

Carbon content and transport investigations on W7-X with Charge Exchange Recombination Spectroscopy

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magnetic

field lines

magnetic

flux surfaces

Fusion on Earth

Most viable reaction for fusion reactors:

 $D + T \rightarrow {}^{4}He (3.52 MeV) + n (14.1 MeV)$



- For energy producing n_iT_iτ_E needs to be above a certain limit.
- → Magnetic confinement: torus shaped plasma, twisted magnetic lines

$$\label{eq:ni} \begin{split} n_i \propto 10^{19} - 10^{20} \; m^{-3} \\ T_i \propto 10 \; keV \propto 100 \; million \; ^\circ \mathrm{C} \end{split}$$



Wendelstein 7-X stellarator



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Impurities in the plasma:

- Entering the plasma from wall components.
- Performance losses by dilution and radiation. Need for control.
- Understanding of behavior is crucial.



- **Neoclassical transport**: collisional transport in toroidal geometry
- Anomalous transport: additional transport effects (e.g. turbulence)
- Neoclassical models predict impurity accumulation in the core plasma for W7-X, which can affect the performance through power losses. [Burhenn, Nucl. Fusion (2009) 065005]

Is impurity accumulation observed during specific plasma settings? What dominates the impurity transport during different configurations?



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In steady state scenario in the core: $\Gamma_{imp} = 0 \rightarrow \frac{v}{D} = \frac{1}{n_{imp}} \frac{\partial n_{imp}}{\partial r}$

Tool:

[B. Geiger, C. Swee]

1 dimensional impurity transport code pySTRAHL. Changing the input diffusion (D) and convection velocity (v) to match the output carbon density with the experimental profiles.



Impurity transport

•
$$\frac{\partial n_{imp}}{\partial t} = -\nabla \Gamma_{imp} + Q_{sources} + Q_{sinks}$$

• Impurity transport flux:
$$\Gamma_{imp} = -D \frac{\partial n_{imp}}{\partial r} + v n_{imp}$$



- Impurity properties obtained from measured spectral line.
- Coronal Equilibrium: models the temperature dependence of the ionisation stages.



Passive Spectroscopy - Measurement



- Measures the line integrated intensity of C⁵⁺ (529.07 nm) with multichannel mirror spectrometer.
- Inversion (Forward modelling) to assess radial distribution of emission intensity that gives comparable line integrated intensity.



Passive Spectroscopy - Inversion



- Every flux surface is represented by a value of ρ. Array of ρ determined for each LOS given the flux surfaces the LOS looks through.
- The modelled emissivity is parametrised to be a double Gaussian shape
- Atomic processes leading to C⁵⁺:

Electron impact excitation: $C^{5+} + e^- \rightarrow [C^{5+}]^* + e^-$

Charge exchange: $C^{6+} + H^0 \rightarrow [C^{5+}]^* + H^+$



Passive Spectroscopy – Calculation



- Inner peak is radially deeper than either of the source particles is expected.
- Strong radial anomalous transport could carry the C⁵⁺ ions further in before they are ionised.
- Preliminary radiation calculation supports this.



Charge eXchange Recombination Spectroscopy

- In the plasma core: only fully-ionized ions, no radiation to be detected
- Neutral Beam Injection (NBI): controllable external neutral source to the charge exchange process.

$$H^0 + C^{6+} \rightarrow H^+ + [C^{5+}]^*$$

- Spatially localised data available.
- Separation of the active and passive component is needed.



- Main spectrometer: ILS (C^{5+} , H^{α})
- This work focuses on carbon density calculated from the measured C⁵⁺ intensity.

Wendelstein

Carbon density calculation

• Calculation of the carbon density is possible using the H^{α} and C^{5+} intensities measured by the ILS.

$$I^{CX} \propto \int_{LOS} n_C \sum_{E=E,E/2,E/3} Q_E^{CX} n_E^{beam} dl \qquad I_E^{BES} \propto \int_{LOS} n_e Q_E^{BES} n_E^{beam} dl$$

Intersection integral between the NBI and the LOS can be eliminated, as the two detectors share the same LOS-s.

$$n_{C} = \frac{I^{CX}}{\sum_{E=E,E/2,E/3} \frac{I_{E}^{BES} Q_{E}^{CX}}{Q_{E}^{BES}}} n_{e}$$



Outline of results

- ECRH dominated plasma scenarios
- NBI dominated plasma scenarios
- Mixed heating plasma scenarios
- High performance plasmas



ECRH dominated plasma scenario





Flat carbon profile throughout the radius, no impurity accumulation.

ECRH dominated plasma scenario





- Strong anomalous diffusion was found in these plasmas based on pySTRAHL simulations.
 [Geiger, Nucl. Fusion (2019) 59 047009]
- This outward diffusion keeps the profiles flat against the neoclassical convection.

NBI dominated plasma scenario





- Heating switch at $t \approx 1s$
- ECRH turned back on at $t \approx 3s$
- n_e starts peaking in the core continuously.
- n_C peaking takes off around 1s
 into the pure NBI phase.

NBI dominated plasma scenario







- Flat n_C profiles in ECRH phase
- *n_C* peaking starts slowly after heating switch at 1s



pySTRAHL simulation:

- One simulation for the whole discharge
- Background kinetic profiles every 100ms (from experiment)
- Neoclassical D and v profiles corresponding to the given kinetic profiles (calculated from Neotransp)
- Anomalous D profile from ECRH plasma studies
- Edge and source parameters set to match n_C level at the edge of the confined plasma at the beginning of the experiment.

Simulation for ECRH phase





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Simulation for NBI phase with neoclassical D





Gives carbon peaking for $\rho < 0.5$ with similar quantitative increase in time as experiment.

Simulation for NBI phase with an. D for p>0.5





Profile flattens out for $\rho > 0.5$ and stops decreasing in time.

Simulation for NBI phase with an. D for ρ <0.5



- S: ratio of simulated diffusion level and full anomalous diffusion
- No significant peaking observed for S > 5%



Combined ECRH and NBI heating







- Carbon flushed out from the core
- Barrier-like gradient remains at $\rho = 0.4 0.5$
- Core turbulence becomes significant again.

High performance plasma scenario



Achieved with continuous pellet injections.

• Triple product:
$$n_i T_i \tau_E = \frac{n_i}{n_e} n_e T_i \frac{W_{dia}}{P_{heating}}$$

- Comparing the performance of different machines on the way to net energy producing fusion reactors.
- Reported highest triple product for stellarators:

$$n_i T_i \tau_E = 6.46 * 10^{19} \, m^{-3} keV \, s$$

[Pedersen, Plasma Phys. Control. Fusion (2019) 61 014035]

• To measure the amount of impurities: $Z_{eff} = \frac{\sum n_j Z_j^2}{\sum n_j Z_j}$

Supports the recorded record triple product for stellarators.

Triple product with core impurity measurements

$$T_i = 3.08 \ keV$$

 $W_{dia} = 1.14 \ MJ$
 $P_{heating} = 4.65 \ MW$
 $n_e = 8.55 * 10^{19} \ m^{-3}$

Measured global parameters:

20181016.037 0.8 = 1.0 s0.7 1.3 s 0.6 0.5 = 1.75 s[%] 0.4 ∪ 0.3 0.2 0.1 0.0 0.0 0.2 0.4 0.6 0.8 1.0 ρ

With measured carbon concentration:

$$C_C = 0.6\%$$

 $Z_{eff} = 1.17 \longrightarrow n_i T_i \tau_E = 6.26 * 10^{19} m^{-3} keV s$

Assuming same oxygen concentration:

$$C_C = 0.6\%, C_O = 0.6\% \longrightarrow n_i T_i \tau_E = 6.00 * 10^{19} m^{-3} keV s$$

 $Z_{eff} = 1.48$







- Impurity transport needs to be understood for their control to achieve reactor like plasma conditions.
- *n_C* was measured at W7-X by the CXRS diagnostic.
- Carbon transport simulations were done based on the experimental n_c profiles.
- Passive Spectroscopy: only measuring carbon at the plasma edge, line integrated
- **CXRS**: measuring carbon in the plasma core, spatially resolved
 - ECRH plasmas: turbulent dominated impurity transport
 - NBI plasmas: neoclassical impurity transport in the core, causes accumulation
 - Mixed heating plasma: turbulent core keeps the accumulation lower
 - High performance plasmas: n_C profiles used for validation