



Study of finite Larmor radius (FLR) physics in tokamaks and stellarators with the gyrokinetic turbulence code GENE-X

José Capitán, Philipp Ulbl, Baptiste J. Frei, Frank Jenko

HEPP Introductory Talk



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SPAIN

Madrid

Garching

Úbeda



SPAIN





Bachelor → Physics

Master → thesis in **CIEMAT**
doing simulations with stella in
W7-X high mirror and **CIEMAT-QI**
configs.

SPAIN



SPAIN



BEST THING TO DO IN SPAIN IS ~~LEAVING~~ EATING

SPAIN



BEST THING TO DO IN SPAIN IS ~~LEAVING~~ EATING



Spanish
omelette

SPAIN



BEST THING TO DO IN SPAIN IS ~~LEAVING~~ EATING



Spanish
omelette



Ham
“serrano”



SPAIN



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Paella right way to do it? :(



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BEST THING TO DO IN SPAIN IS ~~LEAVING~~ EATING



Spanish
omelette



Ham
“serrano”



Paella right way to do it? :(



Olive oil



SPAIN



random fact
the king of Spain is also



random fact

the king of Spain is also the king of Jerusalem 😐



???

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Madrid

Úbeda

Garching

random fact

the king of Spain is also the king of Jerusalem 😐



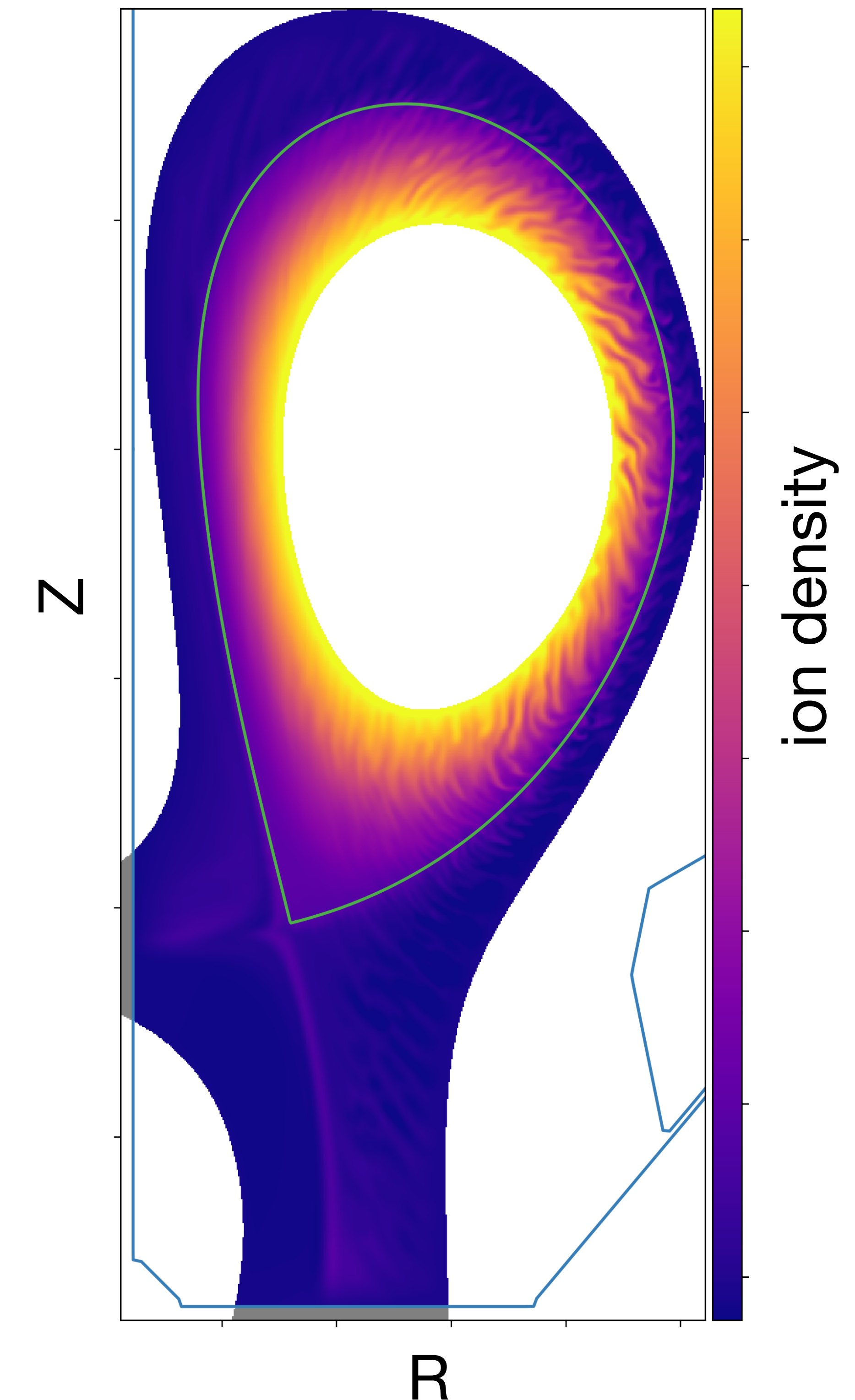
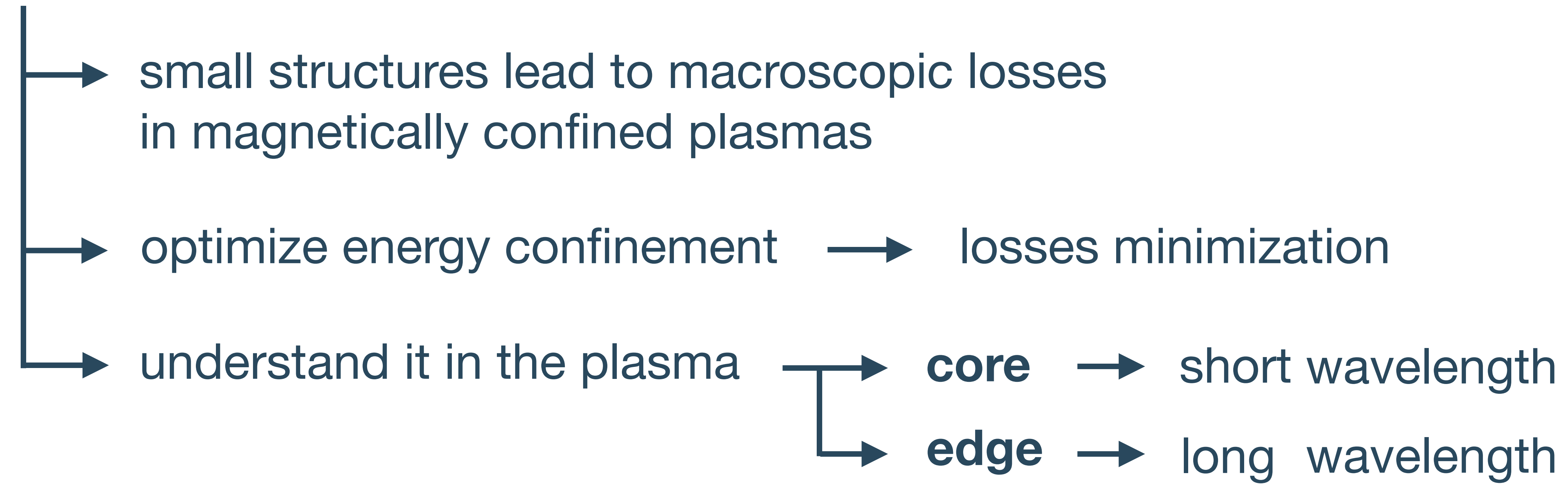
???

now blackouts too !!! 😊



PhD project

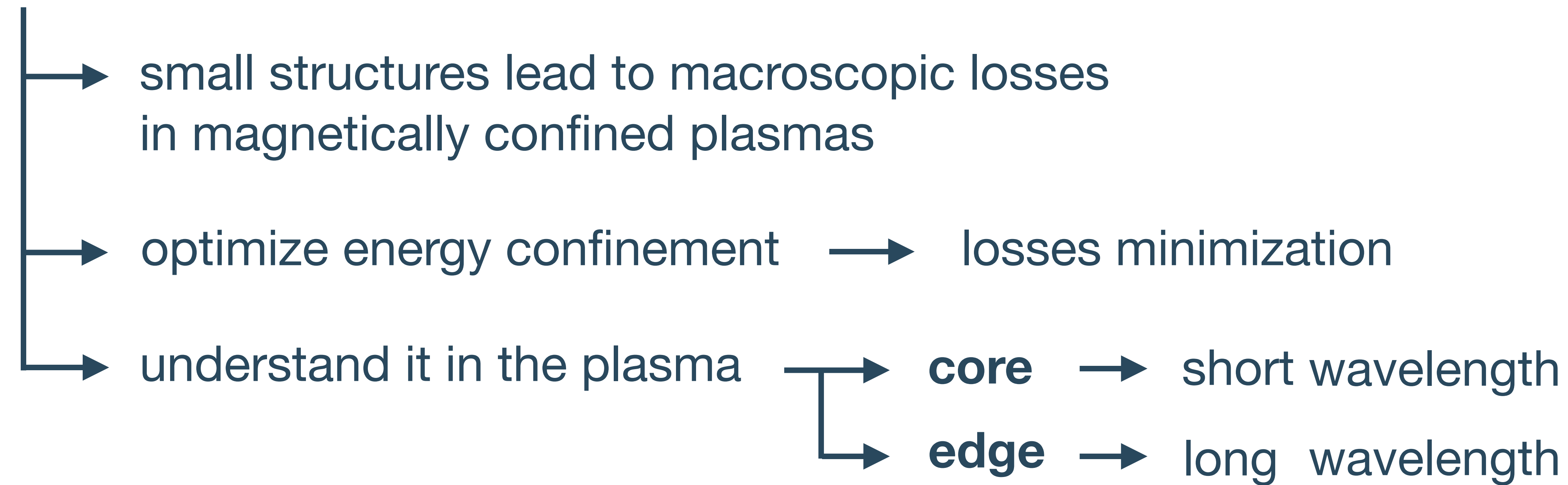
MICRO TURBULENCE



Motivation

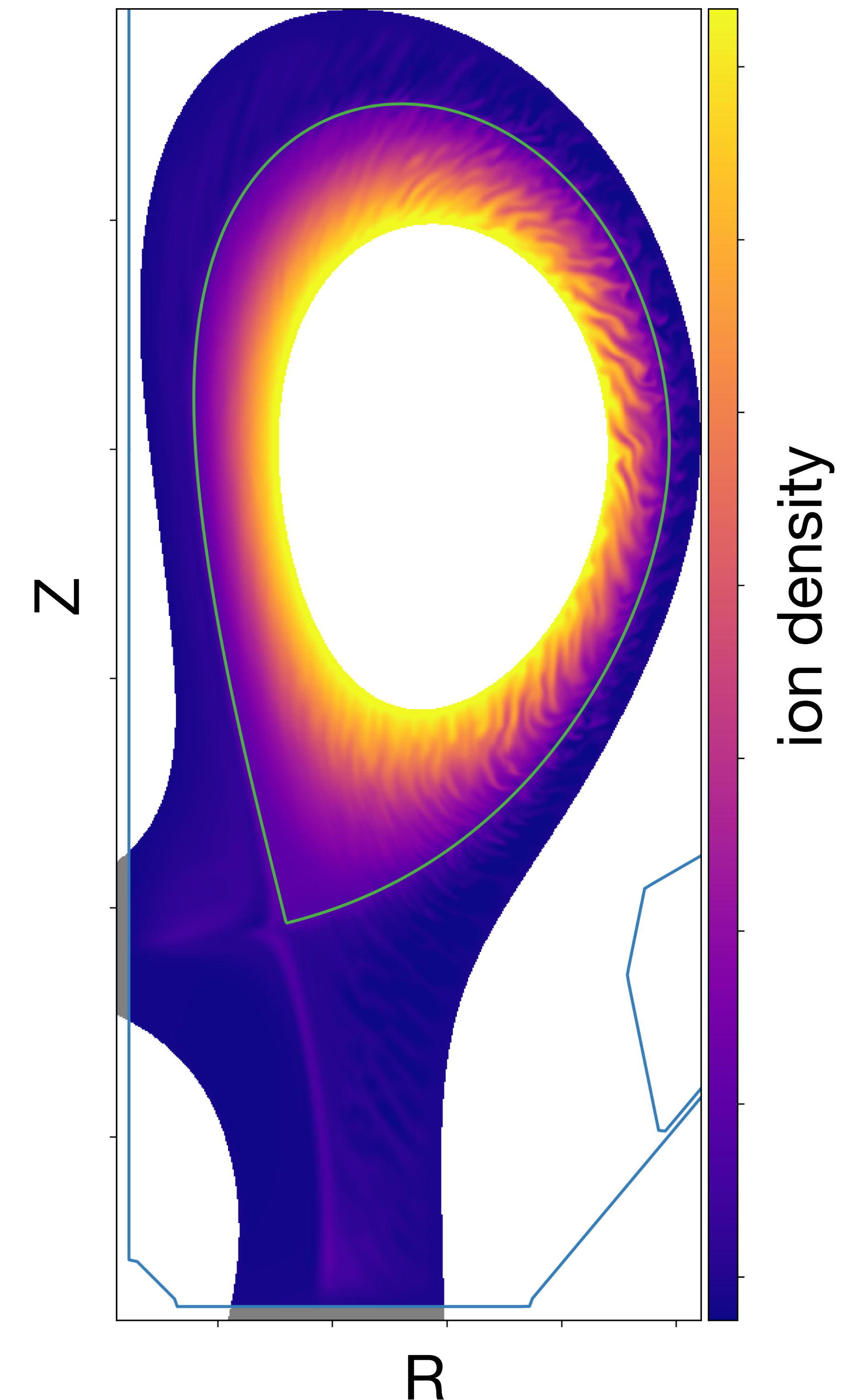


MICRO TURBULENCE



GENE-X

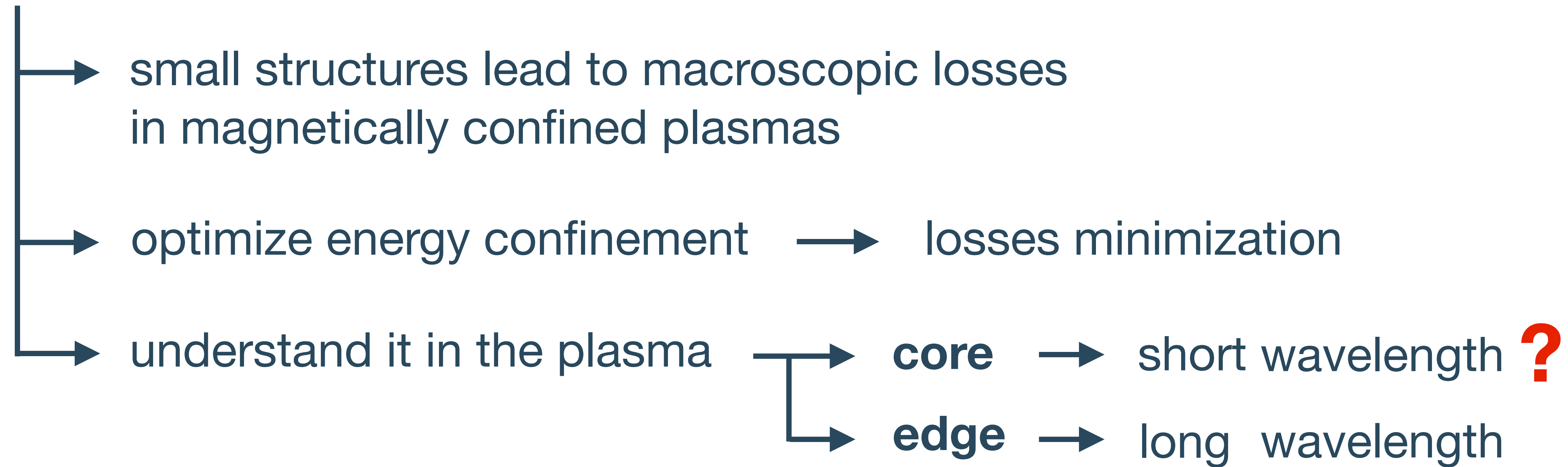
- Eulerian code that solves the gyrokinetic Vlasov eq. on a grid
- collisional, full- f , EM gyrokinetic turbulence model
- addresses the complexities of **edge turbulence** simulations



Motivation



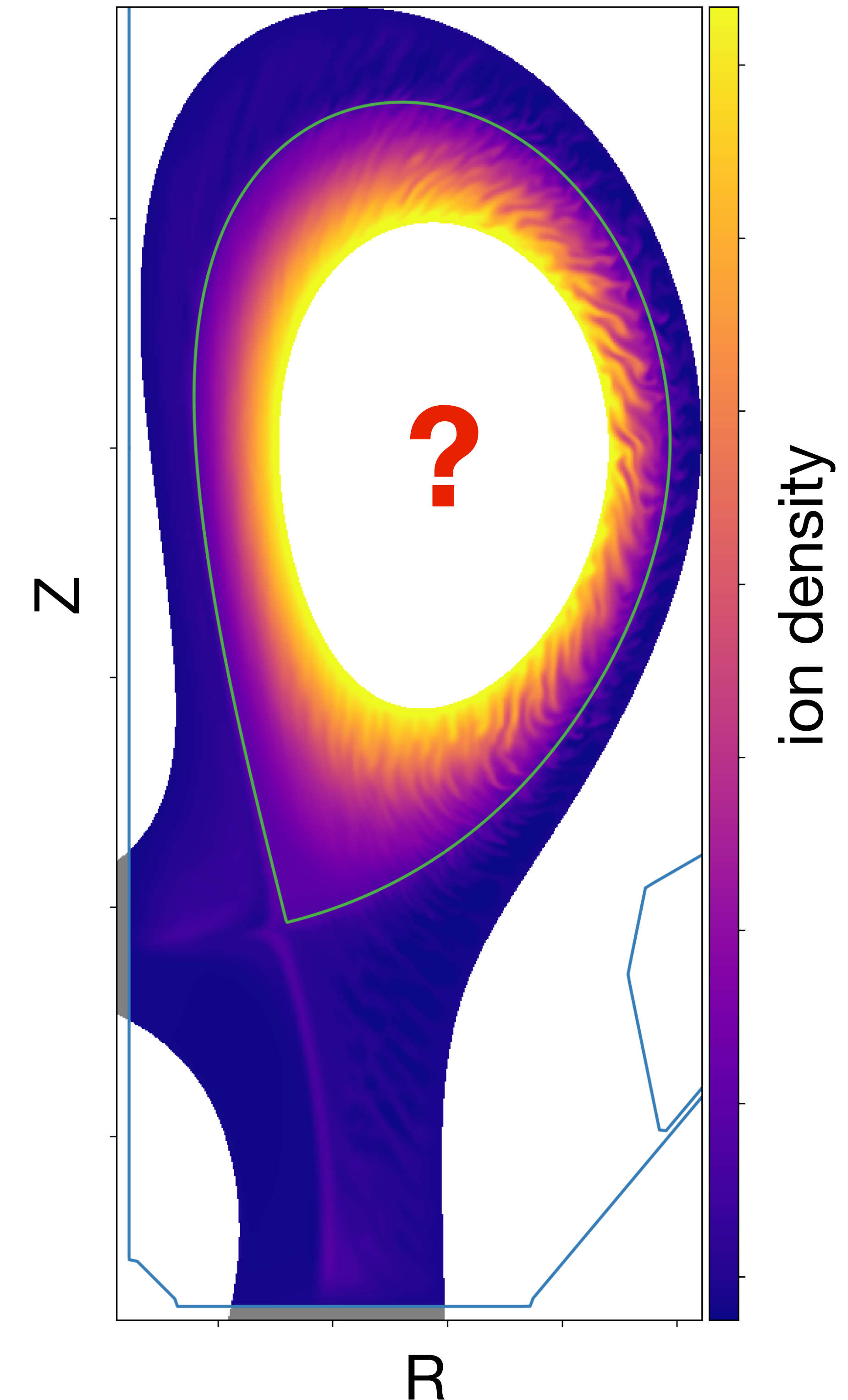
MICRO TURBULENCE



GENE-X

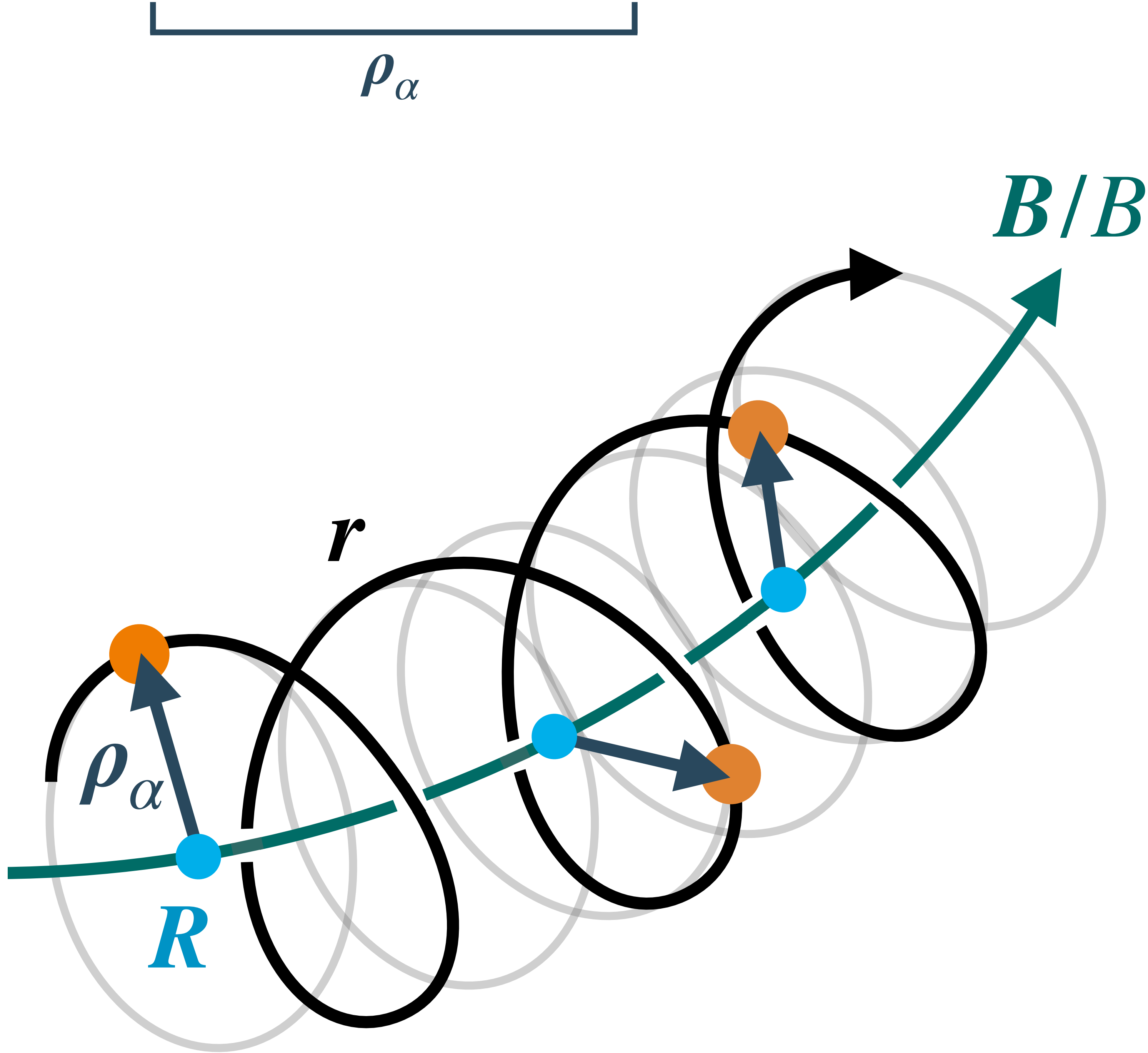
→ optimized for long wavelength turbulence!

- Eulerian code that solves the gyrokinetic Vlasov eq. on a grid
- collisional, full- f , EM gyrokinetic turbulence model
- addresses the complexities of **edge turbulence** simulations



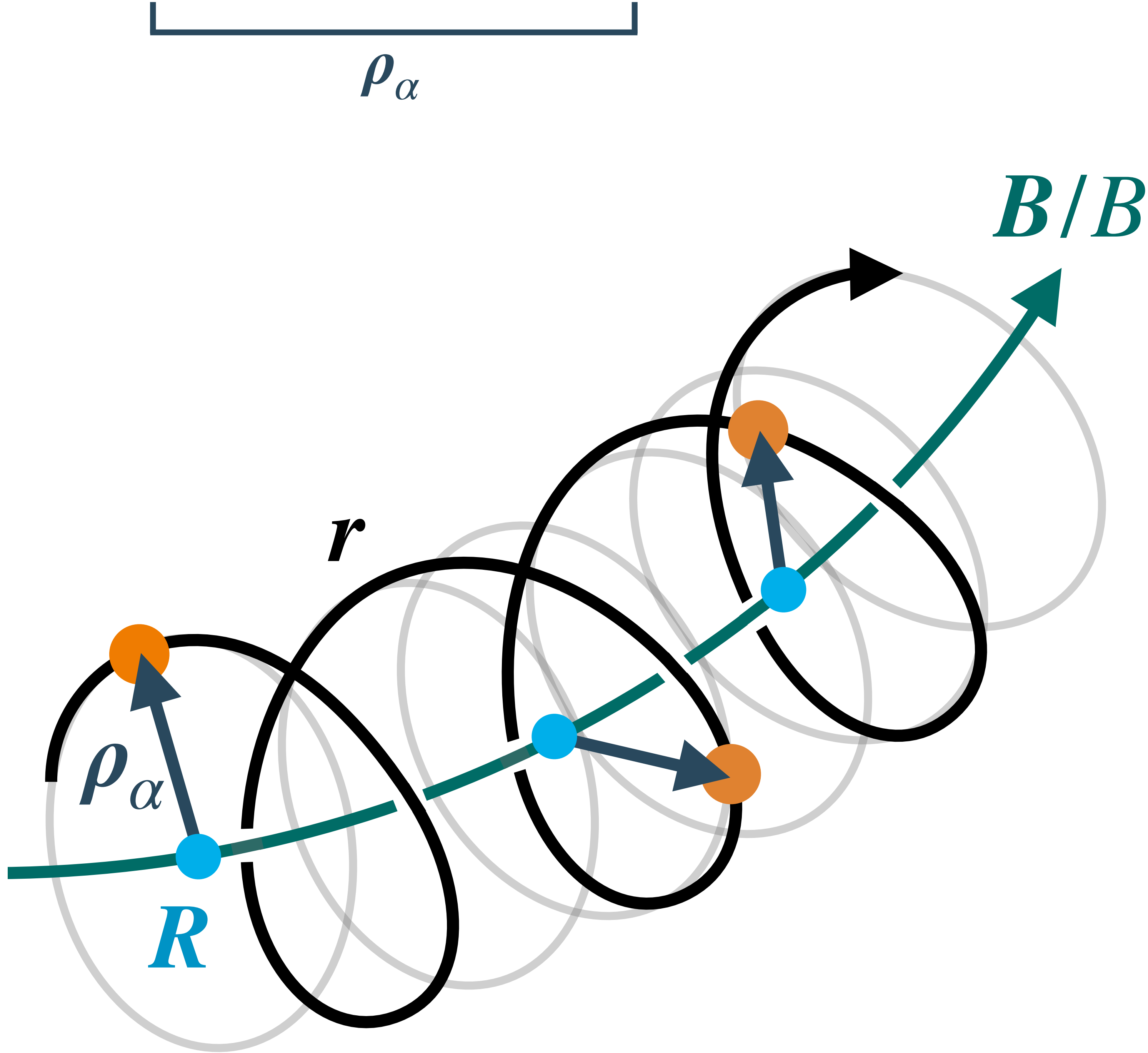
Finite Larmor radius (FLR) effects

short wavelength turbulence



Finite Larmor radius (FLR) effects

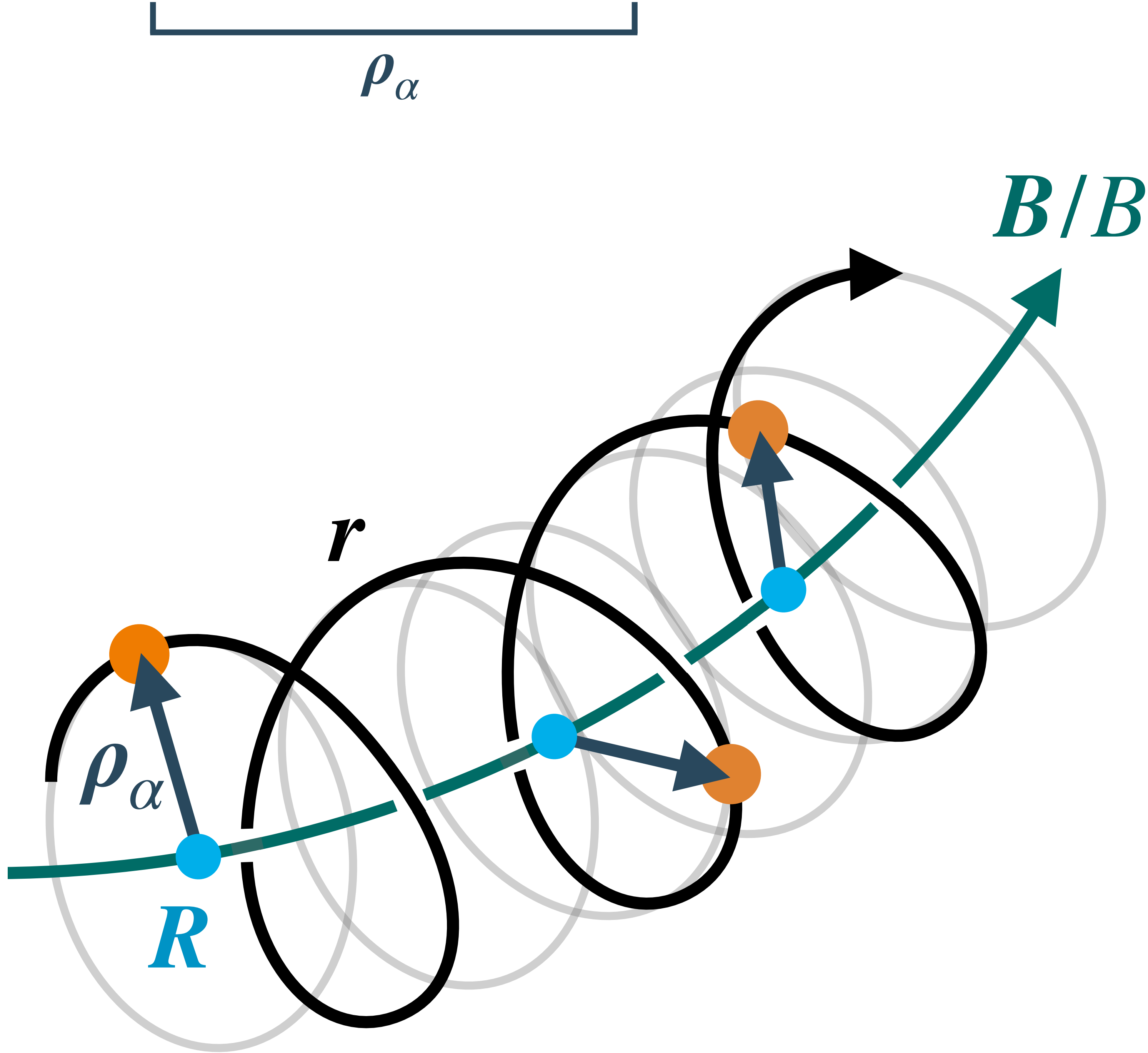
short wavelength turbulence



- ρ_α relevant in
- plasma core
 - core-edge transition
 - near steep gradients

Finite Larmor radius (FLR) effects

short wavelength turbulence



ρ_α relevant in

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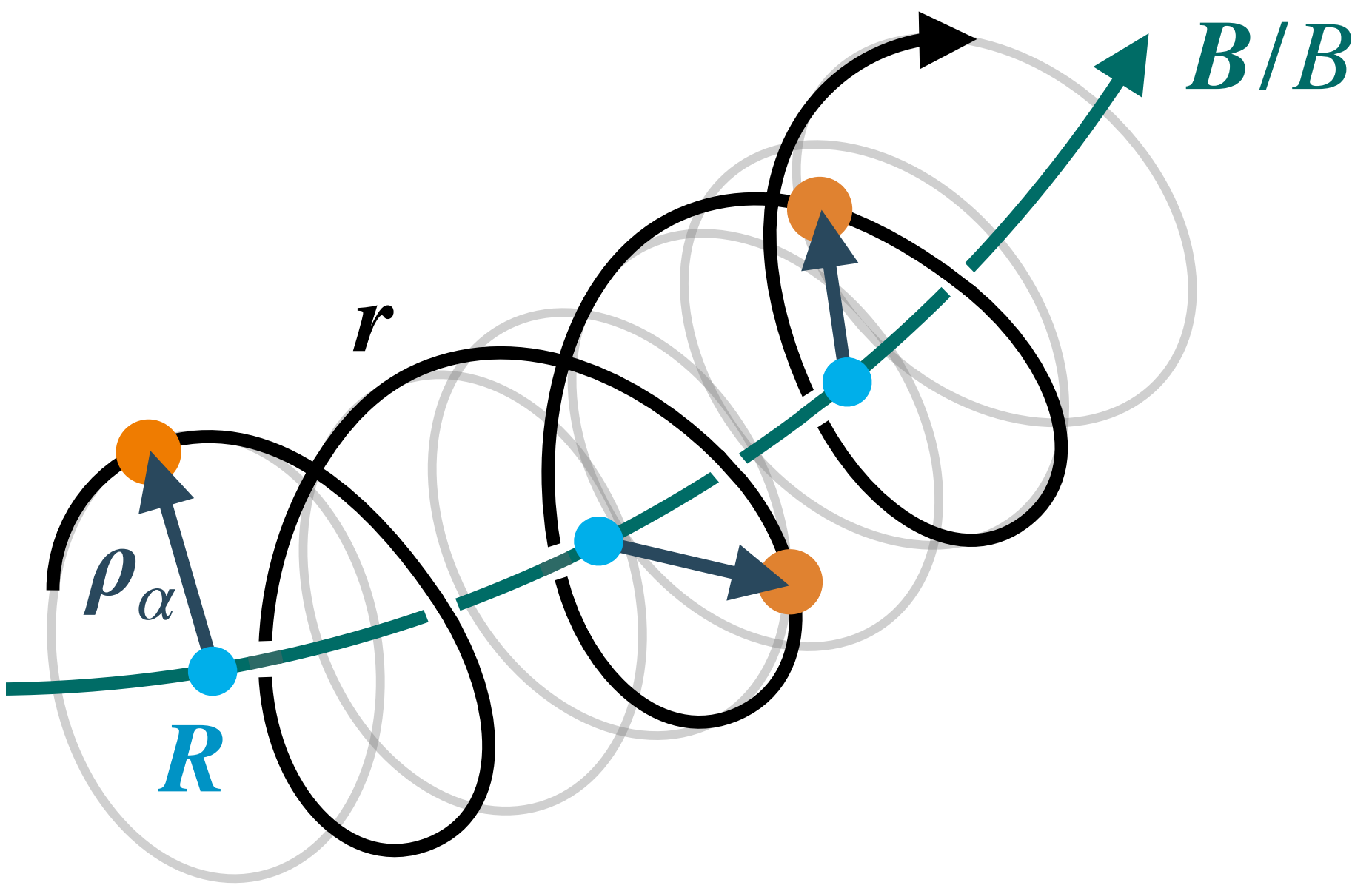
FLR effects

- to
 - simulations from the edge to the core
 - regions where gradients $\sim \rho_\alpha$ order
 - etc.
- through
 - gyro-averages

VLASOV EQ. $\frac{\partial f_\alpha}{\partial t} + \dot{\mathbf{R}} \cdot \nabla f_\alpha + \dot{v}_\parallel \frac{\partial f_\alpha}{\partial v_\parallel} = 0 \quad (\mathbf{B} = B \mathbf{b})$

$$\dot{\mathbf{R}} \sim \mathbf{B} v_\parallel + \frac{m_\alpha c}{q_\alpha} v_\parallel^2 \nabla \times \mathbf{b} + v_\parallel \nabla \times \mathbf{A}_{1\parallel} + \frac{c}{q_\alpha} \mathbf{b} \times \mu \nabla B + c \mathbf{b} \times \nabla \phi_1$$

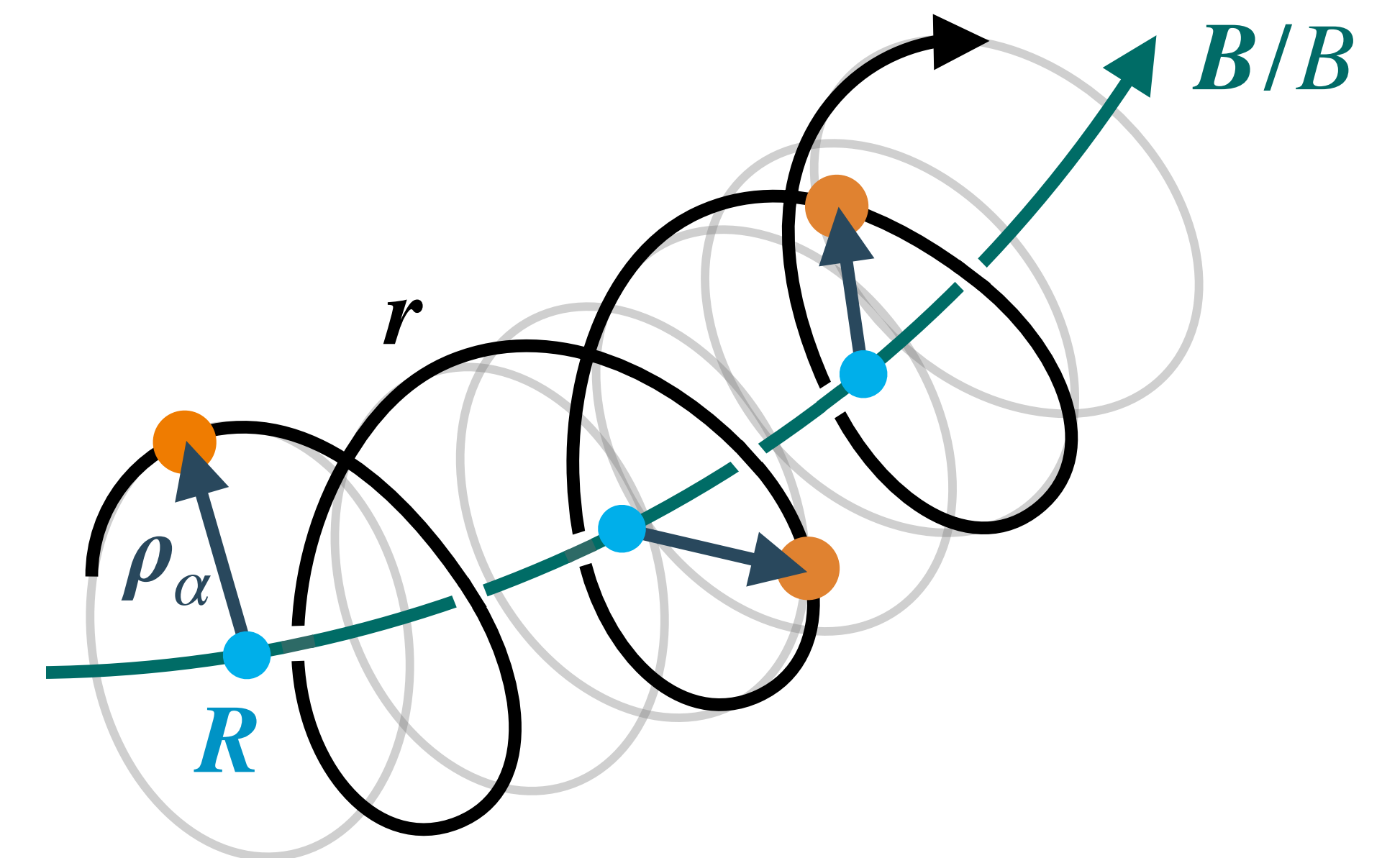
$$\dot{v}_\parallel \sim -Y \left(\mu \nabla B + q_\alpha \nabla \phi_1 \right) - \frac{q_\alpha}{m_\alpha c} \frac{\partial A_{1\parallel}}{\partial t}$$



VLASOV EQ. $\frac{\partial f_\alpha}{\partial t} + \dot{\mathbf{R}} \cdot \nabla f_\alpha + \dot{v}_\parallel \frac{\partial f_\alpha}{\partial v_\parallel} = 0 \quad (B = B \mathbf{b})$

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$\dot{\mathbf{R}} \sim$

$B v_\parallel$

$\frac{m_\alpha c}{q_\alpha} v_\parallel^2 \nabla \times \mathbf{b}$

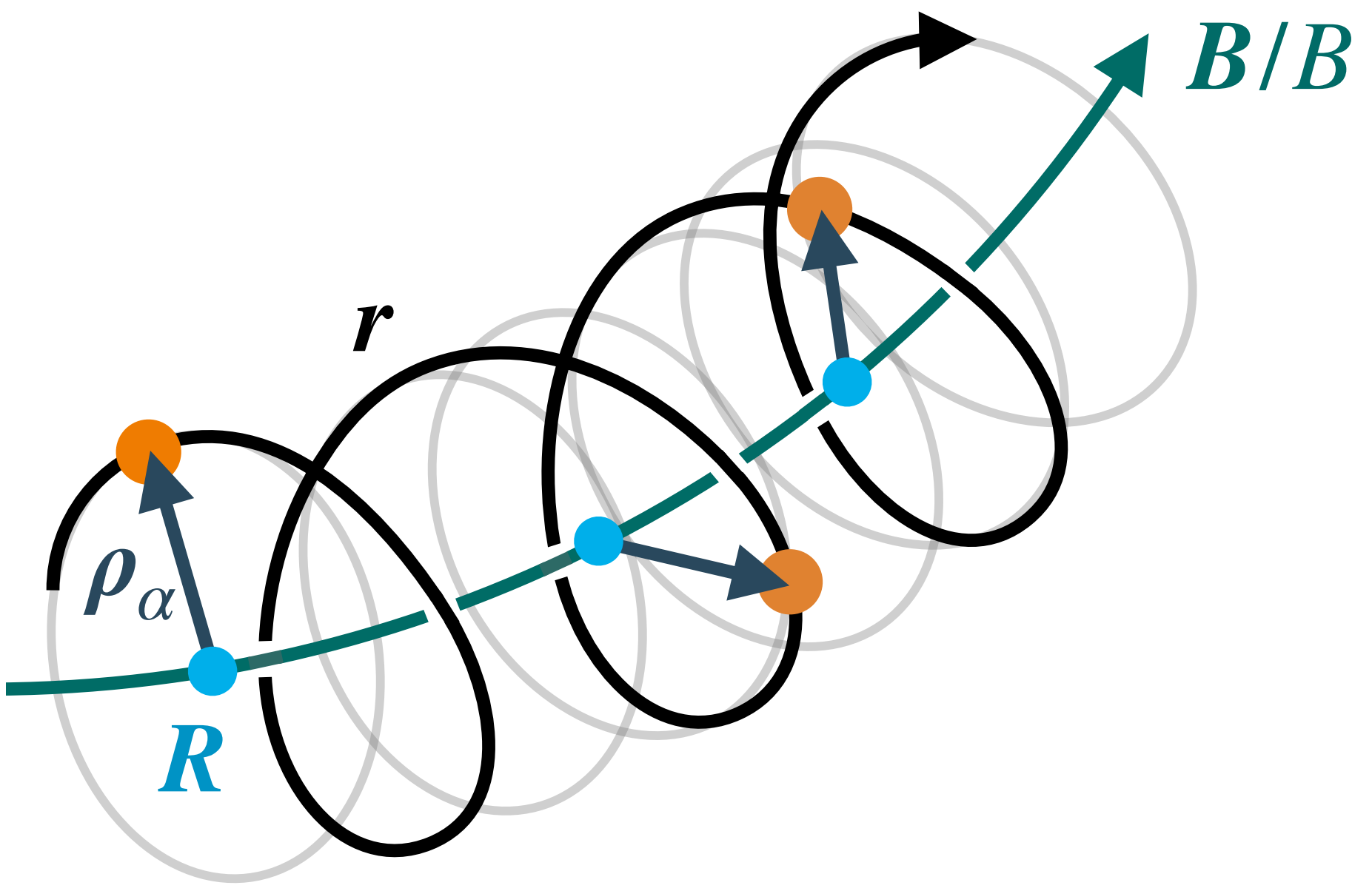
$+ v_\parallel \nabla \times A_{1\parallel}$

$+ \frac{c}{q_\alpha} \mathbf{b} \times \mu \nabla B + c \mathbf{b} \times \nabla \phi_1$

curvature drift

B advection

$\dot{v}_\parallel \sim - Y \left(\mu \nabla B + q_\alpha \nabla \phi_1 \right) - \frac{q_\alpha}{m_\alpha c} \frac{\partial A_{1\parallel}}{\partial t}$



VLASOV EQ. $\frac{\partial f_\alpha}{\partial t} + \dot{\mathbf{R}} \cdot \nabla f_\alpha + \dot{v}_\parallel \frac{\partial f_\alpha}{\partial v_\parallel} = 0 \quad (\mathbf{B} = B \mathbf{b})$

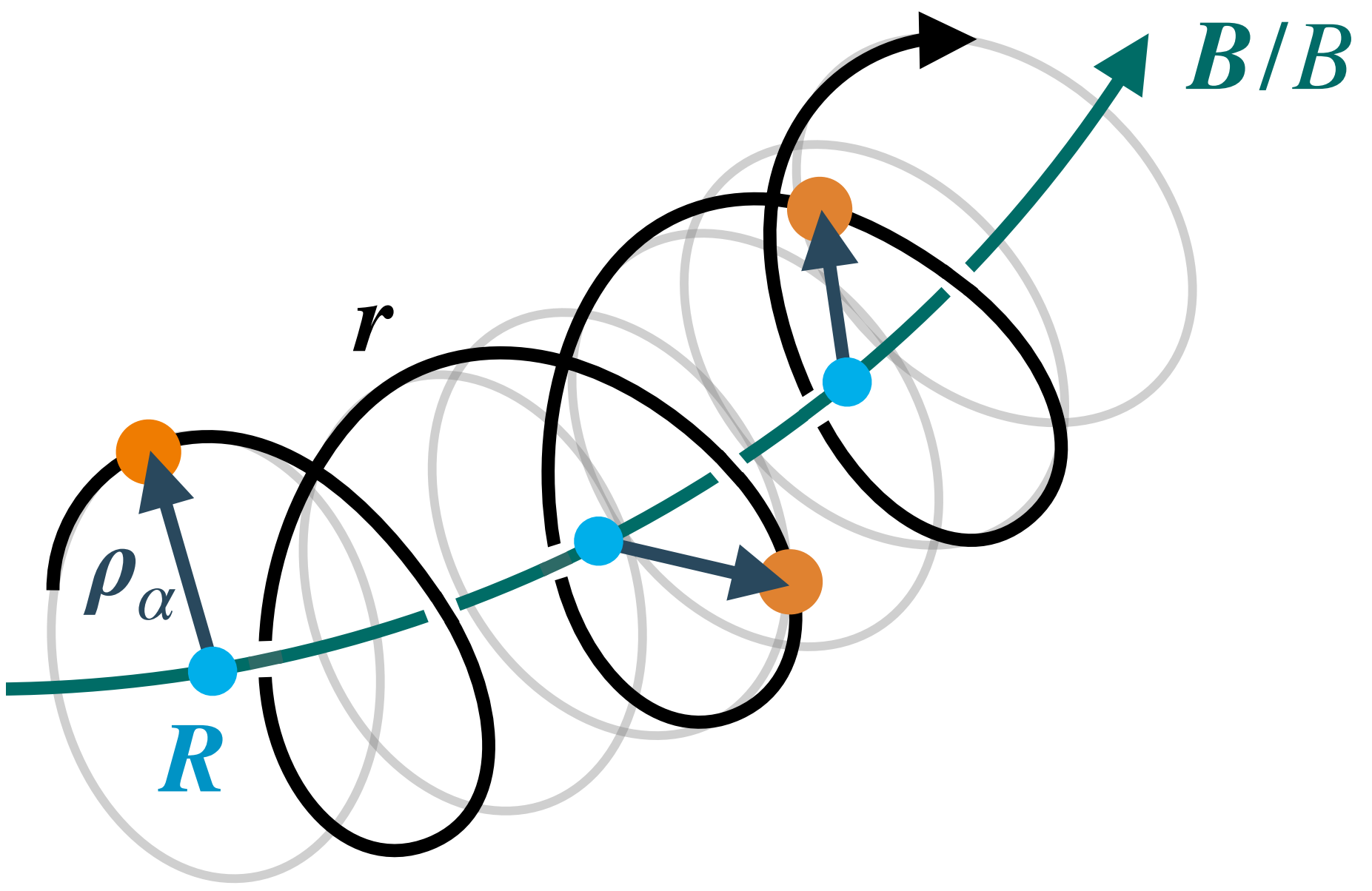
$\dot{\mathbf{R}} \sim \mathbf{B} v_\parallel + \frac{m_\alpha c}{q_\alpha} v_\parallel^2 \nabla \times \mathbf{b} + v_\parallel \nabla \times \mathbf{A}_{1\parallel} + \frac{c}{q_\alpha} \mathbf{b} \times \mu \nabla B + c \mathbf{b} \times \nabla \phi_1$

B advection

curvature drift

$\perp$ EM flutter

$\dot{v}_\parallel \sim -Y \left(\mu \nabla B + q_\alpha \nabla \phi_1 \right) - \frac{q_\alpha}{m_\alpha c} \frac{\partial A_{1\parallel}}{\partial t}$

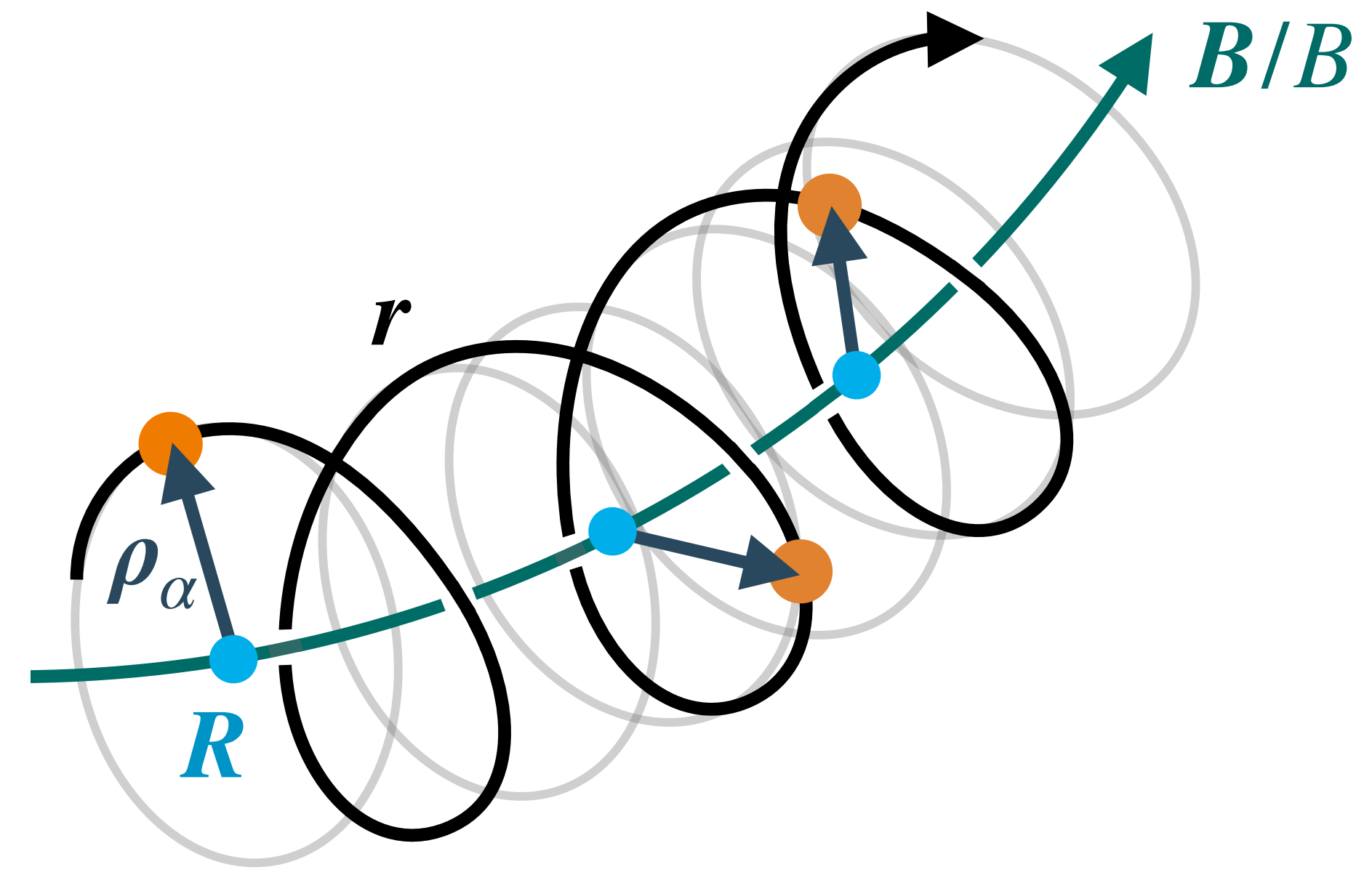


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B advectioncurvature drift $\perp$ EM flutter $\perp$ drift ∇B

$\dot{v}_\parallel \sim -Y \left(\mu \nabla B + q_\alpha \nabla \phi_1 \right) - \frac{q_\alpha}{m_\alpha c} \frac{\partial A_{1\parallel}}{\partial t}$



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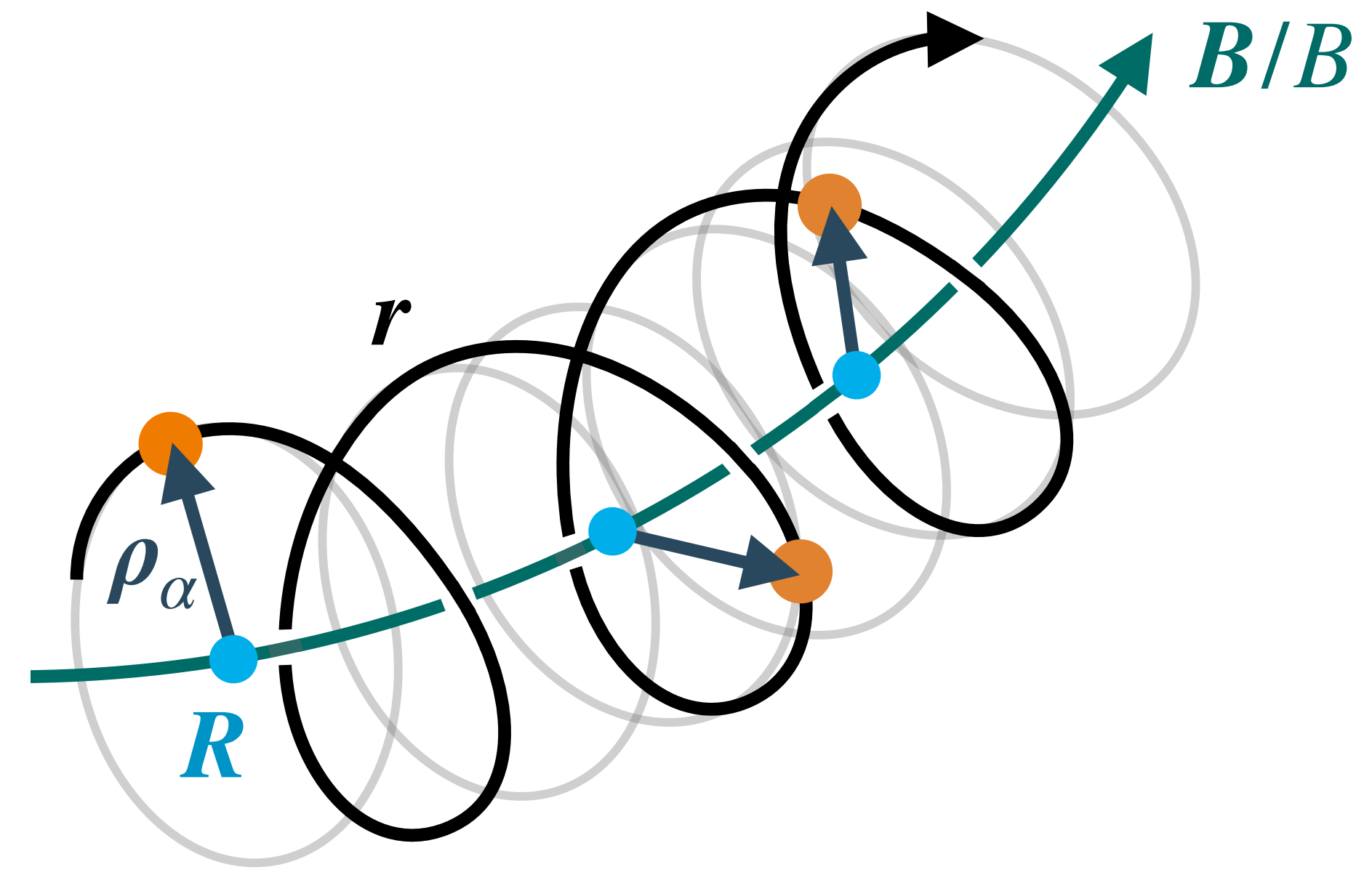
curvature drift

$\perp$ EM flutter

$\perp$ drift ∇B

$\perp$ drift $\mathbf{E} \times \mathbf{B}$

$\dot{v}_\parallel \sim -Y \left(\mu \nabla B + q_\alpha \nabla \phi_1 \right) - \frac{q_\alpha}{m_\alpha c} \frac{\partial A_{1\parallel}}{\partial t}$



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B advection

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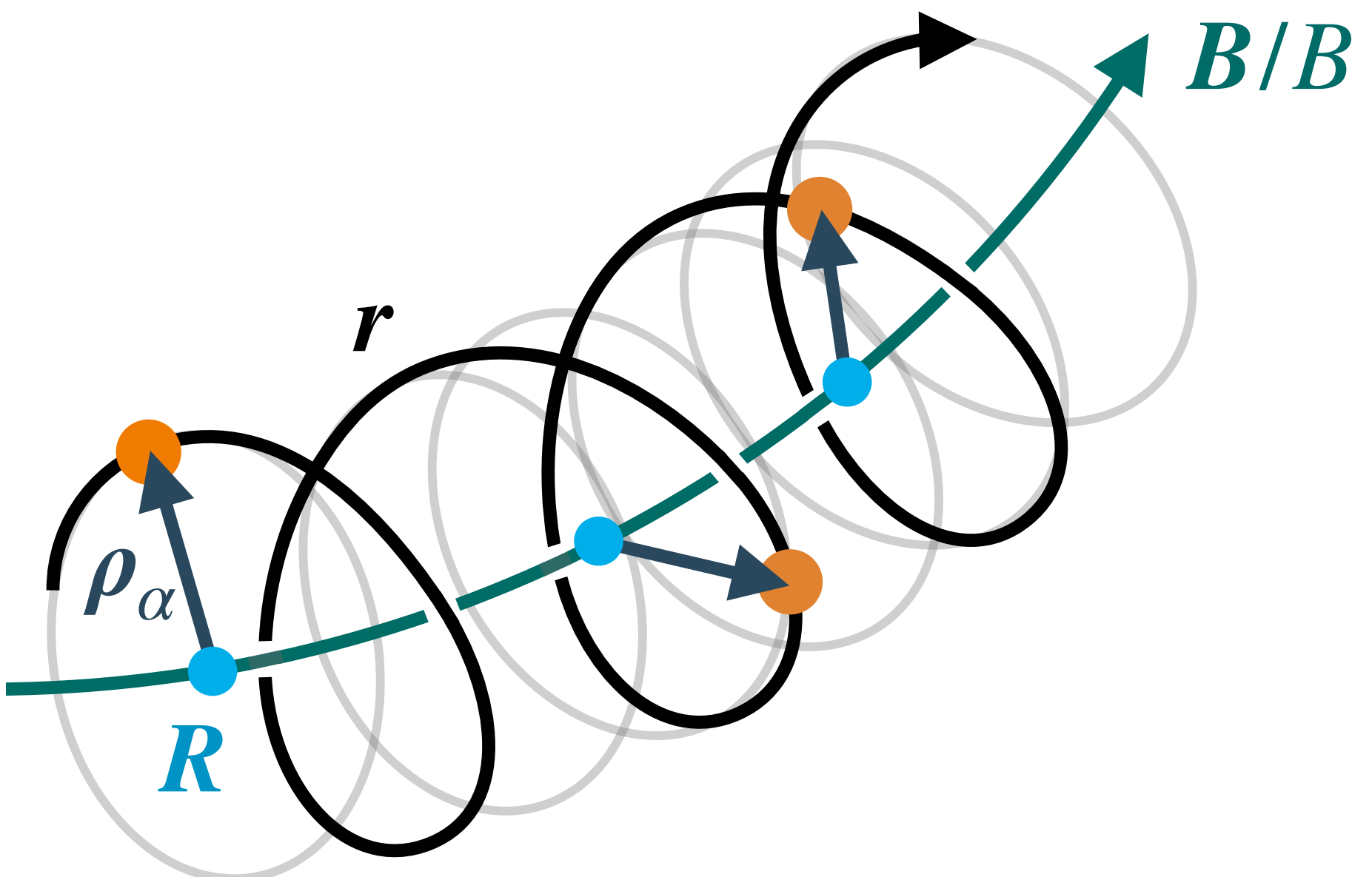
$\perp$ EM flutter

$\perp$ drift ∇B

$\perp$ drift $\mathbf{E} \times \mathbf{B}$

$\dot{v}_\parallel \sim -Y \left(\mu \nabla B + q_\alpha \nabla \phi_1 \right) - \frac{q_\alpha}{m_\alpha c} \frac{\partial A_{1\parallel}}{\partial t}$

mirror force



VLASOV EQ. $\frac{\partial f_\alpha}{\partial t} + \dot{\mathbf{R}} \cdot \nabla f_\alpha + \dot{v}_\parallel \frac{\partial f_\alpha}{\partial v_\parallel} = 0 \quad (\mathbf{B} = B \mathbf{b})$

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B advection

curvature drift

$\perp$ EM flutter

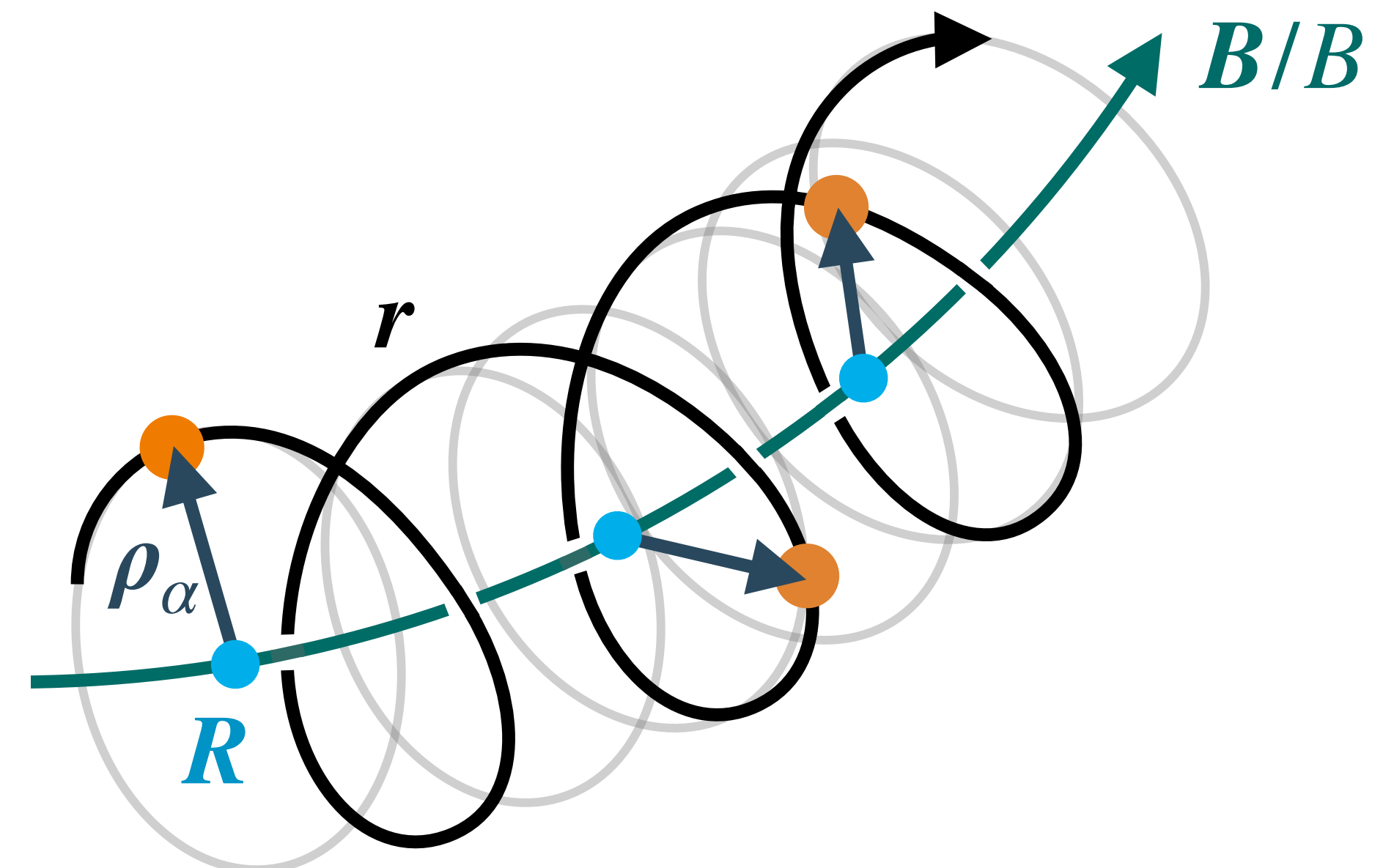
$\perp$ drift ∇B

$\perp$ drift $\mathbf{E} \times \mathbf{B}$

$\dot{v}_\parallel \sim -Y \left(\mu \nabla B + q_\alpha \nabla \phi_1 \right) - \frac{q_\alpha}{m_\alpha c} \frac{\partial A_{1\parallel}}{\partial t}$

$\mathbf{E}$ field acceler.

mirror force



VLASOV EQ. $\frac{\partial f_\alpha}{\partial t} + \dot{\mathbf{R}} \cdot \nabla f_\alpha + \dot{v}_\parallel \frac{\partial f_\alpha}{\partial v_\parallel} = 0 \quad (B = B \mathbf{b})$

$\dot{\mathbf{R}} \sim$

$B v_\parallel$

$+$

$\frac{m_\alpha c}{q_\alpha} v_\parallel^2 \nabla \times \mathbf{b}$

$+$

$v_\parallel \nabla \times A_{1\parallel}$

$+$

$\frac{c}{q_\alpha} \mathbf{b} \times \mu \nabla B$

$+$

$c \mathbf{b} \times \nabla \phi_1$

B advection

curvature drift

$\perp$ EM flutter

$\perp$ drift ∇B

$\perp$ drift $E \times B$

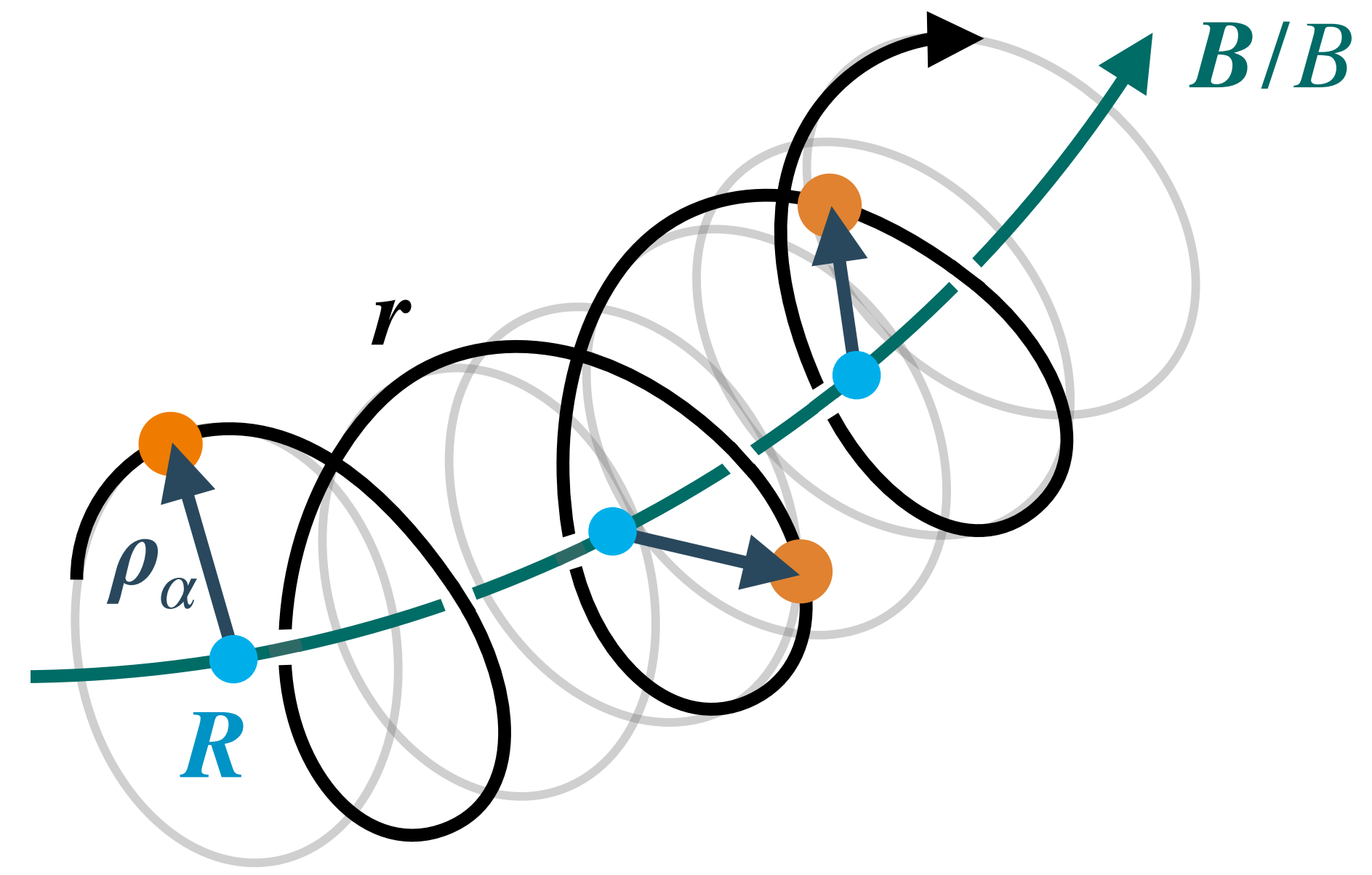
$\dot{v}_\parallel \sim$

$-Y \left(\mu \nabla B + q_\alpha \nabla \phi_1 \right) - \frac{q_\alpha}{m_\alpha c} \frac{\partial A_{1\parallel}}{\partial t}$

E field acceler.

mirror force

induction



VLASOV EQ.

$$\frac{\partial f_\alpha}{\partial t} + \dot{\mathbf{R}} \cdot \nabla f_\alpha + \dot{v}_\parallel \frac{\partial f_\alpha}{\partial v_\parallel} = 0 \quad (\mathbf{B} = B \mathbf{b})$$

gyro-average

$$\langle \phi(\mathbf{r}) \rangle_R = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_\alpha \cdot \nabla} \phi(\mathbf{R})$$

$\dot{\mathbf{R}} \sim$

$B v_\parallel$

$\frac{m_\alpha c}{q_\alpha} v_\parallel^2 \nabla \times \mathbf{b}$

$v_\parallel \nabla \times \mathbf{A}_{1\parallel}$

$\frac{c}{q_\alpha} \mathbf{b} \times \mu \nabla B$

$c \mathbf{b} \times \nabla \phi_1$

B advection

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$\perp$ EM flutter

$\perp$ drift ∇B

$\perp$ drift $\mathbf{E} \times \mathbf{B}$

$\dot{v}_\parallel \sim$

$\mu \nabla B$

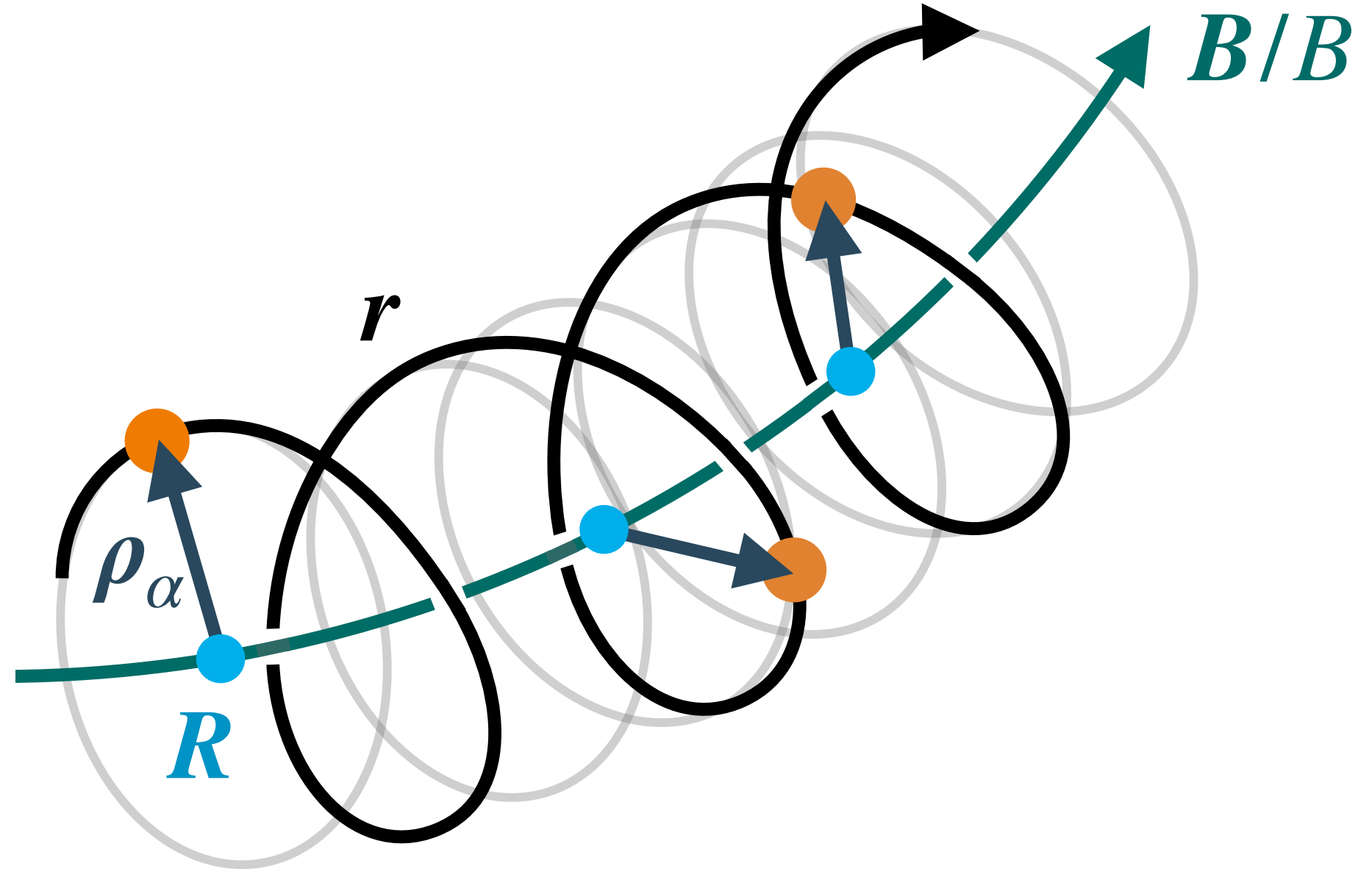
$q_\alpha \nabla \phi_1$

$-\frac{q_\alpha}{m_\alpha c} \frac{\partial A_{1\parallel}}{\partial t}$

$\mathbf{E}$ field acceler.

mirror force

induction



Gyro-averages and GENE-X equations



VLASOV EQ.

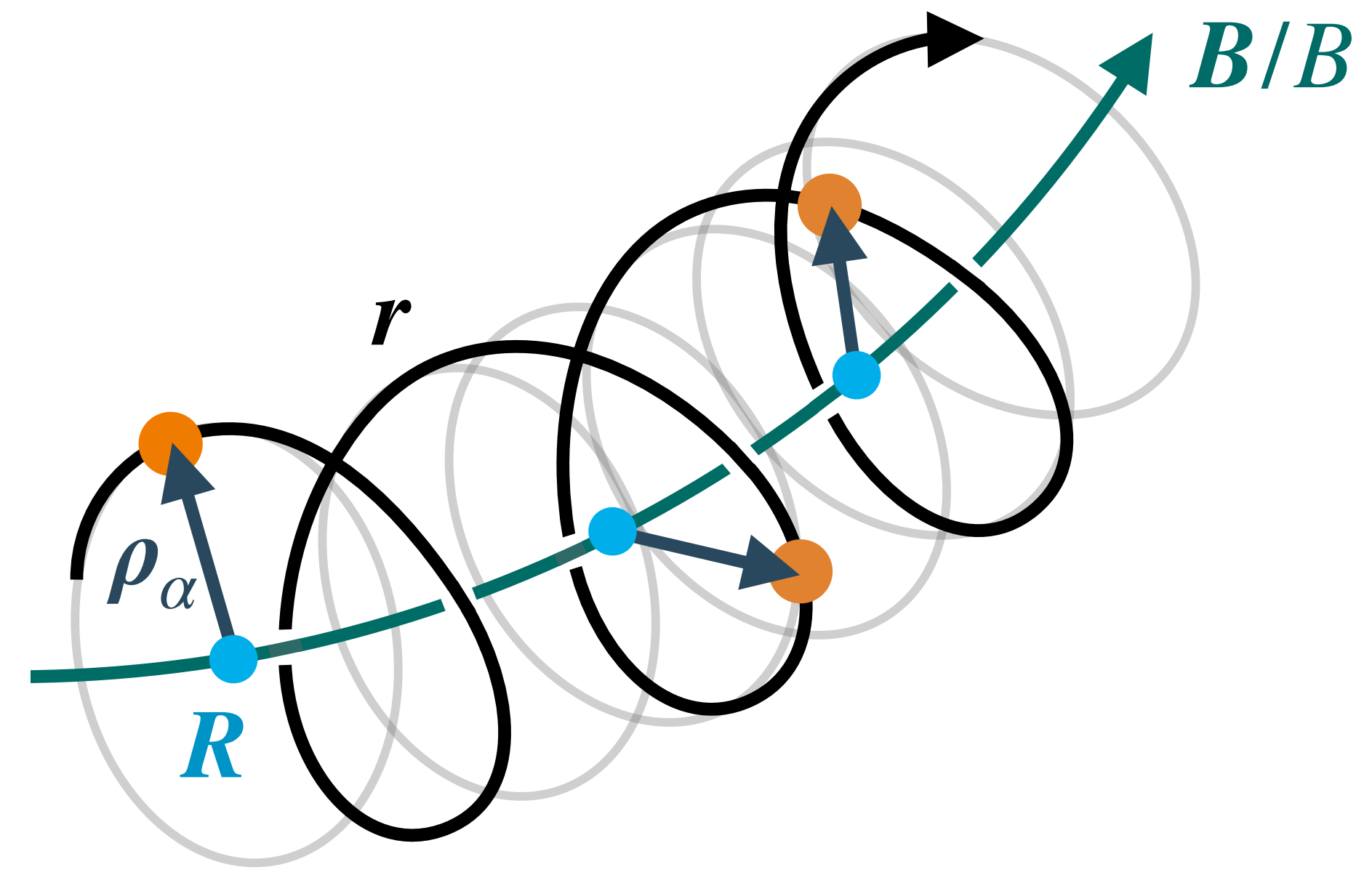
$$\frac{\partial f_\alpha}{\partial t} + \dot{\mathbf{R}} \cdot \nabla f_\alpha + \dot{v}_\parallel \frac{\partial f_\alpha}{\partial v_\parallel} = 0 \quad (\mathbf{B} = B \mathbf{b})$$

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$$\dot{\mathbf{R}} \sim \underbrace{B v_\parallel}_{B \text{ advection}} + \underbrace{\frac{m_\alpha c}{q_\alpha} v_\parallel^2 \nabla \times \mathbf{b}}_{\text{curvature drift}} + \underbrace{v_\parallel \langle \nabla \times \mathbf{A}_{1\parallel} \rangle_R}_{\perp \text{ EM flutter}} + \underbrace{\frac{c}{q_\alpha} \mathbf{b} \times \mu \nabla B}_{\perp \text{ drift } \nabla B} + \underbrace{c \mathbf{b} \times \nabla \langle \phi_1 \rangle_R}_{\perp \text{ drift } \mathbf{E} \times \mathbf{B}}$$

$$\dot{v}_\parallel \sim -Y \left(\underbrace{\mu \nabla B}_{\text{mirror force}} + \underbrace{q_\alpha \nabla \langle \phi_1 \rangle_R}_{\mathbf{E} \text{ field acceler.}} \right) - \underbrace{\frac{q_\alpha}{m_\alpha c} \frac{\partial \langle \mathbf{A}_{1\parallel} \rangle_R}{\partial t}}_{\text{induction}}$$



Padé approximant

gyro-average

$$\langle \phi(\mathbf{r}) \rangle_{\mathbf{R}} = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_{\alpha} \cdot \nabla} \phi(\mathbf{R})$$

Padé approximant

gyro-average

$$\langle \phi(\mathbf{r}) \rangle_{\mathbf{R}} = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_\alpha \cdot \nabla} \phi(\mathbf{R})$$

↓ $\mathcal{F}$ local gyrokinetics

$$\langle \phi(\mathbf{r}) \rangle_{\mathbf{R}} = \frac{1}{(2\pi)^3} \int d\mathbf{k} J_0(\rho_\alpha k_\perp) e^{i\mathbf{k} \cdot \mathbf{R}} \hat{\phi}(\mathbf{k})$$

Padé approximant

gyro-average

$$\langle \phi(\mathbf{r}) \rangle_R = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_\alpha \cdot \nabla} \phi(\mathbf{R})$$

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↓
FLR operator in $\mathcal{F}$
Bessel func. of the 1st kind

Padé approximant

gyro-average

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↓
FLR operator in $\mathcal{F}$
Bessel func. of the 1st kind

Taylor on $J_0 \xrightarrow{k_\perp \rightarrow \infty} \infty$

⇓

another approx., e.g., **Padé approximant**

$$J_0(\rho_\alpha k_\perp) \stackrel{P0/2}{\approx} \frac{1}{1 + \rho_\alpha^2 k_\perp^2 / 4}$$

gyro-average

$$\langle \phi(\mathbf{r}) \rangle_R = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_\alpha \cdot \nabla} \phi(\mathbf{R})$$

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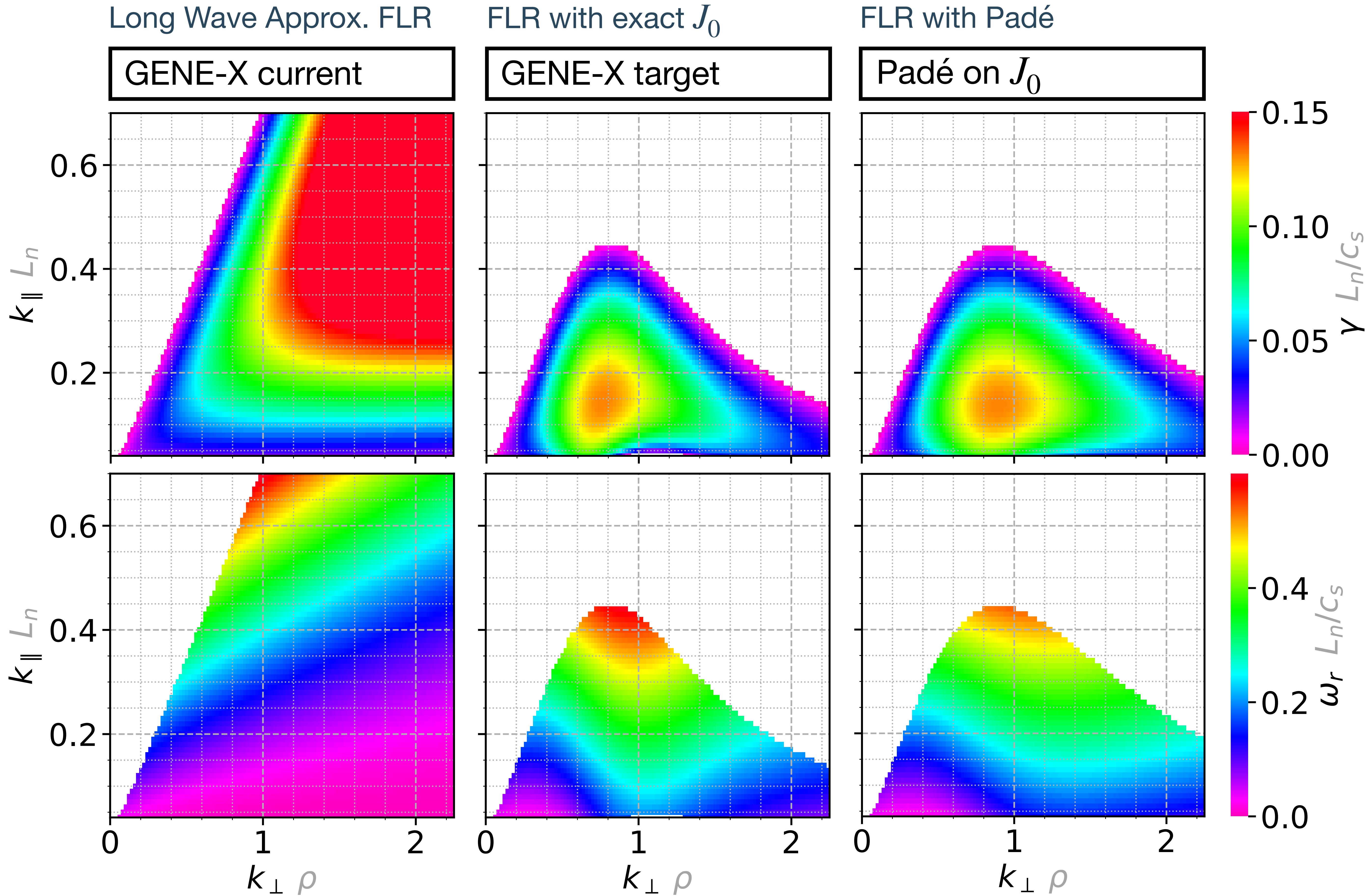
↓ $\mathcal{F}^{-1}$ global gyrokinetics

$$\left(1 - \frac{1}{4} \rho_\alpha^2 \nabla_\perp^2 \right) \langle \phi(\mathbf{r}) \rangle_R = \phi(\mathbf{R})$$

**differential equation
for a gyro-average**

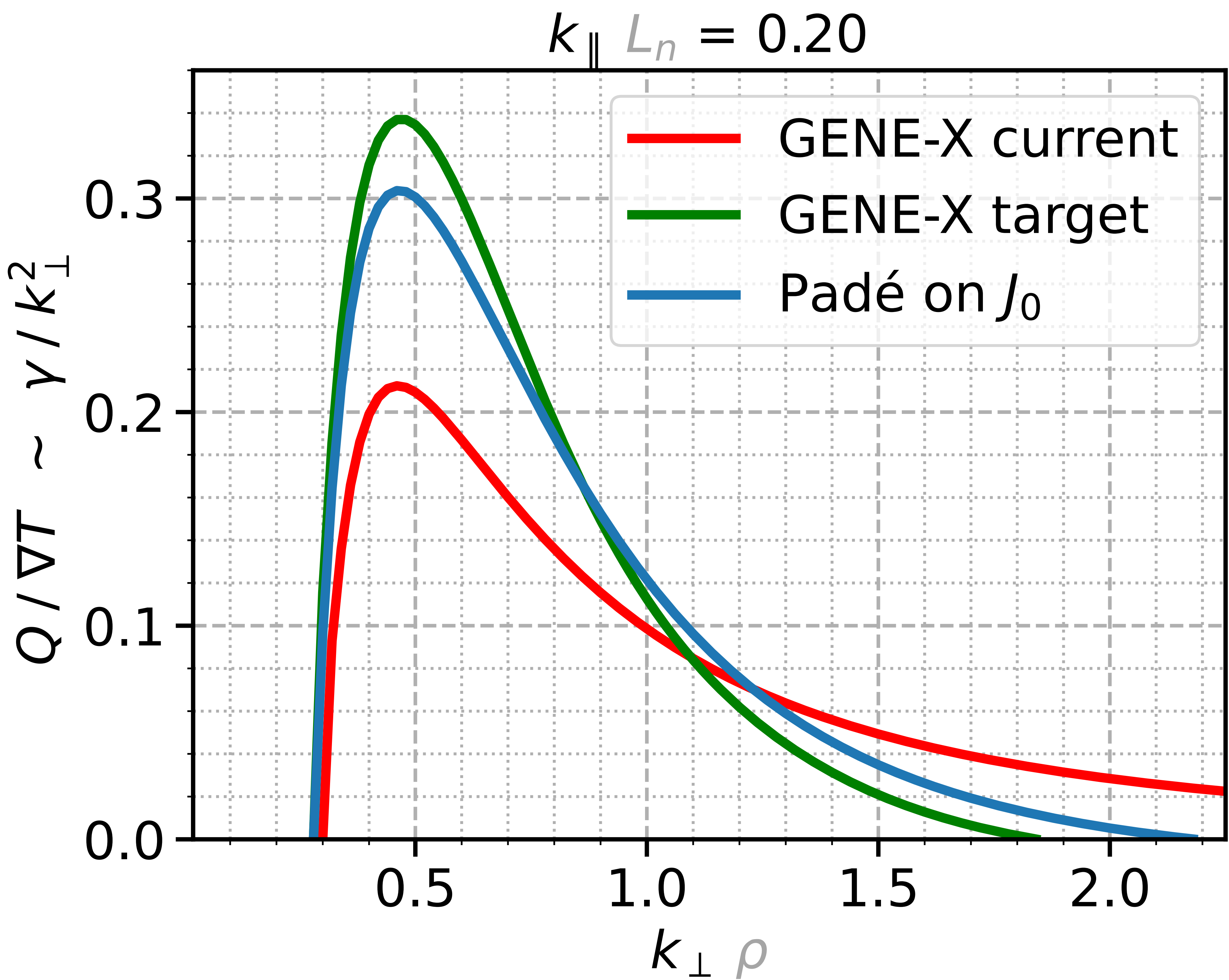
Impact of FLR models

EXAMPLE in the electrostatic slab ion temperature gradient (ITG) dispersion relation in Fourier space with diffusion in $\parallel$ and $\perp$ dirs.



Impact of FLR models – ITG disp. rel.

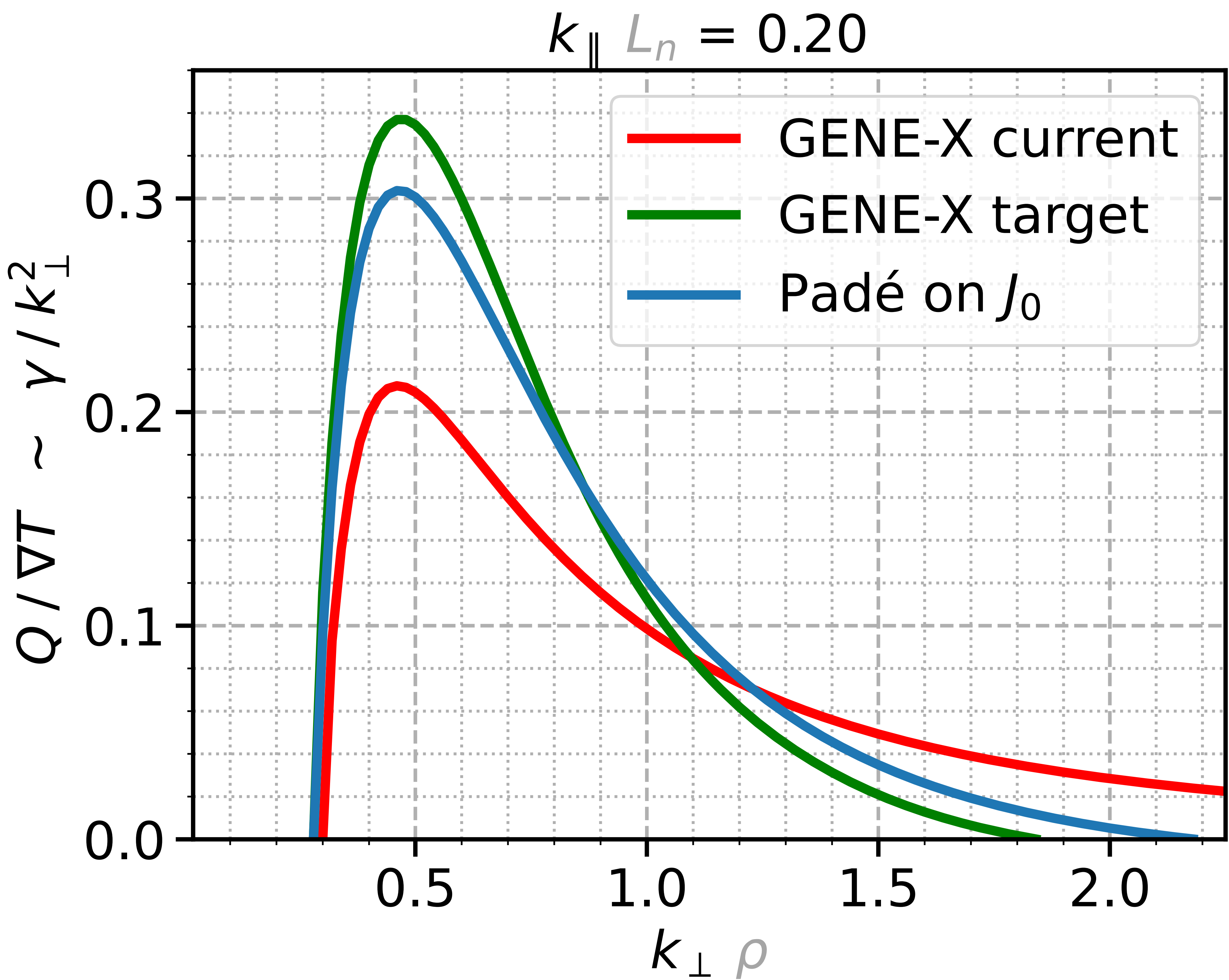
GENE-X current → Long Wave App. FLR
GENE-X target → FLR with exact J_0
Padé on J_0 → FLR with Padé



analysis of quasi-linear flux

Impact of FLR models – ITG disp. rel.

GENE-X current → Long Wave App. FLR
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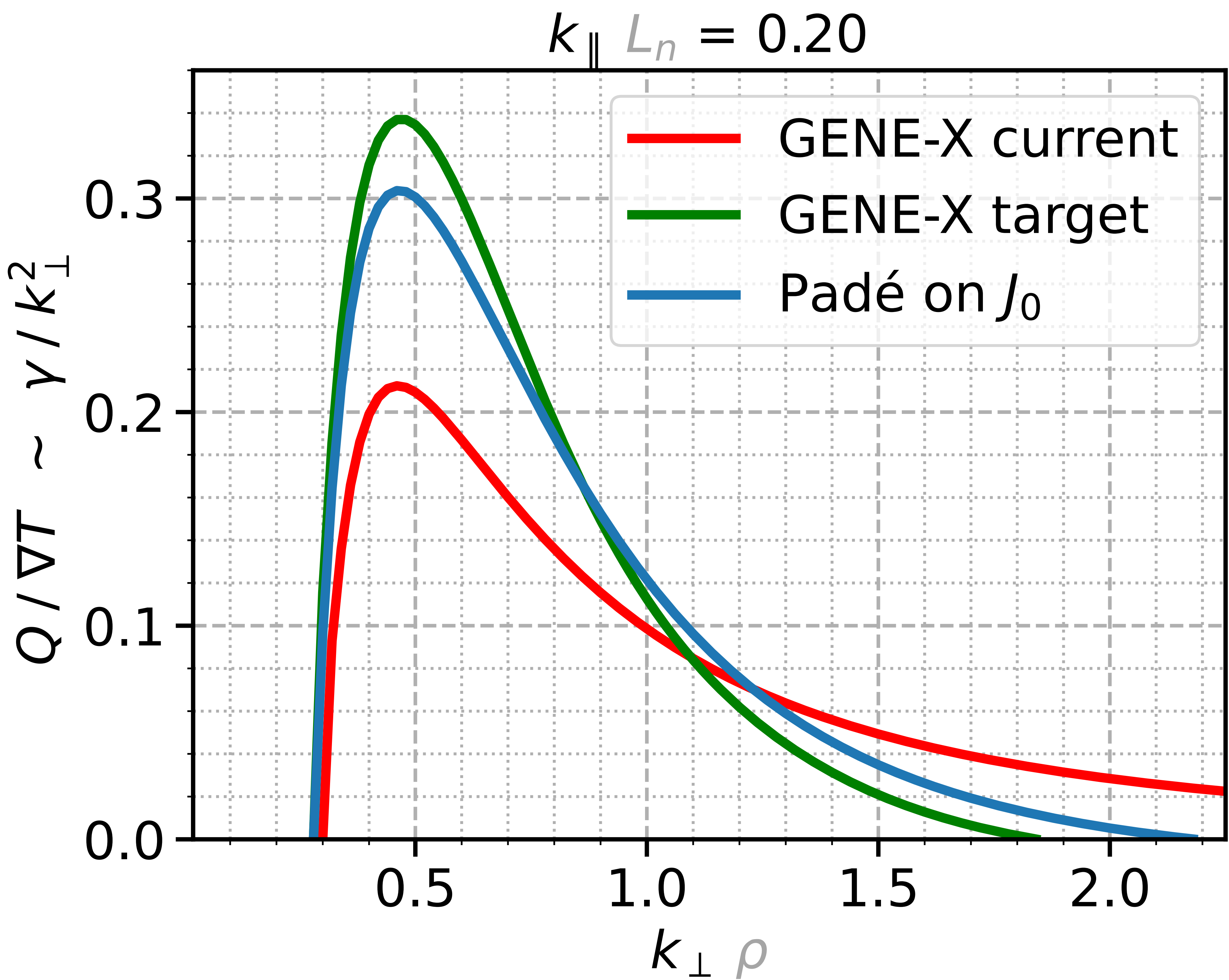


analysis of quasi-linear flux

└ major Physics change from GENE-X **current** to **target**

Impact of FLR models – ITG disp. rel.

GENE-X current → Long Wave App. FLR
GENE-X target → FLR with exact J_0
Padé on J_0 → FLR with Padé



analysis of quasi-linear flux

- major Physics change from GENE-X **current** to **target**
- Padé on J_0** not far-fetched from **target** and provides very convenient algorithm

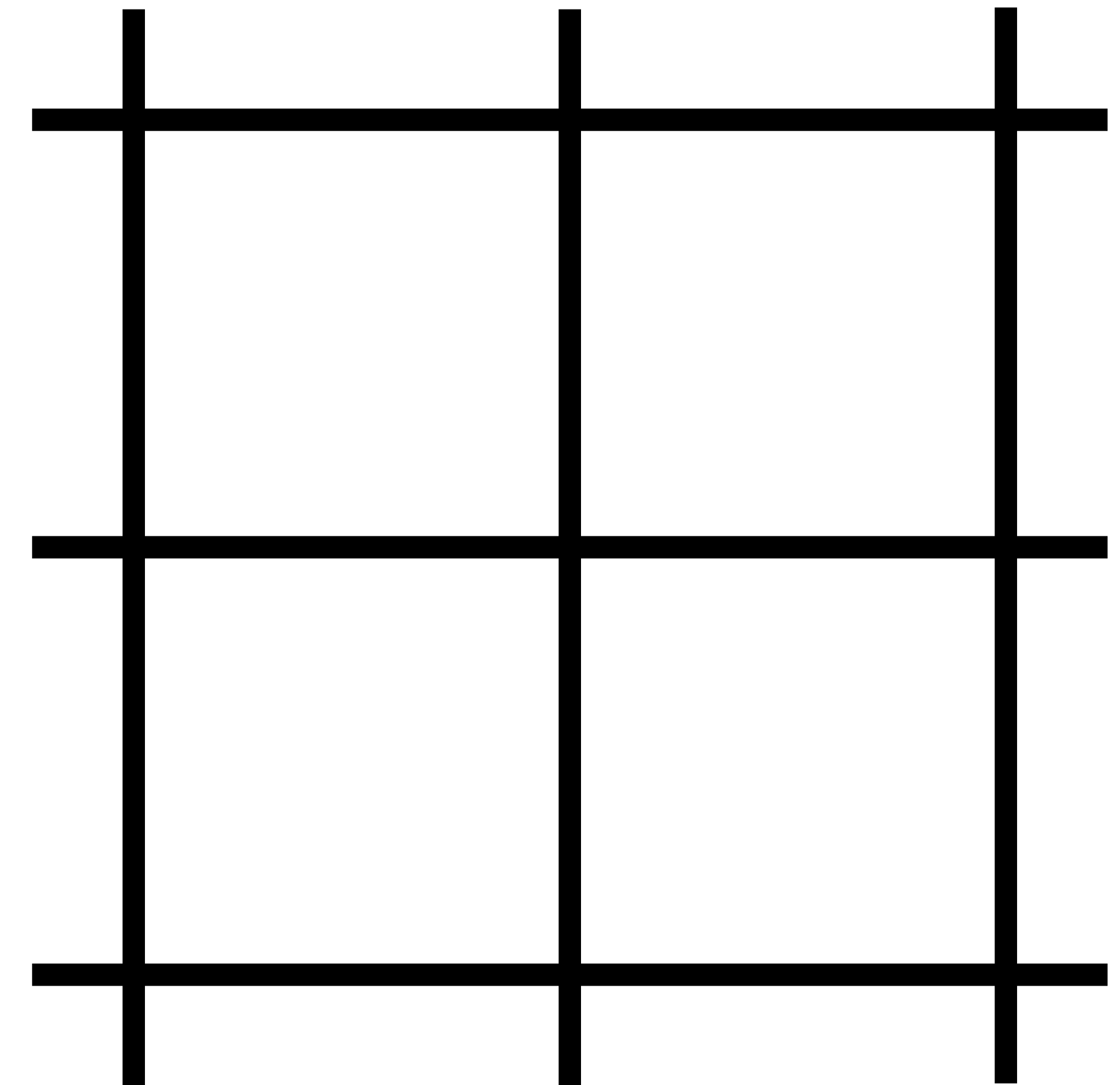
$$\left(1 - \frac{1}{4} \rho_{\alpha}^2 \nabla_{\perp}^2\right) \langle \phi(\mathbf{r}) \rangle_R = \phi(\mathbf{R})$$

differential eq. for gyro-average

Gyro-matrix approach

gyro-average

$$\langle \phi(\mathbf{r}) \rangle_{\mathbf{R}} = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_{\alpha} \cdot \nabla} \phi(\mathbf{R})$$

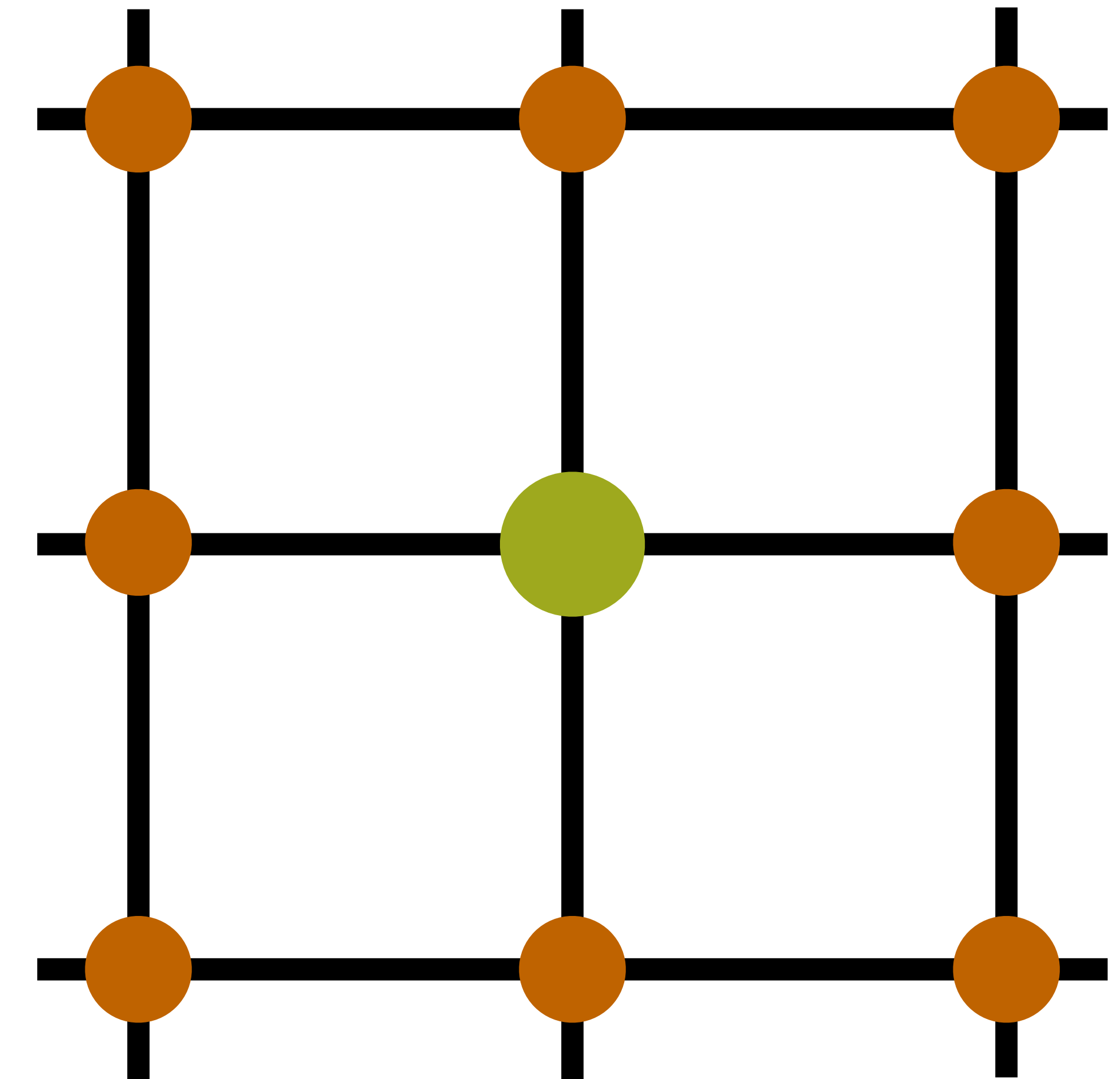


Gyro-matrix approach

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$$\langle \phi(\mathbf{r}) \rangle_{\mathbf{R}} = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_{\alpha} \cdot \nabla} \phi(\mathbf{R})$$

→ grid of **values**



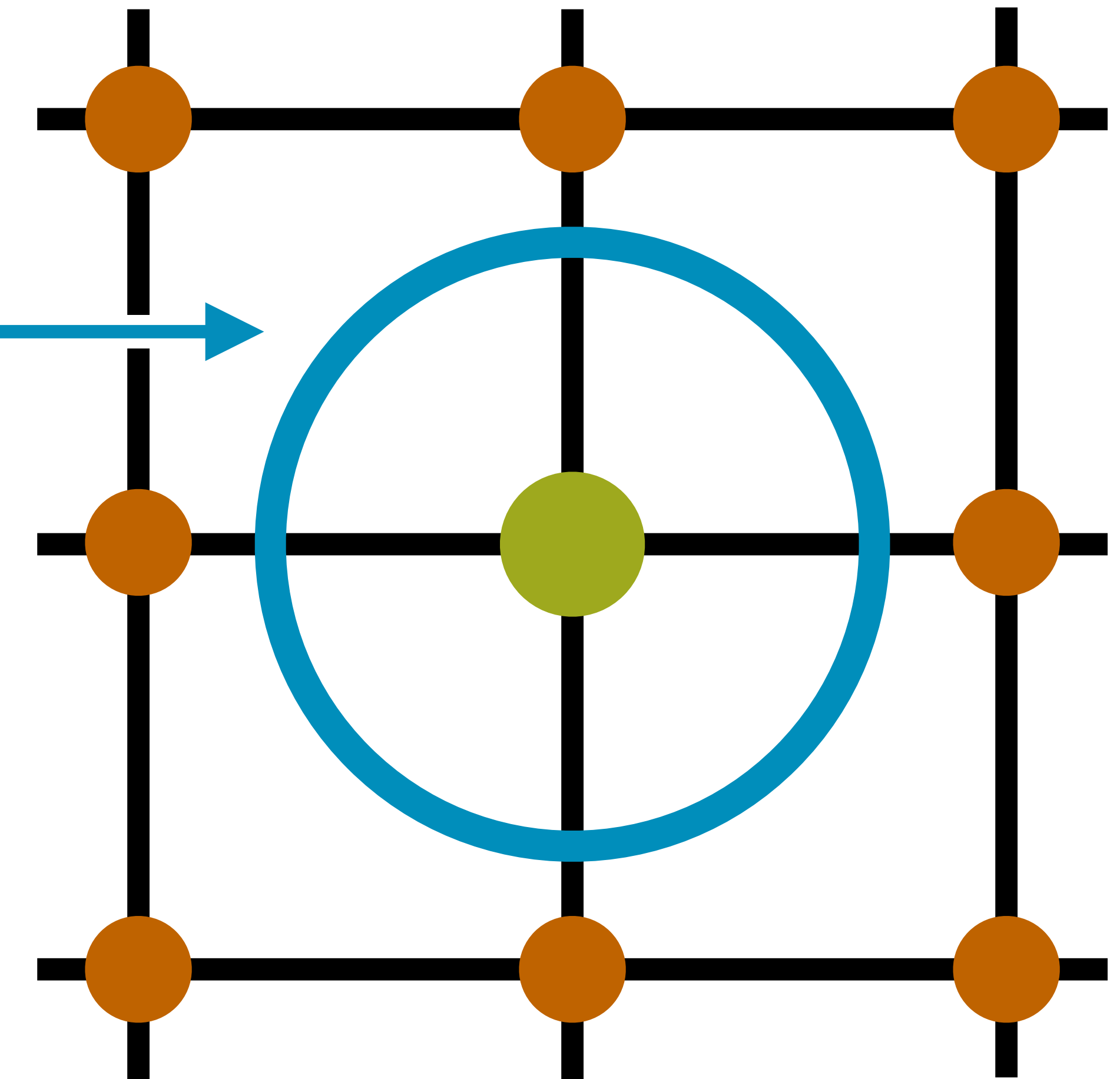
Gyro-matrix approach

gyro-average

$$\langle \phi(\mathbf{r}) \rangle_R = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_\alpha \cdot \nabla} \phi(\mathbf{R})$$

→ grid of **values**

→ **interpolated quantities** around **each point**



Gyro-matrix approach

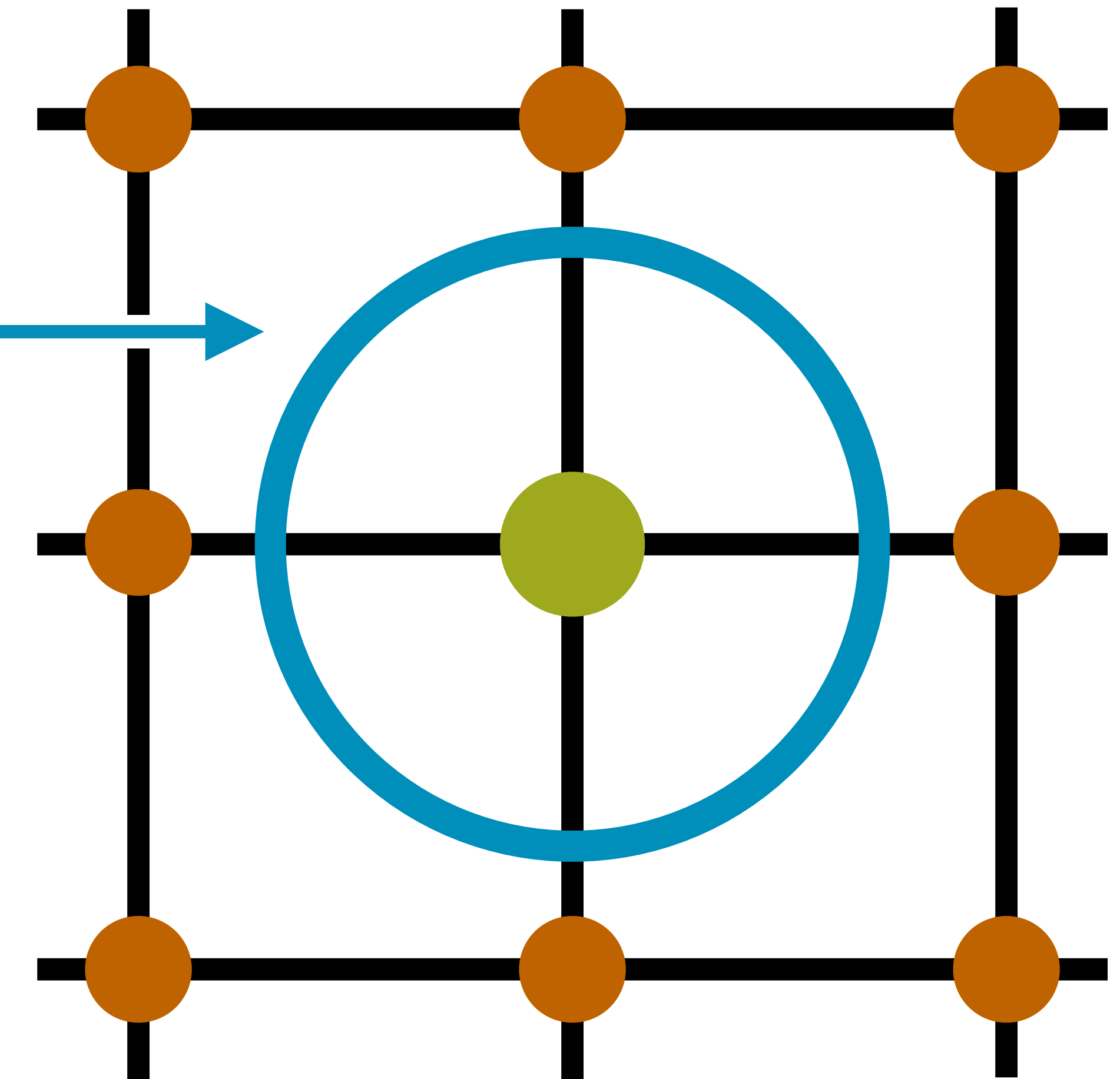
gyro-average

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- grid of **values**
- **interpolated quantities** around **each point**
- matrix combining interpolation and gyro-average

→ **gyro-matrix** $\langle \phi \rangle_R = G(\mu) \phi$

calculates gyro-averages for all orders
up to the discretization error



Gyro-matrix approach

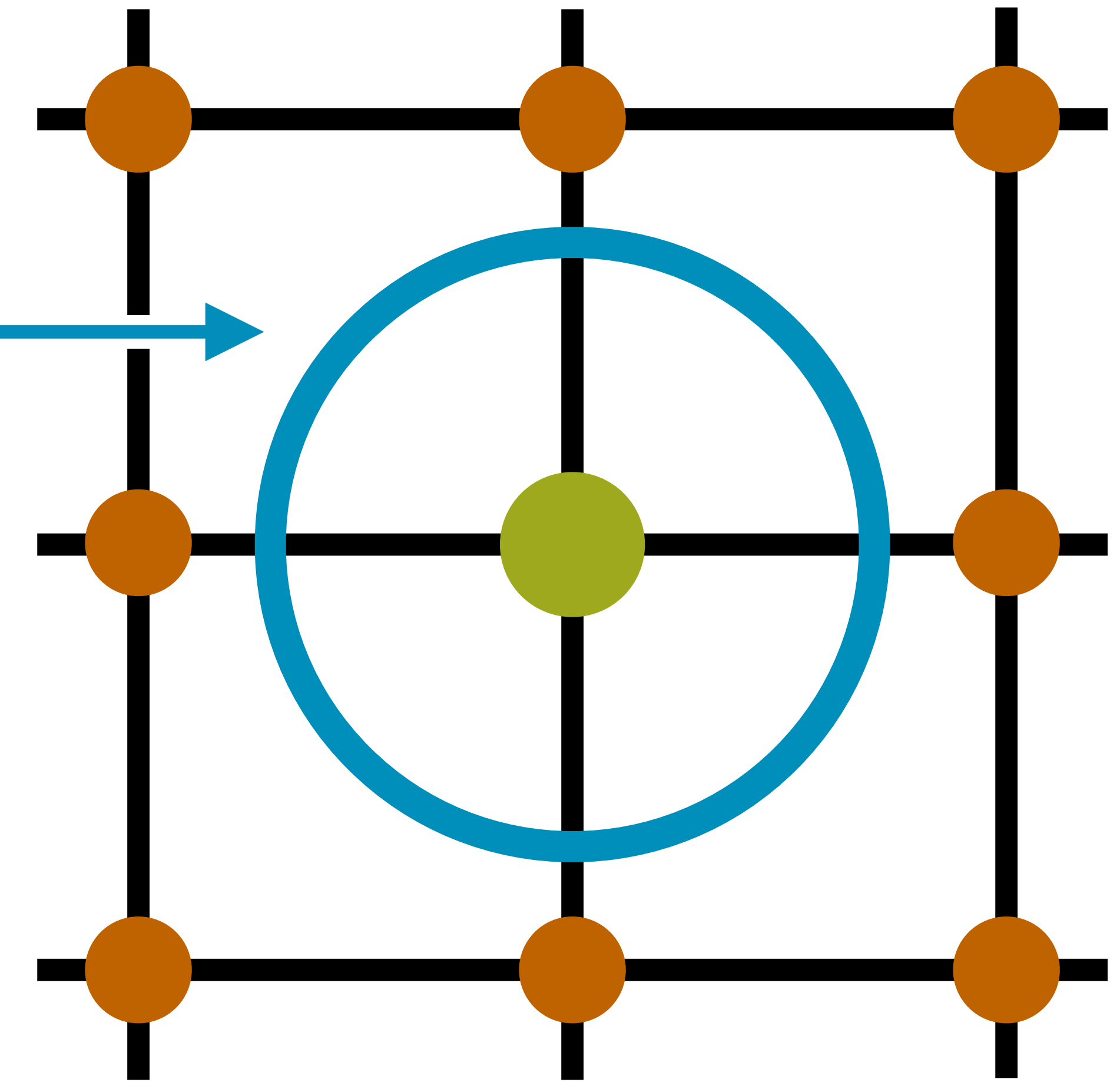
gyro-average

$$\langle \phi(\mathbf{r}) \rangle_R = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_\alpha \cdot \nabla} \phi(\mathbf{R})$$

- grid of **values**
- **interpolated quantities** around **each point**
- matrix combining interpolation and gyro-average

gyro-matrix $\langle \phi \rangle_R = G(\mu) \phi$

calculates gyro-averages for all orders
up to the discretization error



CAVEATS: intricate algorithm, expensive in resources...

Study of finite Larmor radius (FLR) physics in tokamaks and stellarators with the gyrokinetic turbulence code GENE-X



SUMMARY

- Re-derivation of GENE-X eqs. accounting for FLR effects
 - Padé approx. of J_0
- Example of FLR models in the ES slab ITG disp. relation
 - Padé on J_0 is Physically close
and provides a convenient algorithm for FLR effects
- Gyro-matrix approach

THANK YOU



ADDENDUM

gyro-average

$$\langle \phi(\mathbf{r}) \rangle_R = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_\alpha \cdot \nabla} \phi(\mathbf{R})$$

$$\mathcal{L} = \frac{q_\alpha}{c} \mathbf{A} \cdot \dot{\mathbf{R}} + \frac{q_\alpha}{c} \langle A_{1\parallel} \rangle_R \mathbf{b} \cdot \dot{\mathbf{R}} + m_\alpha v_\parallel \mathbf{b} \cdot \dot{\mathbf{R}} + \frac{\mu B}{\Omega_\alpha} \dot{\theta} - q_\alpha \langle \phi_1 \rangle_R - \frac{1}{2} m_\alpha v_\parallel^2 - \mu B + \frac{m_\alpha c^2}{2 B^2} |\nabla_\perp \phi_1|^2$$

VLASOV EQ. $\frac{\partial f_\alpha}{\partial t} + \dot{\mathbf{R}} \cdot \nabla f_\alpha + \dot{v}_\parallel \frac{\partial f_\alpha}{\partial v_\parallel} = 0$ $\left(B^* \sim B + \frac{m_\alpha c}{q_\alpha} v_\parallel \nabla \times \mathbf{b} + \langle \nabla \times \mathbf{A}_1 \rangle \right)$

using the GENE-X Poisson bracket [1] for our coordinates

$$\begin{aligned} \dot{\mathbf{R}} &= \frac{B^*}{B_\parallel^*} v_\parallel + \frac{c}{q_\alpha B_\parallel^*} \mathbf{b} \times \left(\mu \nabla B + q_\alpha \nabla \langle \phi_1 \rangle_R \right) \\ \dot{v}_\parallel &= - \frac{B^*}{m_\alpha B_\parallel^*} \left(\mu \nabla B + q_\alpha \nabla \langle \phi_1 \rangle_R \right) - \frac{q_\alpha}{m_\alpha c} \frac{\partial \langle A_{1\parallel} \rangle}{\partial t} \end{aligned}$$

Gyro-averages and GENE-X equations

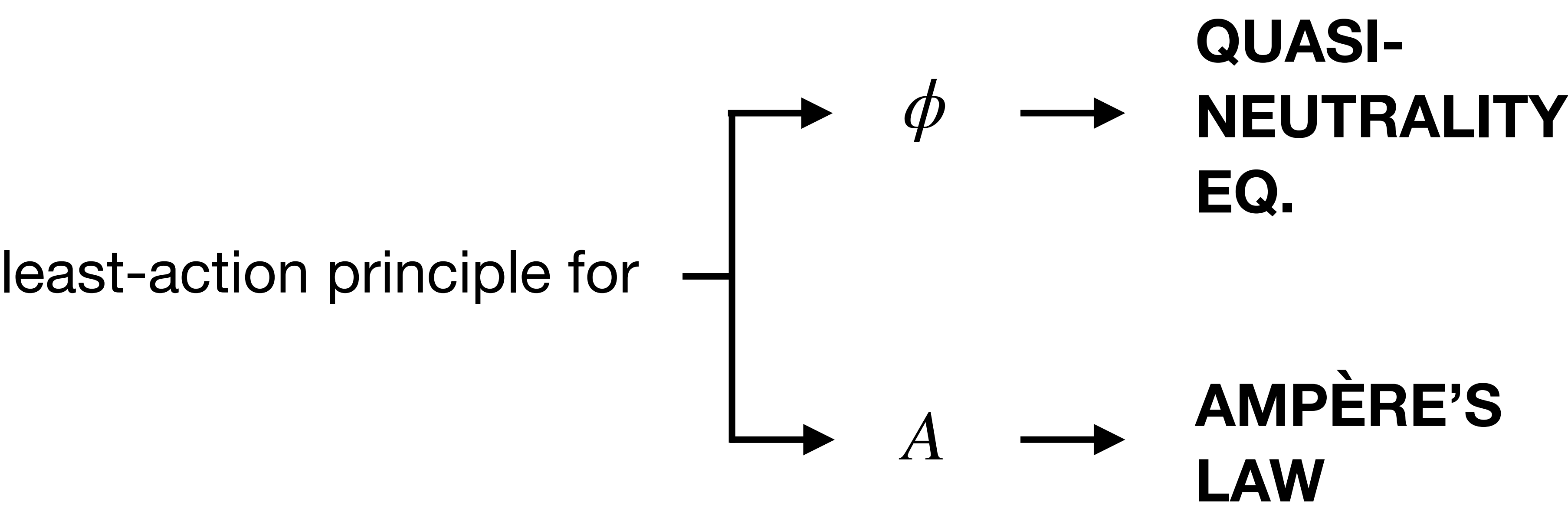
gyro-average

$$\langle \phi(\mathbf{r}) \rangle_{\mathbf{R}} = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_{\alpha} \cdot \nabla} \phi(\mathbf{R})$$

adjoint
operator

$$\langle \phi(\mathbf{R}) \rangle_{\mathbf{r}}^{\dagger} = \int_0^{2\pi} \frac{d\theta}{2\pi} \int d\mathbf{R} \, \delta(\mathbf{R} + \rho_{\alpha} - \mathbf{r}) \, \phi(\mathbf{R})$$

$$\mathcal{L} = \frac{q_{\alpha}}{c} \mathbf{A} \cdot \dot{\mathbf{R}} + \frac{q_{\alpha}}{c} \left\langle A_{1\parallel} \right\rangle_{\mathbf{R}} \mathbf{b} \cdot \dot{\mathbf{R}} + m_{\alpha} v_{\parallel} \mathbf{b} \cdot \dot{\mathbf{R}} + \frac{\mu B}{\Omega_{\alpha}} \dot{\theta} - q_{\alpha} \left\langle \phi_1 \right\rangle_{\mathbf{R}} - \frac{1}{2} m_{\alpha} v_{\parallel}^2 - \mu B + \frac{m_{\alpha} c^2}{2 B^2} \left| \nabla_{\perp} \phi_1 \right|^2$$



$$\sum_{\alpha} q_{\alpha} \left\langle n_{\alpha} \right\rangle_{\mathbf{r}}^{\dagger} = - \nabla \cdot \left(\sum_{\alpha} n_{\alpha} \frac{m_{\alpha} c^2}{B^2} \nabla_{\perp} \phi_1 \right)$$

$$\sum_{\alpha} \frac{q_{\alpha}}{c} \left\langle \int dW f_{\alpha} v_{\parallel} \right\rangle_{\mathbf{r}}^{\dagger} = - \frac{1}{4\pi} \left| \nabla_{\perp}^2 A_{1\parallel} \right|$$

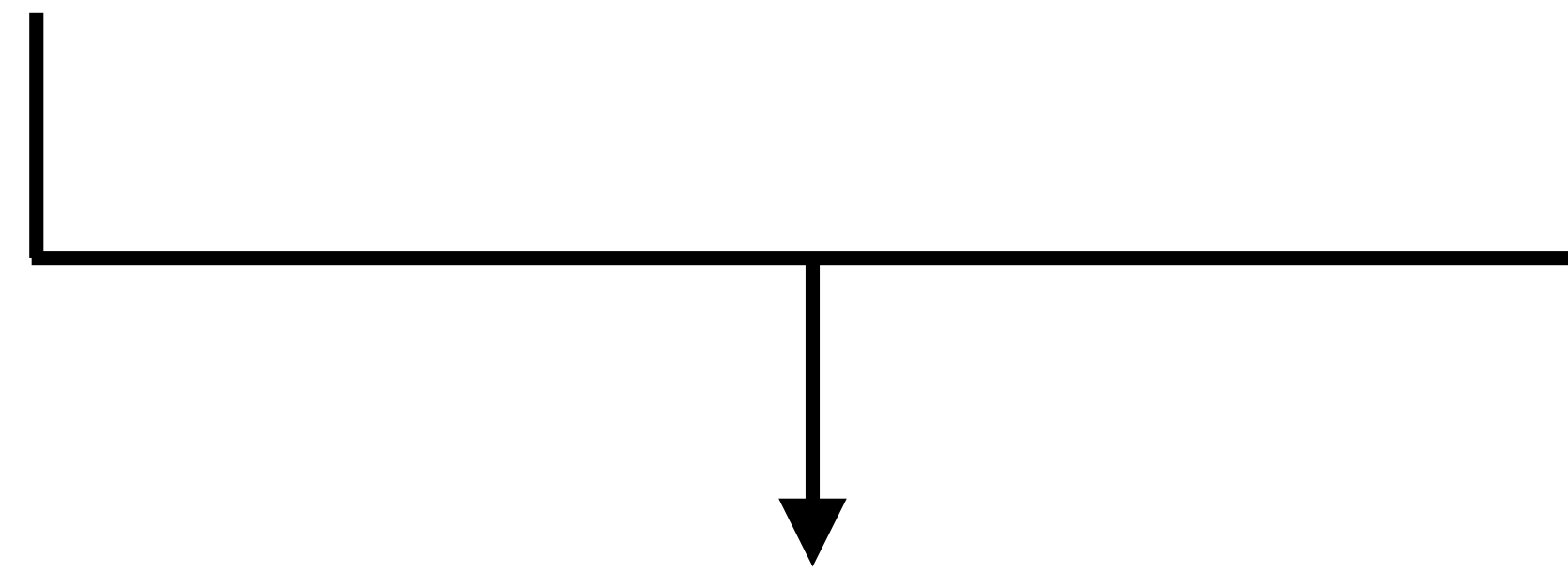
Gyro-averages and GENE-X equations

gyro-average

$$\langle \phi(\mathbf{r}) \rangle_{\mathbf{R}} = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_\alpha \cdot \nabla} \phi(\mathbf{R})$$

adjoint
operator

$$\langle \phi(\mathbf{R}) \rangle_{\mathbf{r}}^\dagger = \int_0^{2\pi} \frac{d\theta}{2\pi} \int d\mathbf{R} \delta(\mathbf{R} + \rho_\alpha - \mathbf{r}) \phi(\mathbf{R})$$



ENERGY CONSERVATION

$$\sum_{\alpha} \int dV \int dW q_{\alpha} f_{\alpha} \left\langle \frac{\partial \phi_1}{\partial t} - \frac{v_{\parallel}}{c} \frac{\partial A_{1\parallel}}{\partial t} \right\rangle_{\mathbf{R}} = \sum_{\alpha} \int dV q_{\alpha} \left(\left\langle \int dW f_{\alpha} \right\rangle_{\mathbf{r}}^\dagger \frac{\partial \phi_1}{\partial t} - \left\langle \int dW f_{\alpha} \frac{v_{\parallel}}{c} \right\rangle_{\mathbf{r}}^\dagger \frac{\partial A_{1\parallel}}{\partial t} \right)$$

Gyro-averages and GENE-X equations

gyro-average

$$\langle \phi(\mathbf{r}) \rangle_R = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{\rho_\alpha \cdot \nabla} \phi(\mathbf{R})$$

adjoint
operator

$$\langle \phi(\mathbf{R}) \rangle_r^\dagger = \int_0^{2\pi} \frac{d\theta}{2\pi} \int d\mathbf{R} \, \delta(\mathbf{R} + \rho_\alpha - \mathbf{r}) \, \phi(\mathbf{R})$$

[1]

$$\mathcal{L} \sim \left| \begin{array}{c} \phi_1 \\ A_{1\parallel} \end{array} \right|$$

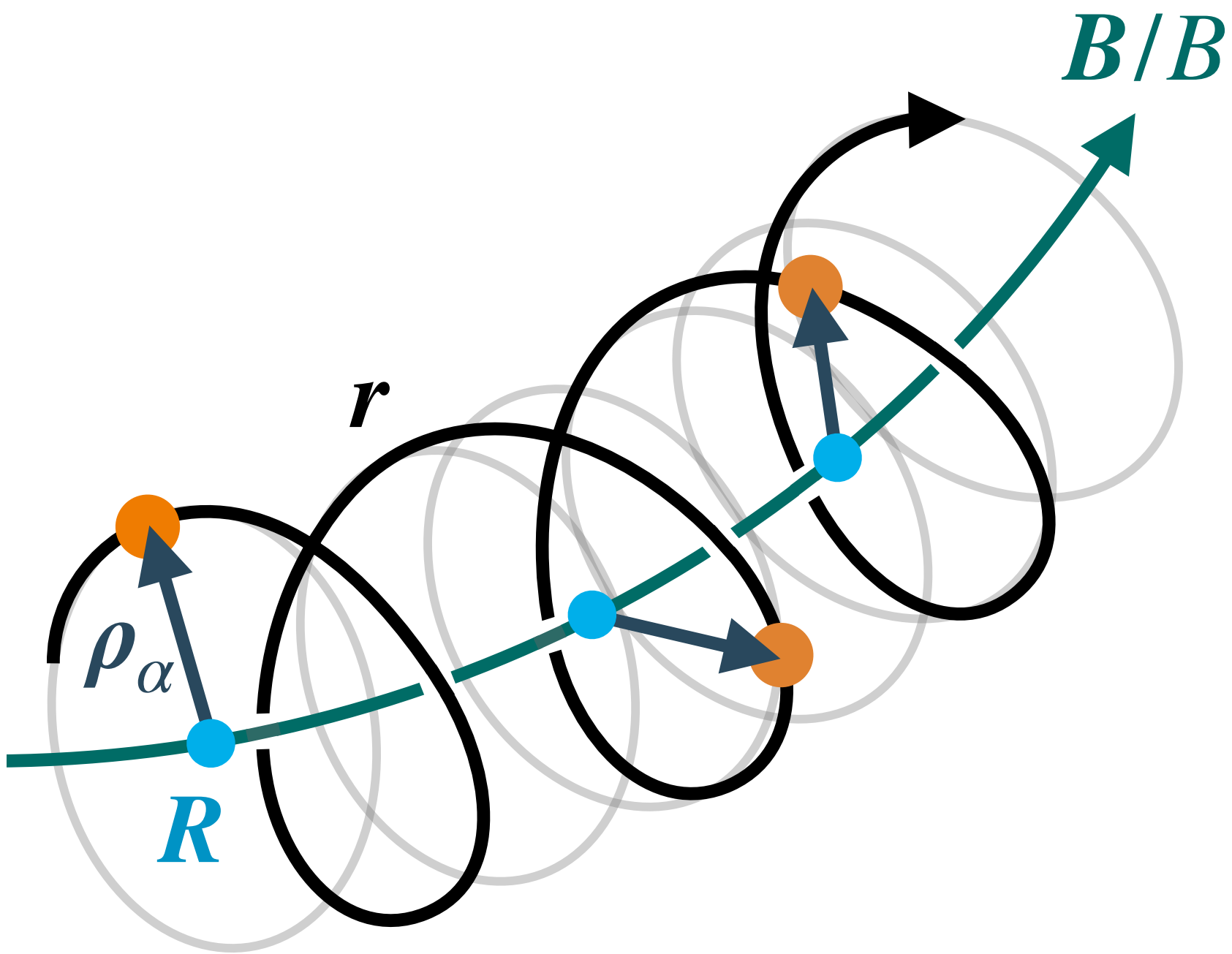
ENERGY CONSERVATION $f_1 \left(\langle \dots \rangle_R \right) = f_2 \left(\langle \dots \rangle_r^\dagger \right)$

VLASOV EQ. $\frac{\partial f_\alpha}{\partial t} + \dot{\mathbf{R}} \cdot \nabla f_\alpha + \dot{v}_\parallel \frac{\partial f_\alpha}{\partial v_\parallel} = 0$

$$\begin{aligned} \dot{\mathbf{R}} &\sim \langle \phi_1 \rangle_R \\ \dot{v}_\parallel &\sim \langle \phi_1 \rangle_R, \langle A_{1\parallel} \rangle \end{aligned}$$

least-action principle for

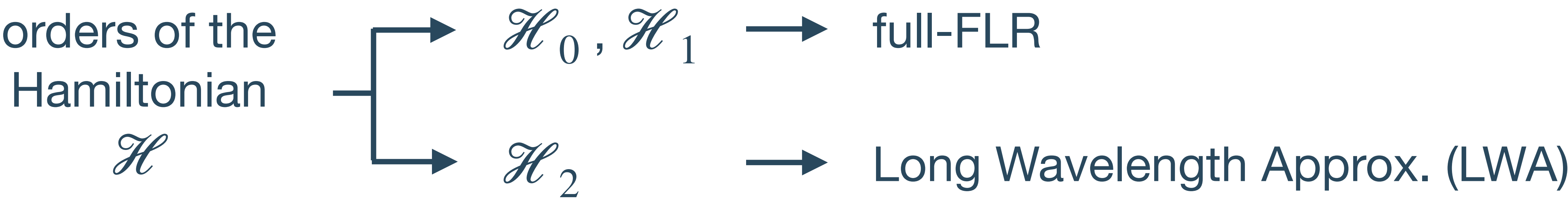
$$\begin{aligned} \phi &\rightarrow \text{QUASI-NEUTRALITY} \\ A &\rightarrow \text{AMPÈRE'S LAW} \end{aligned}$$



$$\sim \langle n_\alpha \rangle_r^\dagger$$

$$\sim \left\langle \int dW f_\alpha v_\parallel \right\rangle_r^\dagger$$

FLR... where?



FLR IN →	$\mathcal{H}_0, \mathcal{H}_1$	$\mathcal{H}_2$
GENE-X current	LWA	LWA
GENE-X target	FULL	LWA
FULL FLR	FULL	FULL

tested in the **electrostatic slab ion temperature gradient (ITG)** dispersion relation in Fourier space
with diffusion in $\parallel$ and $\perp$ dirs.

$$\sum_{\alpha} \frac{1}{T_{\alpha}} \Gamma_0 + \sum_{\alpha} \frac{1}{T_{\alpha}} \Gamma_0 \xi_{\text{eff}} Z - \sum_{\alpha} \frac{1}{\sqrt{2} T_{\alpha}} \frac{k_{\perp}}{k_{\parallel}} \left(\Gamma_0 Z + \eta \left(\Gamma_0 \xi_{\text{eff}} (1 + \xi_{\text{eff}} Z) - \frac{1}{2} \Gamma_0 Z + a (\Gamma_1 - \Gamma_0) Z \right) \right)$$

+ polarization term = 0

$$\begin{aligned} &\rightarrow \text{if full-FLR in } \mathcal{H}_2 \Rightarrow \equiv \sum_{\alpha} \tau_{\alpha}^{-1} (1 - \Gamma_0) \\ &\rightarrow \text{if LWA in } \mathcal{H}_2 \Rightarrow \equiv \sum_{\alpha} k_{\perp}^2 \end{aligned}$$

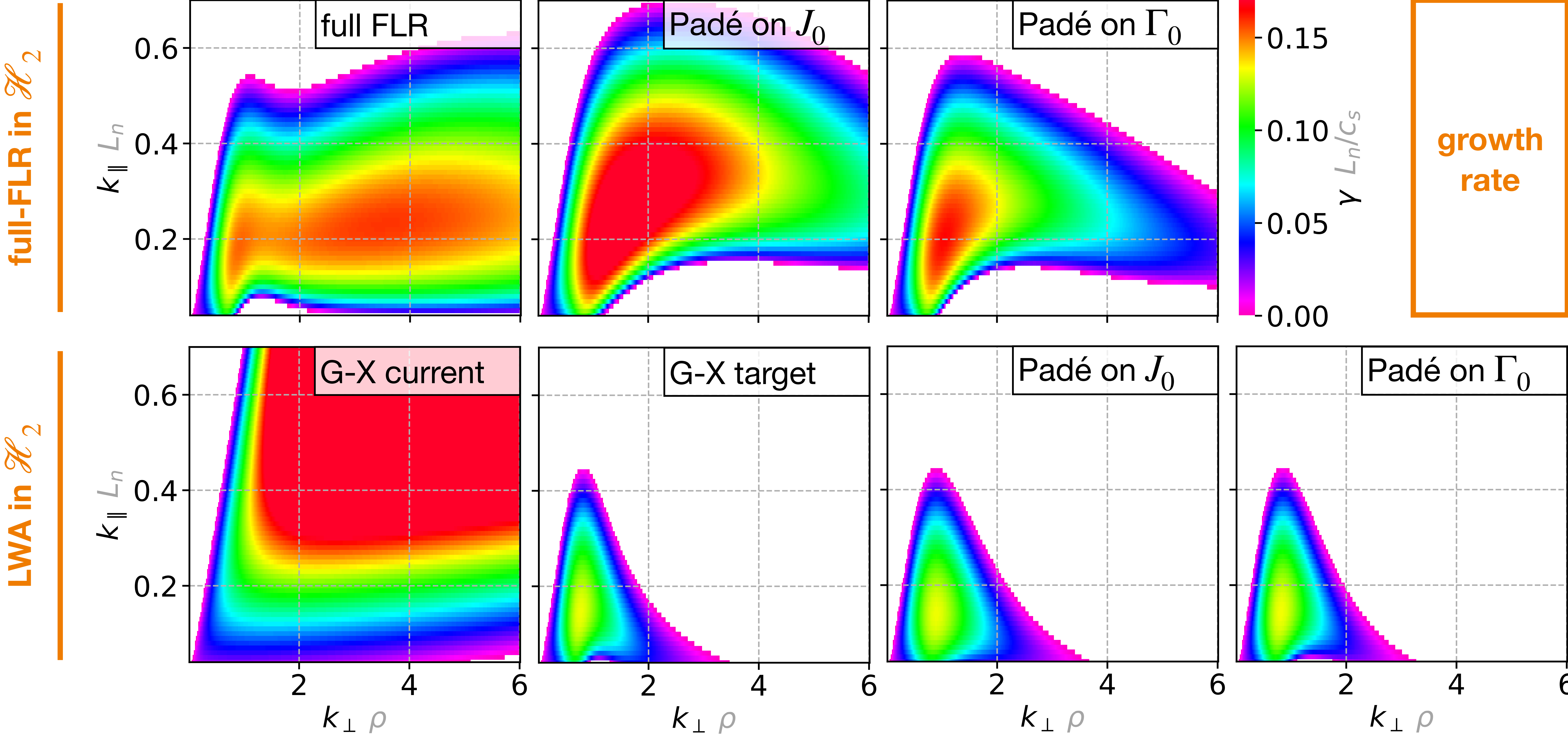
merging Vlasov eq. + quasi-neutrality eq.
 $\Rightarrow J_0$ respective operators combine to $\Gamma_0 \sim J_0^2$

Physics in Γ_0 !

Padé on J_0 from where we get Γ_0
or Padé on Γ_0 directly

FLR models

tested in the **electrostatic slab ion temperature gradient (ITG)**
dispersion relation in Fourier space **with diffusion** in $\parallel$ and $\perp$ dirs.

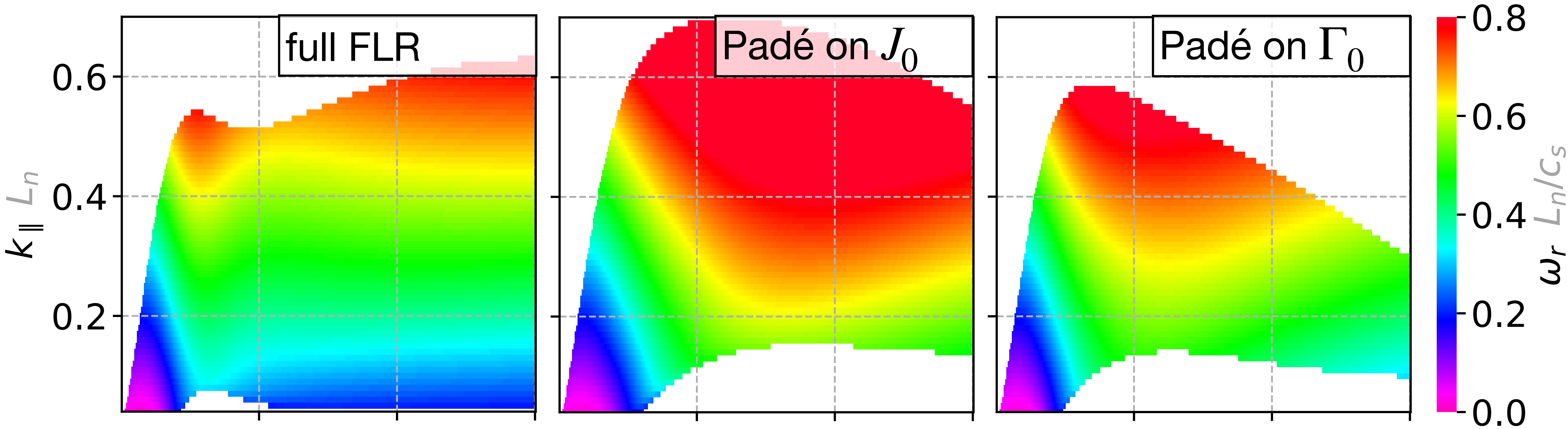


FLR models

tested in the **electrostatic slab ion temperature gradient (ITG)**
dispersion relation in Fourier space **with diffusion** in $\parallel$ and $\perp$ dirs.

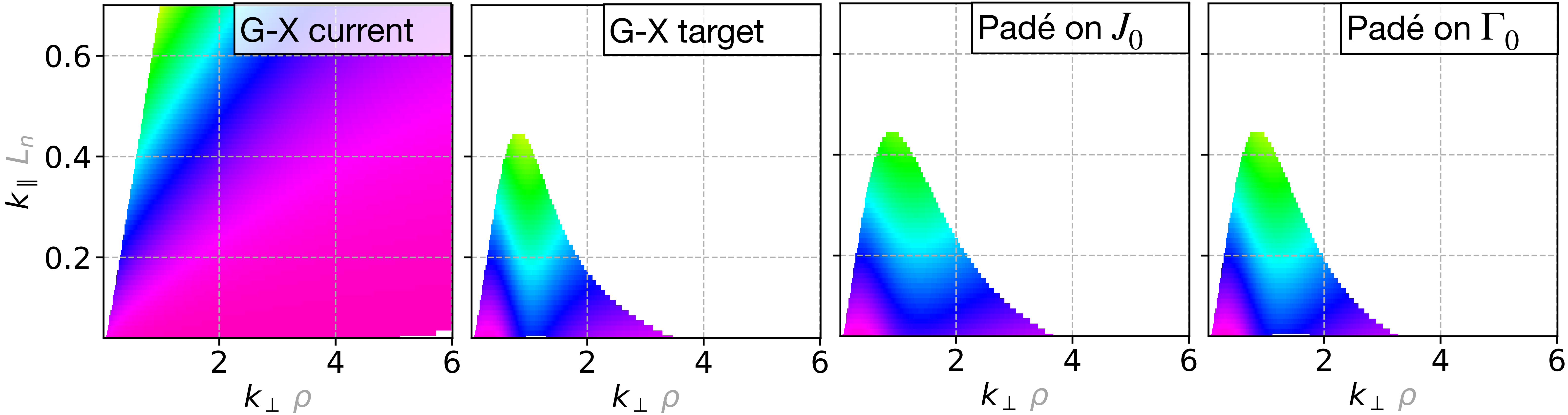


full-FLR in $\mathcal{H}_2$



real
freq.

LWA in $\mathcal{H}_2$



LWA in $\mathcal{H}_0, \mathcal{H}_1$

full-FLR in $\mathcal{H}_0, \mathcal{H}_1$

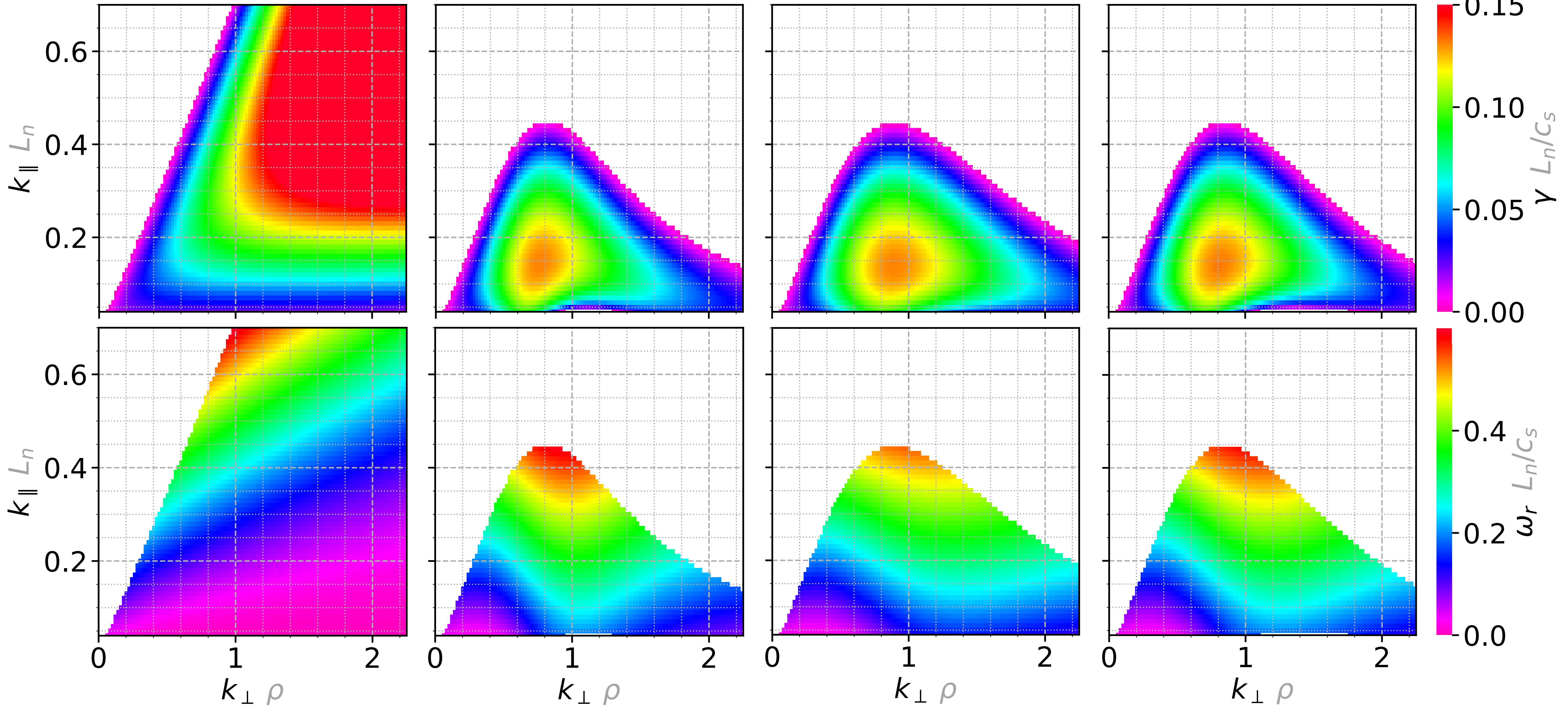
Padé

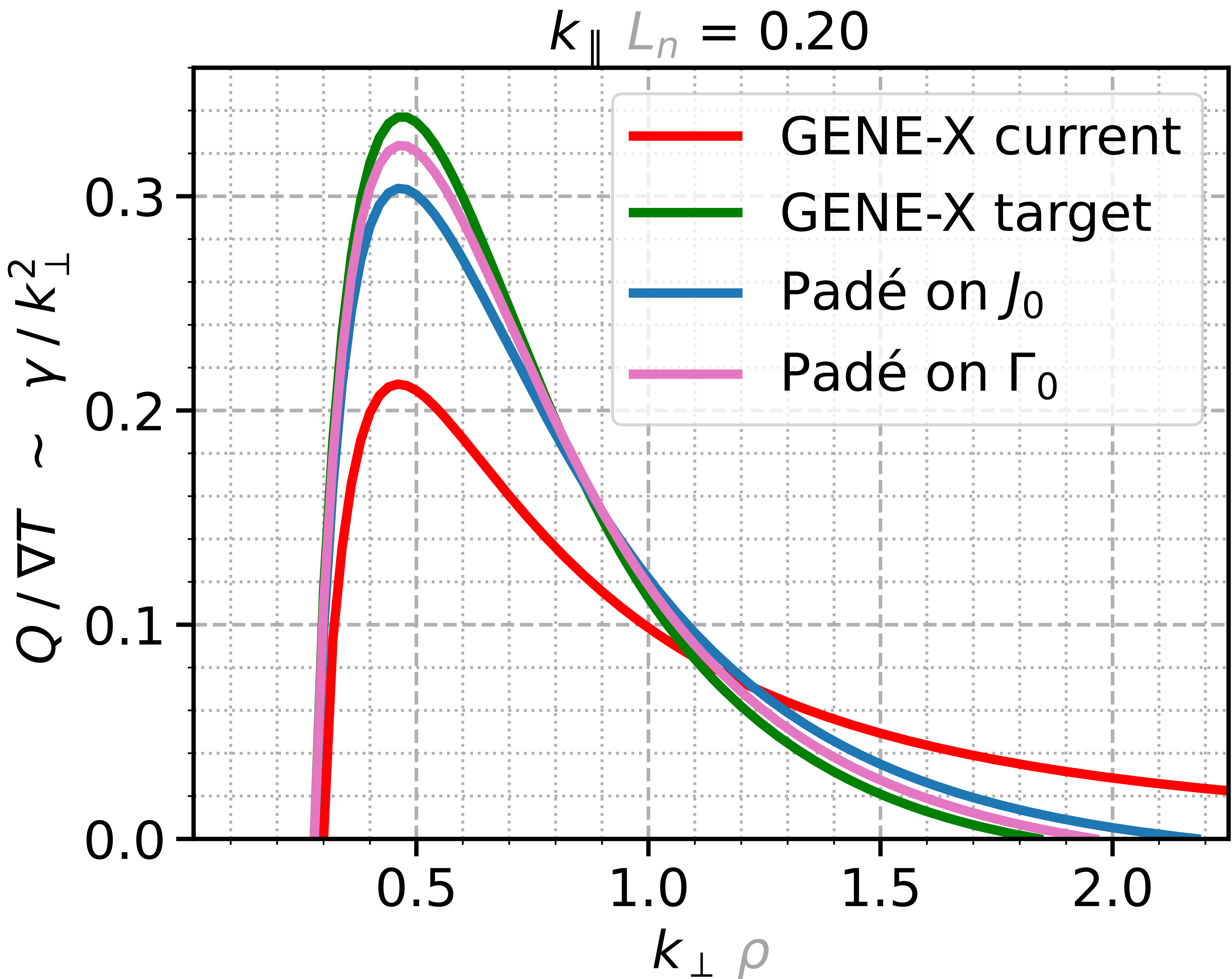
GENE-X current

GENE-X target

Padé on J_0

Padé on Γ_0





analysis of quasi-linear flux [1]

- major Physics change from GENE-X **current** to **target**
- Padé on Γ_0** more Physically sensible
- Padé on J_0** not far-fetched and provides very convenient algorithm

$$\left(1 - \frac{1}{4} \rho_{\alpha}^2 \nabla_{\perp}^2\right) \langle \phi(\mathbf{r}) \rangle_{\theta} = \phi(\mathbf{R})$$

differential eq. for gyro-average



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HEPP Introductory Talk



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