Employing values of P_{heat} , n_e and W_{dia} at instance when W_{dia} is maximal (red dashed line)

Calculation of the renormalisation factor for temporal analysis

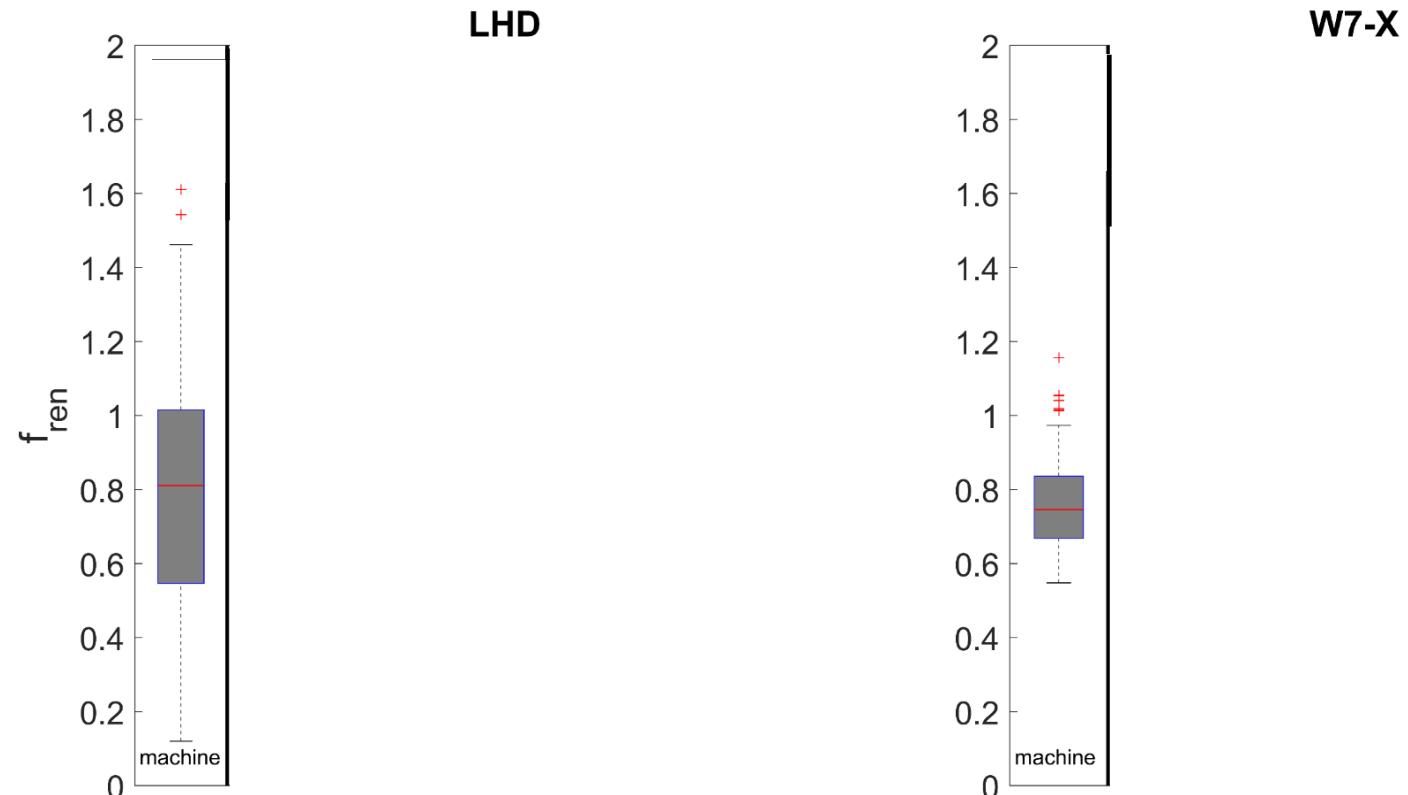
$$\tau_E = \frac{W_{dia}}{P_{heat} - \frac{dW_{dia}}{dt}} \quad \& \quad \tau^{ISS04} = 0.134 a^{2.28} R^{0.64} B^{0.84} t_2^{0.41} P_{heat}^{-0.61} n_e^{0.54}$$

$$\frac{dW^{ISS04}(t)}{dt} + \frac{W^{ISS04}(t)}{\tau^{ISS04}(t)} = P_{heat}(t)$$

Using integrating factors: $I(t) = \exp \left(\int_{t_0}^t \frac{dt'}{\tau^{ISS04}(t')} \right)$

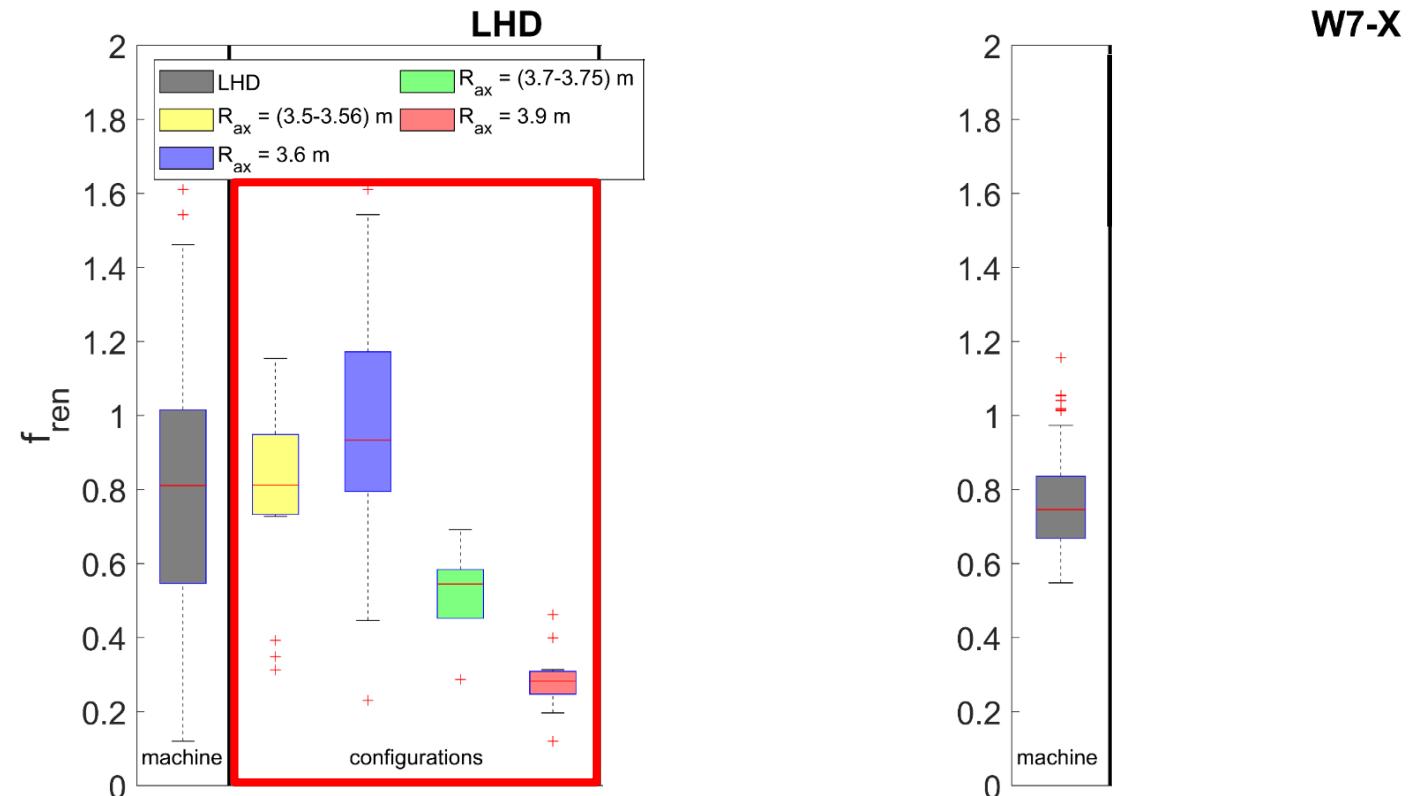
$$W^{ISS04}(t') = \frac{\int_{t_0}^t dt' I(t') P_{heat}(t')}{I(t')}$$

Existence of dependencies in τ_E



LHD and W7-X datasets show similar confinement quality

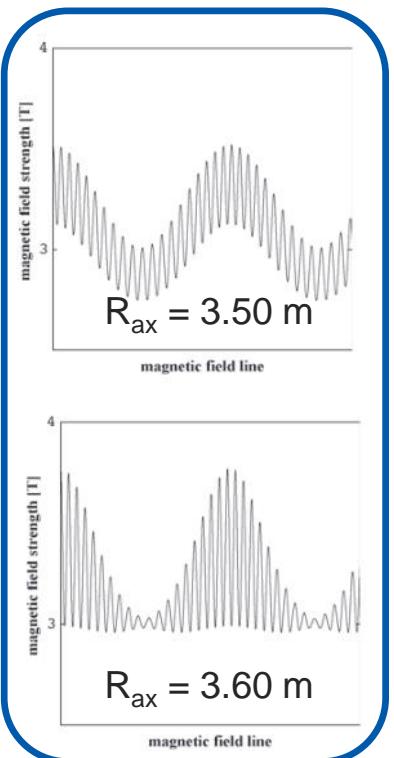
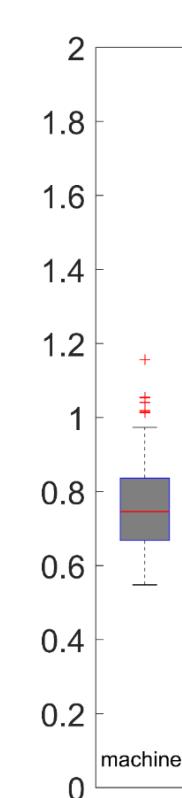
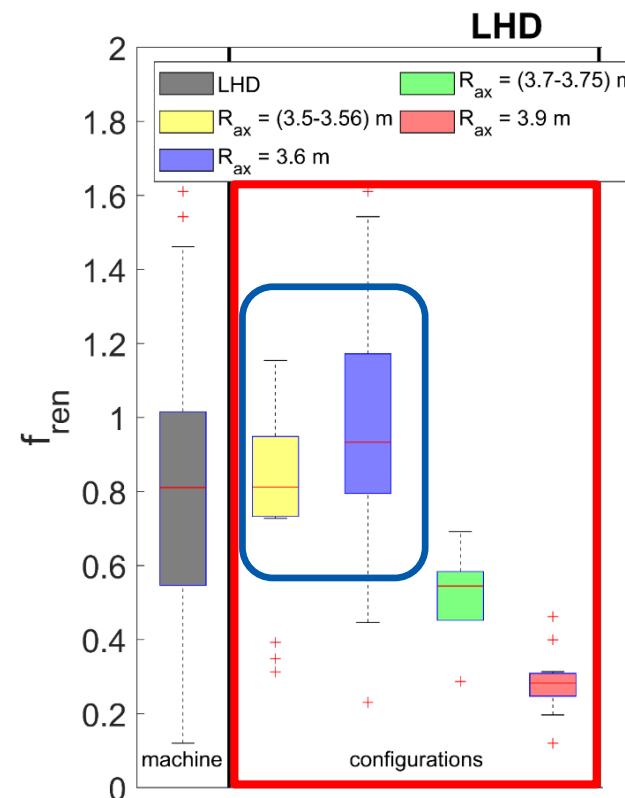
Existence of dependencies in τ_E



Dependency of the confinement on magnetic configuration in LHD dataset: optimal at $R_{ax} = 3.60$ m [1]

[1] S. Murakami et al 2002, *Nucl. Fusion* **42** L19

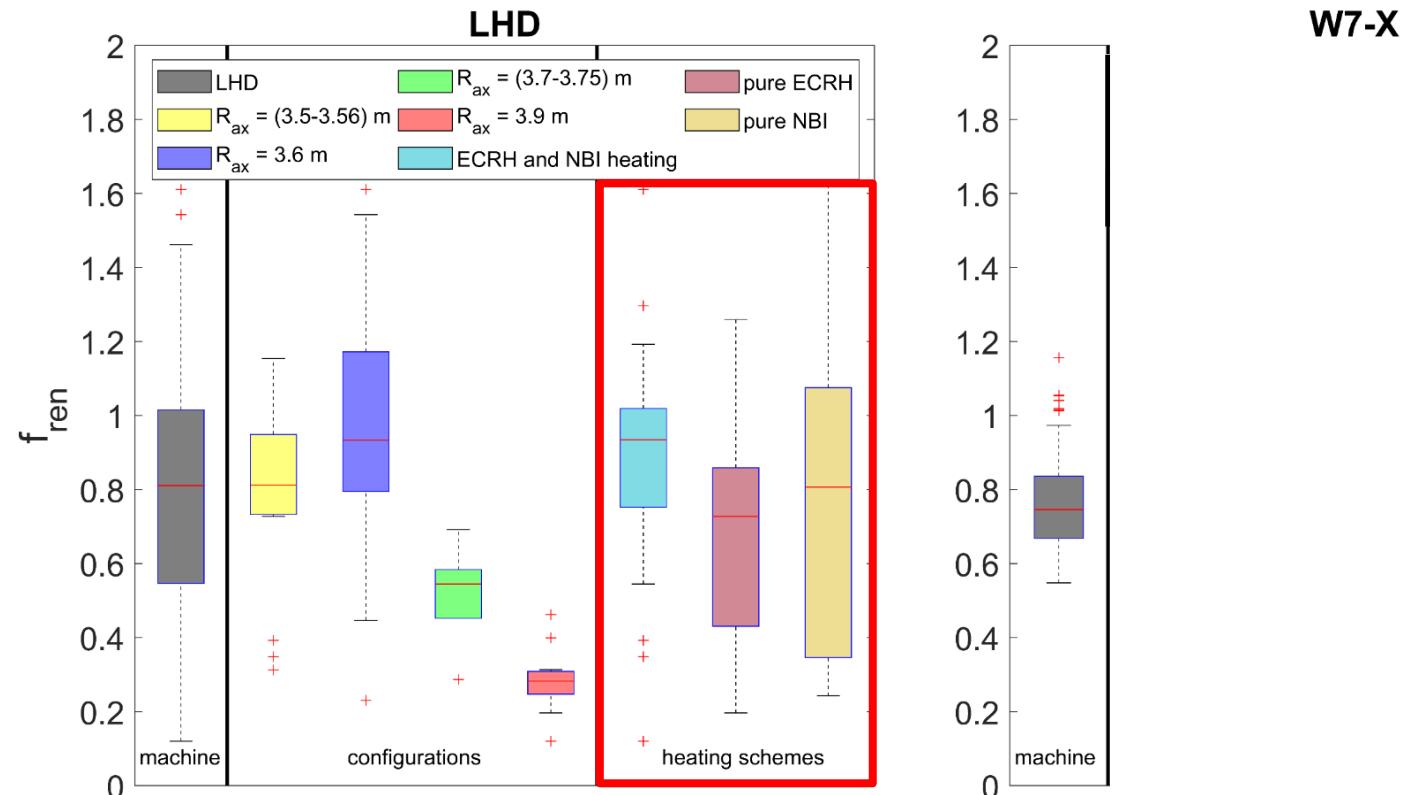
Existence of dependencies in τ_E



Dependency of the confinement on magnetic configuration in LHD dataset: optimal at $R_{ax} = 3.60 m$ [1] due to improved neoclassical confinement

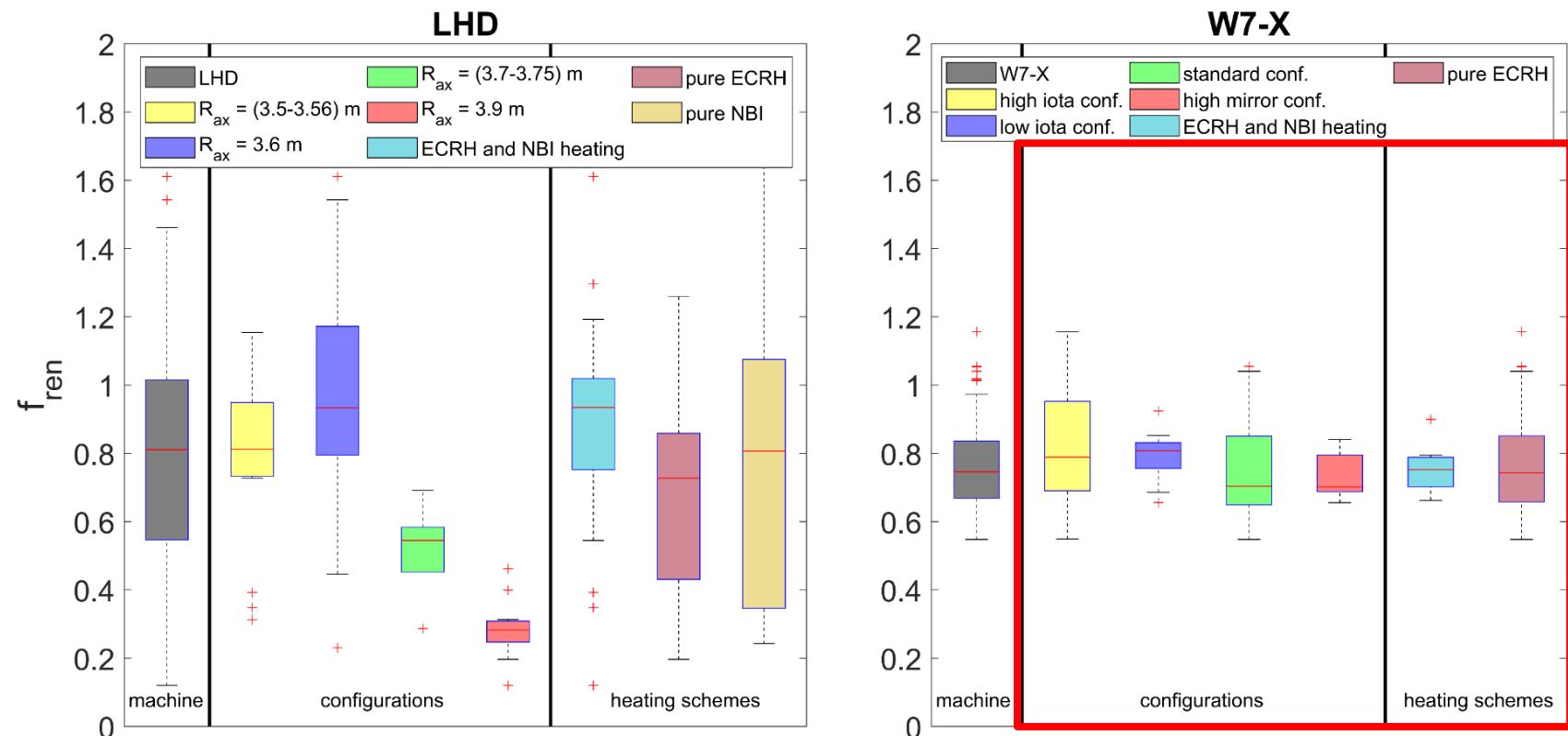
[1] S. Murakami et al 2002, *Nucl. Fusion* **42** L19

Existence of dependencies in τ_E



No dependency of the confinement on the heating method in LHD dataset

Existence of dependencies in τ_E

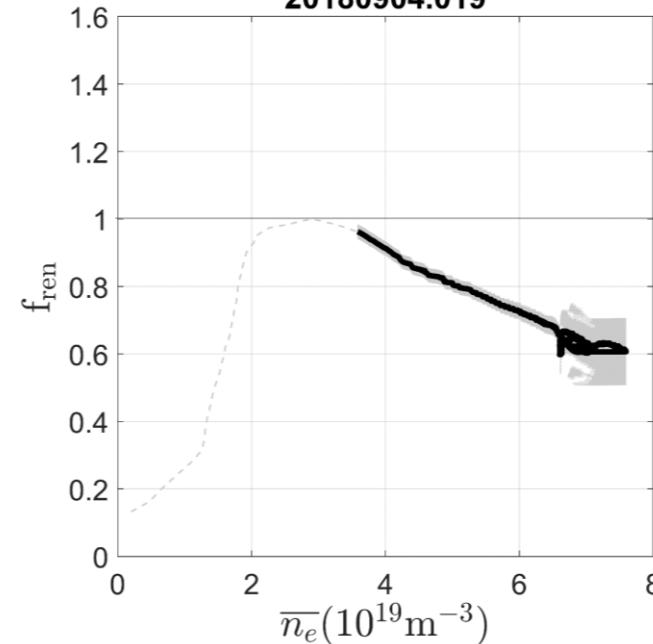


**No dependency on magnetic configuration or heating method
in W7-X dataset**

Density dependency of confinement in typical ECR heated discharges

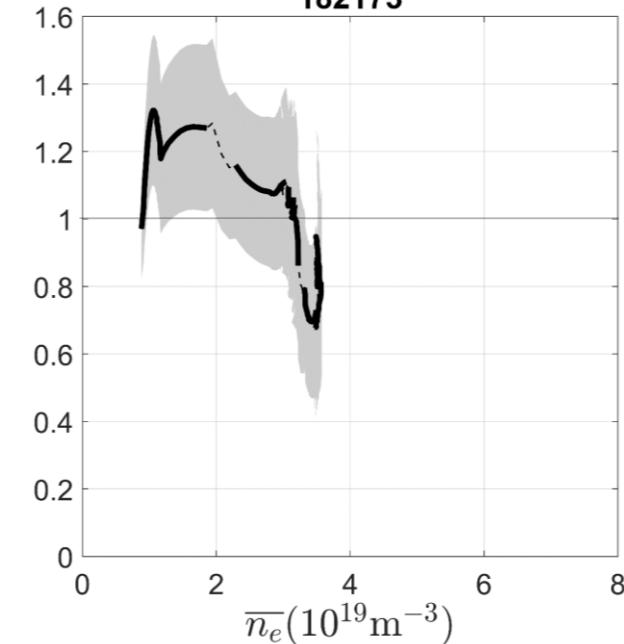
W7-X

20180904.019



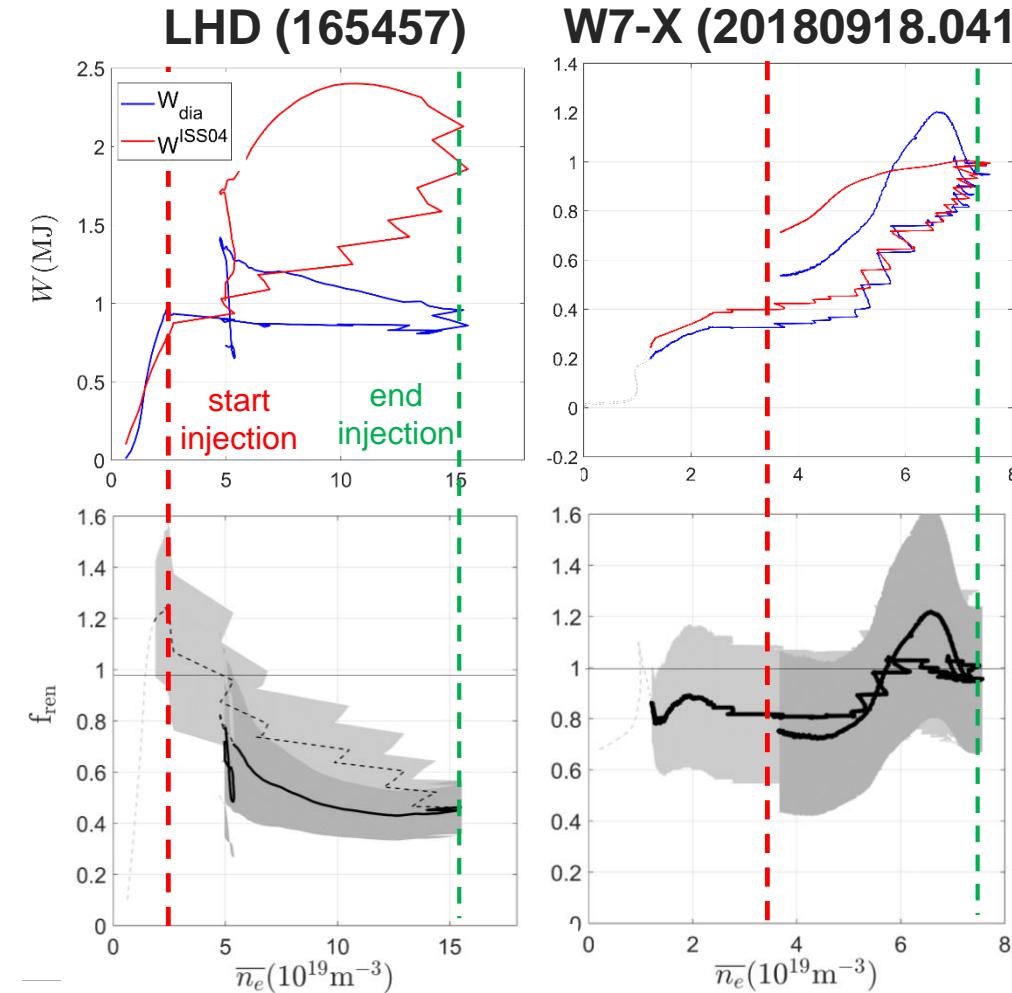
LHD

182173



Linear decrease of f_{ren} for increasing density in purely ECR heated stationary discharges

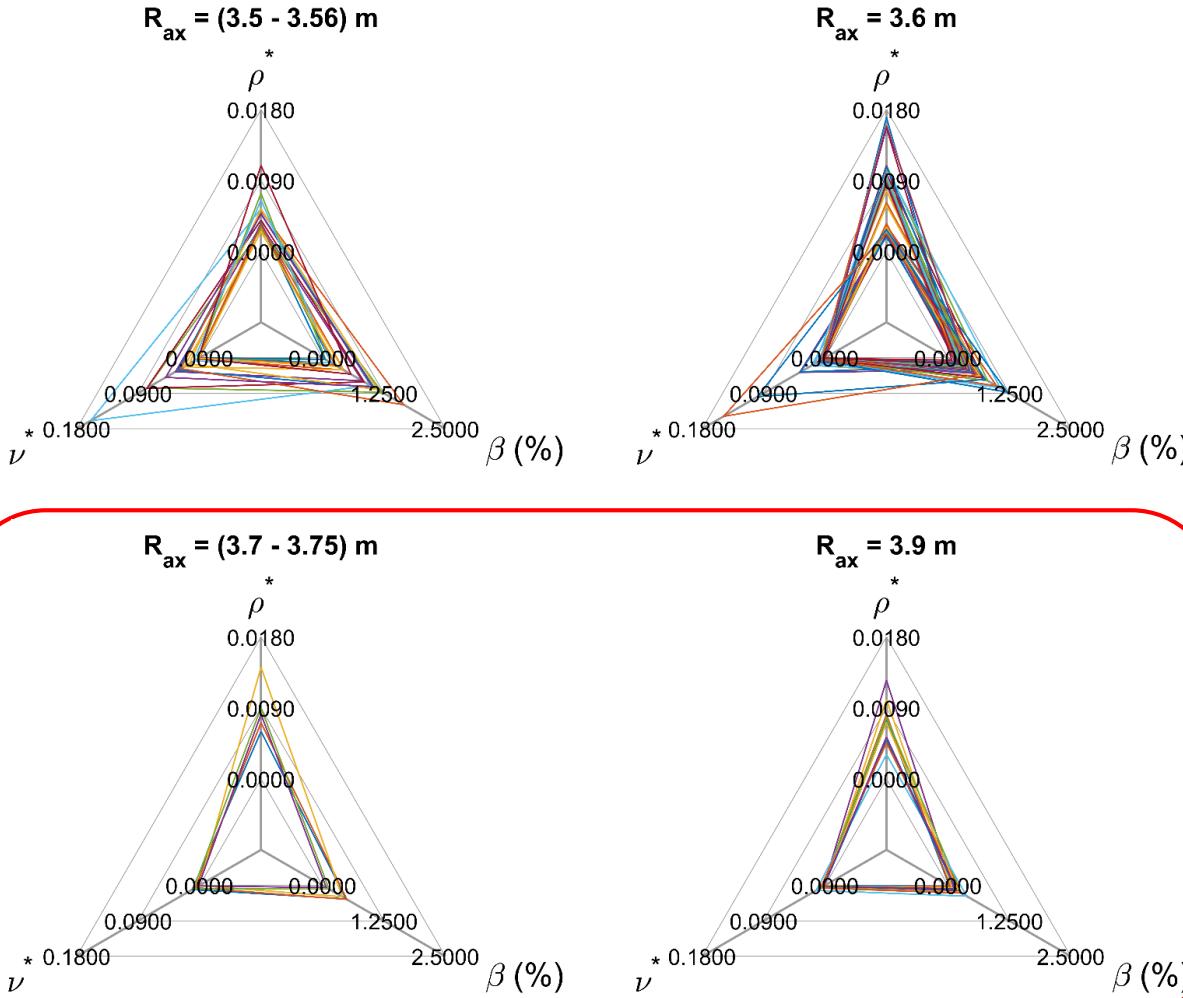
Machine dependency on confinement improvement by pellet-injection



Pellet injections lead to enhancement of plasma energy

Confinement improvement different for LHD and W7-X

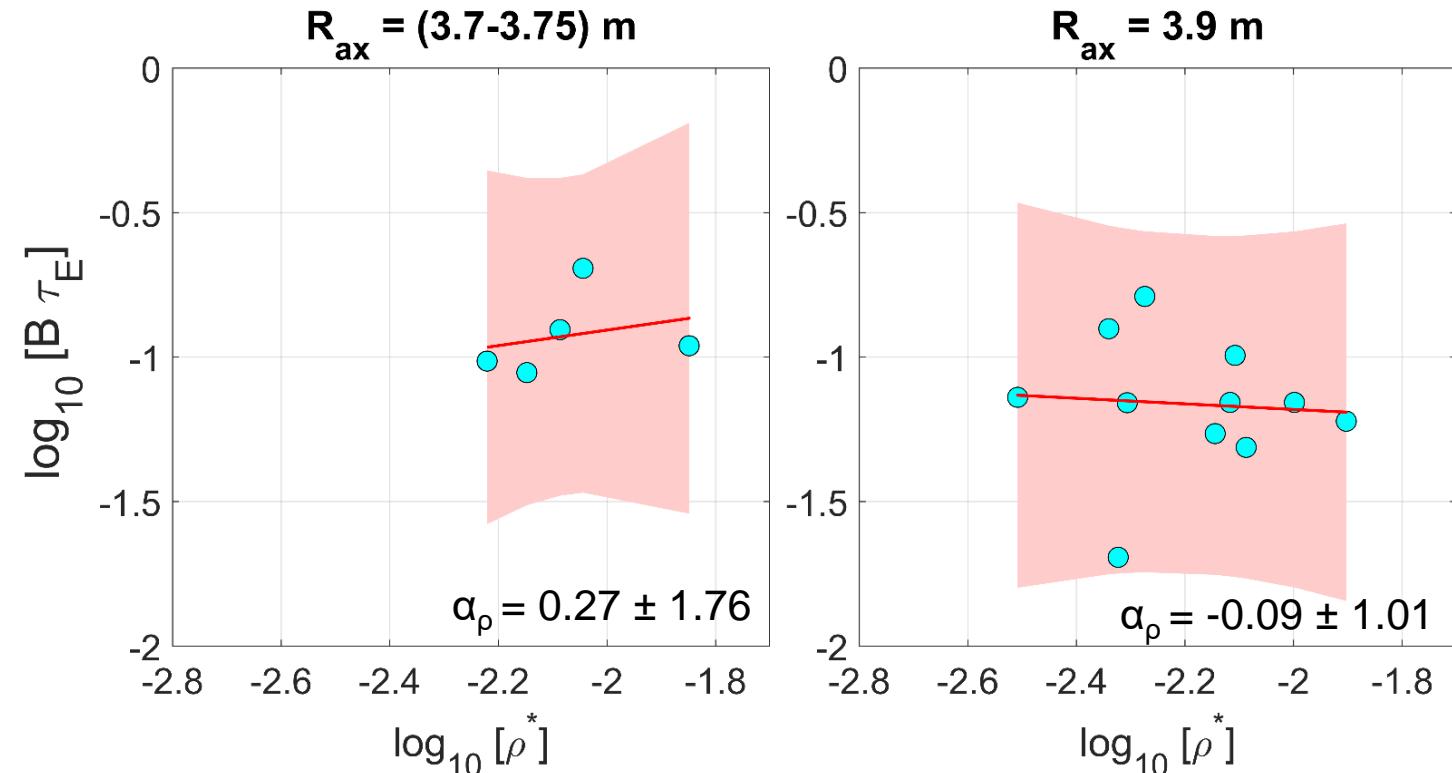
Scaling with dimensionless parameters of the LHD dataset



Possible scaling with one dimensionless parameter:

Configuration	β	ρ^*	ν^*
$R_{ax} = (3.5 - 3.56) \text{ m}$			
$R_{ax} = 3.6 \text{ m}$			
$R_{ax} = (3.7 - 3.75) \text{ m}$		X	
$R_{ax} = 3.9 \text{ m}$		X	

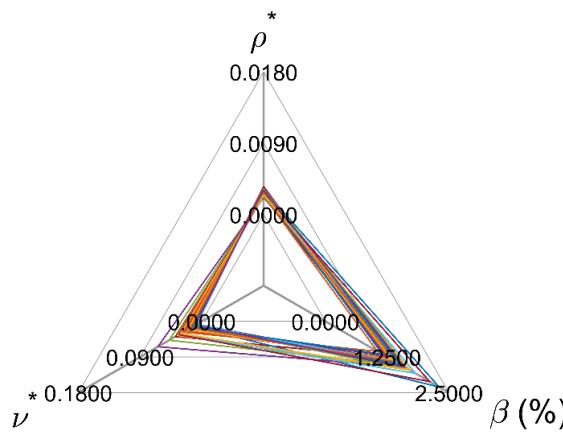
Confinement scaling with ρ^* of the LHD dataset



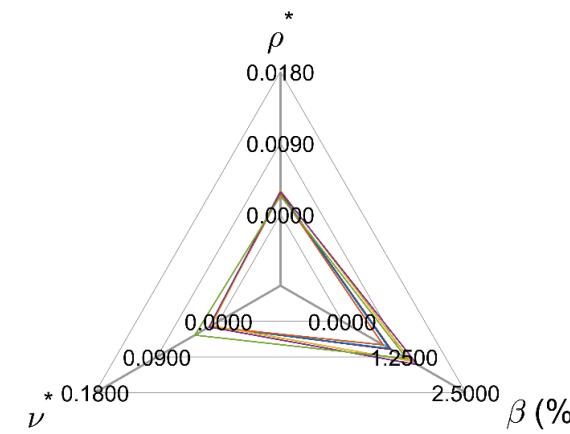
No reliable scaling of the confinement time with ρ^* in the LHD dataset due to the large error and small dataset

Scaling with dimensionless parameters of the W7-X dataset

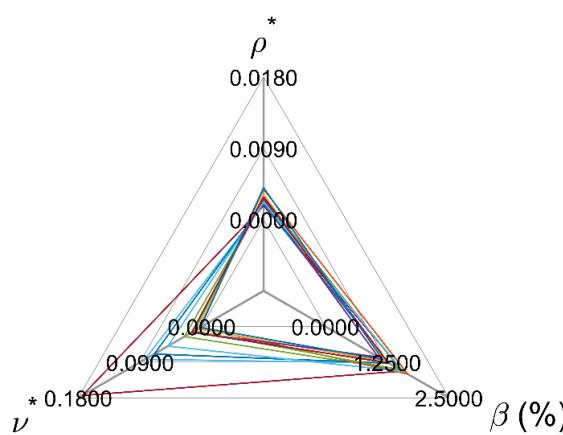
standard configuration



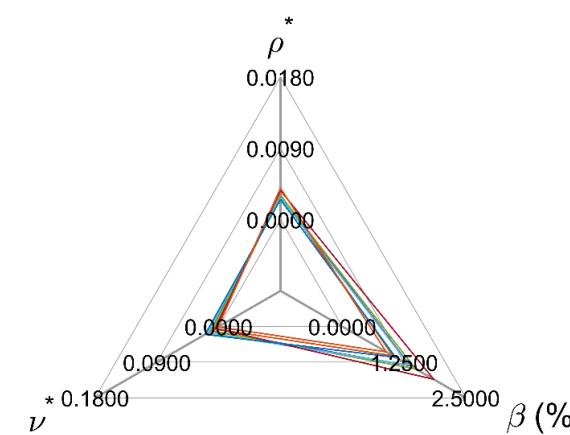
high mirror configuration



high iota configuration



low iota configuration



Possible scaling with one dimensionless parameter:

Configuration	β	ρ^*	ν^*
standard	✗		
high-mirror	✗		
high-iota			✗
low-iota	✗		

No ρ^* scaling possible:
further ρ^* scan experiments in
W7-X needed

Proposal: Variation of magnetic field for ρ^* - scans in W7-X

High- and low magnetic field for operation with O2 at high fields and X3 at low fields for similar deposition profiles:

$$B_H = 2.5 \text{ T} \quad B_L = \frac{2}{3} B_H = 1.67 \text{ T}$$

Keeping β and ν^* constant over both operation scenarios by matching high and low field density n and temperature T using the magnetic field:

$$\left. \begin{array}{l} \beta \propto \frac{nT}{B^2} = \text{const} \\ \nu^* \propto \frac{na}{T^2} = \text{const} \end{array} \right\} \quad \begin{array}{l} \frac{n_L}{n_H} = \left(\frac{B_L}{B_H}\right)^{\frac{4}{3}} = 0.58 \\ \frac{T_L}{T_H} = \left(\frac{B_L}{B_H}\right)^{\frac{2}{3}} = 0.76 \end{array} \quad \xrightarrow{\rho^* \propto \frac{\sqrt{T}}{aB}} \quad \text{possible } \rho^* \text{ variation:}$$

$$\boxed{\frac{\rho_H^*}{\rho_L^*} = \left(\frac{B_H}{B_L}\right)^{-\frac{2}{3}} = 0.76}$$

Comparison of low and high field discharges provides insight in ρ^* - scaling

Summary

- Investigation of confinement quality for selected datasets from LHD and W7-X with respect to magnetic configuration, heating method and discharge scenario
- Unified approach for the assessment of LHD and W7-X confinement data
- Use of integrating factors for a temporal analysis of transient confinement

	LHD dataset	W7-X dataset
Magnetic configuration	Improved τ_E at $R_{ax} = 3.6m$ (reduced neoclassic transport)	No dependency found
Heating methods	No apparent dependency found	No dependency found
Pellet injection	Enhanced W_{dia} with $f_{ren} < 1$	Enhanced W_{dia} with $f_{ren} > 1$

Outlook: ρ^* - scan experiment in W7-X (achievable by magnetic field variation)